

# Chapter 1

## Basics of statistical physics

$$\dot{\mathbf{r}}_i = \frac{\partial \mathcal{H}(\mathbf{q})}{\partial \mathbf{p}_i} \quad \text{and} \quad \dot{\mathbf{p}}_i = -\frac{\partial \mathcal{H}(\mathbf{q})}{\partial \mathbf{r}_i}, \quad \forall i = 1 \dots N, \quad (1.1)$$

### 1.1 Liouville's theorem

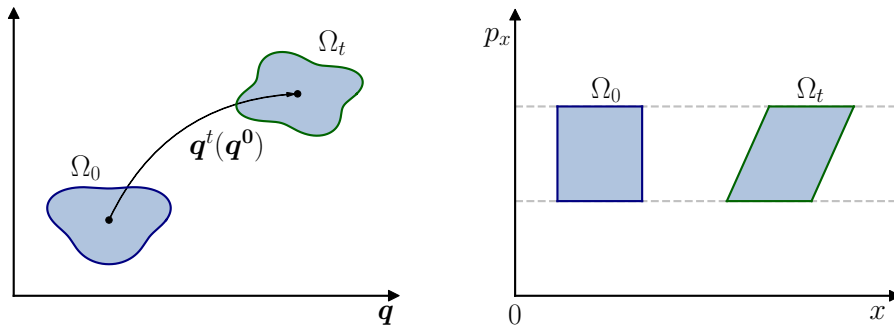
$$\int_{\Omega_0} d^{6N} \mathbf{q} = \int_{\Omega_t} d^{6N} \mathbf{q} = \int_{\Omega_t} d^{6N} \mathbf{q}^t. \quad (1.2)$$

$$\int_M f(\mathbf{y}) d^N \mathbf{y} = \int_{\phi^{-1}(M)} f(\phi(\mathbf{x})) \left| \det \frac{D\phi}{Dx} \right| d^N \mathbf{x}. \quad (1.3)$$

$$\frac{D\phi}{Dx} = \begin{bmatrix} \frac{\partial \phi_1}{\partial x_1} & \frac{\partial \phi_2}{\partial x_1} & \cdots & \frac{\partial \phi_N}{\partial x_1} \\ \frac{\partial \phi_1}{\partial x_2} & \frac{\partial \phi_2}{\partial x_2} & \cdots & \frac{\partial \phi_N}{\partial x_2} \\ \vdots & \vdots & \ddots & \vdots \end{bmatrix}. \quad (1.4)$$

$$\int_{\Omega_t} d^{6N} \mathbf{q} = \int_{\Omega_t} d^{6N} \mathbf{q}^t = \int_{\Omega_0} \left| \det \frac{Dq^t}{Dq^0} \right| d^{6N} \mathbf{q}^0. \quad (1.5)$$

$$\int_{\Omega_t} 1 d^{6N} \mathbf{q} = \int_{\Omega_0} 1 d^{6N} \mathbf{q}. \quad (1.6)$$



$$J^t \equiv \left| \det \frac{Dq^t}{Dq^0} \right|, \quad (1.7)$$

$$\frac{dJ_{ij}^t}{dt} = \frac{\partial}{\partial q_j^0} \frac{dq_i^t}{dt} = \frac{\partial \dot{q}_i^t}{\partial q_j^0} = \sum_{k=1}^{6N} \frac{\partial \dot{q}_i^t}{\partial q_k^t} \frac{\partial q_k^t}{\partial q_j^0} = \sum_{k=1}^{6N} \frac{\partial \dot{q}_i^t}{\partial q_k^t} J_{kj}^t, \quad (1.8)$$

$$J_{kj}^t = \frac{\partial q_k^t}{\partial q_j^0}. \quad (1.9)$$

$$\frac{dJ^t}{dt} = \sum_{i,j} \frac{\partial J^t}{\partial J_{ij}^t} \frac{dJ_{ij}^t}{dt} = \sum_{i,j,k} \frac{\partial J^t}{\partial J_{ij}^t} J_{kj}^t \frac{\partial \dot{q}_i^t}{\partial q_k^t}. \quad (1.10)$$

$$\sum_j \frac{\partial J^t}{\partial J_{ij}^t} J_{kj}^t = \delta_{ik} J^t \quad (1.11)$$

$$\frac{dJ^t}{dt} = J^t \sum_{i=1}^{6N} \frac{\partial \dot{q}_i^t}{\partial q_i^t}. \quad (1.12)$$

$$\frac{dJ^t}{dt} = J^t \sum_{i=1}^{3N} \left( \frac{\partial r_i^t}{\partial r_i^t} + \frac{\partial p_i^t}{\partial p_i^t} \right) = J^t \sum_{i=1}^{3N} \left( \frac{\partial^2 \mathcal{H}}{\partial r_i^t \partial p_i^t} - \frac{\partial^2 \mathcal{H}}{\partial p_i^t \partial r_i^t} \right) \quad (1.13)$$

$$B = \frac{1}{\det A} \text{adj} A, \quad (1.14)$$

$$(\text{adj} A)^{ij} \equiv (-1)^{i+j} \det(A^{ji}). \quad (1.15)$$

$$b_{ij} = \frac{(-1)^{i+j}}{a} \det(A^{ji}),$$

$$a = \sum_i (-1)^{i+j} a_{ij} \det A^{ij}.$$

$$\frac{\partial a}{\partial a_{ij}} = (-1)^{i+j} \det A^{ij} = ab^{ji}. \quad (1.16)$$

$$\sum_j \frac{\partial a}{\partial a_{ij}} a_{kj} = \sum_j ab^{ji} a_{kj} = a \sum_j b^{ji} a_{kj} = a(BA)_{ki} = a\delta_{ik}, \quad (1.17)$$

## 1.2 Statistical description of a physical system

$$\int_{-\infty}^{+\infty} D_N d\omega = \mathcal{C}, \quad (1.18)$$

$$\frac{\partial}{\partial t} \int_{\Omega} D_N d\omega = - \oint_{\partial\Omega} D_N \dot{\mathbf{q}} \cdot d\mathbf{\Sigma} \quad (1.19)$$

$$\int_{\Omega} \frac{\partial D_N}{\partial t} d\omega = - \int_{\Omega} \nabla \cdot (D_N \dot{\mathbf{q}}) d\omega. \quad (1.20)$$

$$\frac{\partial D_N}{\partial t} = -\nabla \cdot (D_N \dot{\mathbf{q}}) . \quad (1.21)$$

$$\frac{\partial D_N}{\partial t} = -\sum_{i=1}^{6N} \left[ \frac{\partial D_N}{\partial q_i} \dot{q}_i + D_N \frac{\partial \dot{q}_i}{\partial q_i} \right] = -\sum_{i=1}^{3N} \left[ \frac{\partial D_N}{\partial r_i} \frac{\partial \mathcal{H}}{\partial p_i} - \frac{\partial D_N}{\partial p_i} \frac{\partial \mathcal{H}}{\partial r_i} + D_N \left( \frac{\partial \mathcal{H}}{\partial r_i \partial p_i} - \frac{\partial \mathcal{H}}{\partial p_i \partial r_i} \right) \right] . \quad (1.22)$$

$$\frac{\partial D_N}{\partial t} = \{\mathcal{H}, D_N\}_P . \quad (1.23)$$

$$f^1(\mathbf{r}_1, \mathbf{p}_1) = \int D_N(\mathbf{r}_1, \mathbf{p}_1, \dots, \mathbf{r}_N, \mathbf{p}_N) d^3 \mathbf{r}_2 d^3 \mathbf{p}_2 \dots d^3 \mathbf{r}_N d^3 \mathbf{p}_N .$$

$$f^2(\mathbf{r}_2, \mathbf{p}_2) = \int D_N(\mathbf{r}_1, \mathbf{p}_1, \dots, \mathbf{r}_N, \mathbf{p}_N) d^3 \mathbf{r}_1 d^3 \mathbf{p}_1 d^3 \mathbf{r}_3 d^3 \mathbf{p}_3 \dots d^3 \mathbf{r}_N d^3 \mathbf{p}_N$$

$$f(\mathbf{r}, \mathbf{p}) = \sum_{i=1}^N f^i(\mathbf{r}, \mathbf{p}) = N f^1(\mathbf{r}, \mathbf{p}) = N f^2(\mathbf{r}, \mathbf{p}) = \dots$$

### 1.3 Reduced distribution functions

$$f(\mathbf{r}, \mathbf{p}) \equiv \int D_N(\mathbf{r}, \mathbf{p}, \mathbf{r}_2, \mathbf{p}_2, \dots, \mathbf{r}_N, \mathbf{p}_N) d\tau_{N-1} , \quad (1.24)$$

$$d\tau_{N-1} \equiv \frac{d^3 \mathbf{r}_2 d^3 \mathbf{p}_2 \dots d^3 \mathbf{r}_N d^3 \mathbf{p}_N}{(N-1)!} \quad (1.25)$$

$$A(\mathbf{r}_1, \mathbf{p}_1, \dots, \mathbf{r}_N, \mathbf{p}_N) = \sum_{i=1}^N a(\mathbf{r}_i, \mathbf{p}_i) , \quad (1.26)$$

$$\langle A \rangle \equiv \int A D_N d\tau_N = \int a(\mathbf{r}, \mathbf{p}) f(\mathbf{r}, \mathbf{p}) d^3 \mathbf{r} d^3 \mathbf{p} . \quad (1.27)$$

$$f_2(\mathbf{r}_1, \mathbf{p}_1, \mathbf{r}_2, \mathbf{p}_2) \equiv \int D_N d\tau_{N-2} , \quad (1.28)$$

$$B(\mathbf{r}_1, \mathbf{p}_1, \dots, \mathbf{r}_N, \mathbf{p}_N) = \sum_{i=1}^N \sum_{j=1, j \neq i}^N b(\mathbf{r}_i, \mathbf{p}_i, \mathbf{r}_j, \mathbf{p}_j) , \quad (1.29)$$

$$\langle B \rangle = \int b(\mathbf{r}_1, \mathbf{p}_1, \mathbf{r}_2, \mathbf{p}_2) f_2(\mathbf{r}_1, \mathbf{p}_1, \mathbf{r}_2, \mathbf{p}_2) d^3 \mathbf{r}_1 d^3 \mathbf{p}_1 d^3 \mathbf{r}_2 d^3 \mathbf{p}_2 . \quad (1.30)$$

$$\mathcal{H} = \sum_{i=1}^N \frac{\mathbf{p}_i^2}{2m} + \sum_{i=1}^N V(\mathbf{r}_i) + \sum_{i=1}^N \sum_{\substack{j=1 \\ j \neq i}}^N w(\mathbf{r}_i, \mathbf{r}_j) . \quad (1.31)$$

$$\langle \mathcal{H} \rangle = \int h(\mathbf{r}, \mathbf{p}) f(\mathbf{r}, \mathbf{p}) d^3 \mathbf{r} d^3 \mathbf{p} + \int w(\mathbf{r}_1, \mathbf{r}_2) f_2(\mathbf{r}_1, \mathbf{p}_1, \mathbf{r}_2, \mathbf{p}_2) d^3 \mathbf{r}_1 d^3 \mathbf{p}_1 d^3 \mathbf{r}_2 d^3 \mathbf{p}_2 , \quad (1.32)$$

$$h(\mathbf{r}, \mathbf{p}) \equiv \frac{p^2}{2m} + V(\mathbf{r}) . \quad (1.33)$$

$$w(\mathbf{r}_i, \mathbf{r}_j) = -\frac{1}{2} \frac{Gm^2}{|\mathbf{r}_i - \mathbf{r}_j|} \quad (1.34)$$

## 1.4 Boltzmann equation

$$\int \frac{\partial D_N}{\partial t} d\tau_{N-1} = \int \sum_{i=1}^N \left( \frac{\partial D_N}{\partial \mathbf{p}_i} \cdot \frac{\partial h}{\partial \mathbf{r}_i} - \frac{\partial D_N}{\partial \mathbf{r}_i} \cdot \frac{\partial h}{\partial \mathbf{p}_i} \right) d\tau_{N-1} + \int \sum_{i,j} \frac{\partial w(\mathbf{r}_i, \mathbf{r}_j)}{\partial \mathbf{r}_i} \cdot \frac{\partial D_N}{\partial \mathbf{p}_i} d\tau_{N-1} . \quad (1.35)$$

$$\int \frac{\partial D_N}{\partial t} d\tau_{N-1} = \frac{\partial}{\partial t} \int D_N d\tau_{N-1} = \frac{\partial f(\mathbf{r}, \mathbf{p})}{\partial t} . \quad (1.36)$$

$$\int \frac{\partial D_N}{\partial \mathbf{p}} \cdot \frac{\partial h}{\partial \mathbf{r}} d\tau_{N-1} - \int \frac{\partial D_N}{\partial \mathbf{r}} \cdot \frac{\partial h}{\partial \mathbf{p}} d\tau_{N-1} = \frac{\partial h}{\partial \mathbf{r}} \cdot \frac{\partial f}{\partial \mathbf{p}} - \frac{\partial h}{\partial \mathbf{p}} \cdot \frac{\partial f}{\partial \mathbf{r}} . \quad (1.37)$$

$$I_1 \equiv \int \int \left( \frac{\partial D_N}{\partial \mathbf{p}_2} \cdot \frac{\partial h}{\partial \mathbf{r}_2} - \frac{\partial D_N}{\partial \mathbf{r}_2} \cdot \frac{\partial h}{\partial \mathbf{p}_2} \right) d^3 \mathbf{r}_2 d^3 \mathbf{p}_2 . \quad (1.38)$$

$$\frac{\partial f}{\partial t} + \frac{\partial h}{\partial \mathbf{p}} \cdot \frac{\partial f}{\partial \mathbf{r}} - \frac{\partial h}{\partial \mathbf{r}} \cdot \frac{\partial f}{\partial \mathbf{p}} = \int \sum_{i,j} \frac{\partial w(\mathbf{r}_i, \mathbf{r}_j)}{\partial \mathbf{r}_i} \cdot \frac{\partial D_N}{\partial \mathbf{p}_i} d\tau_{N-1} . \quad (1.39)$$

$$\frac{\partial f}{\partial t} + \frac{\partial h}{\partial \mathbf{p}} \cdot \frac{\partial f}{\partial \mathbf{r}} - \frac{\partial h}{\partial \mathbf{r}} \cdot \frac{\partial f}{\partial \mathbf{p}} = \left( \frac{\partial f}{\partial t} \right)_{\text{coll}} . \quad (1.40)$$

$$f_2(\mathbf{r}, \mathbf{p}, \mathbf{r}', \mathbf{p}') = f(\mathbf{r}, \mathbf{p}) f(\mathbf{r}', \mathbf{p}') + g_2(\mathbf{r}, \mathbf{p}, \mathbf{r}', \mathbf{p}') . \quad (1.41)$$

$$\Phi(\mathbf{r}) \equiv \langle W \rangle(\mathbf{r}) = \int w(\mathbf{r}, \mathbf{r}') f(\mathbf{r}', \mathbf{p}') d^3 \mathbf{r}' d^3 \mathbf{p}' = - \int \frac{Gm^2}{|\mathbf{r} - \mathbf{r}'|} f(\mathbf{r}', \mathbf{p}') d^3 \mathbf{r}' d^3 \mathbf{p}' \quad (1.42)$$

$$\frac{\partial f}{\partial t} + \frac{\partial h}{\partial \mathbf{p}} \cdot \frac{\partial f}{\partial \mathbf{r}} - \nabla [V(\mathbf{r}) + \Phi(\mathbf{r})] \cdot \frac{\partial f}{\partial \mathbf{p}} = 0 , \quad (1.43)$$

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - \frac{1}{m} \nabla [V(\mathbf{r}) + \Phi(\mathbf{r})] \cdot \frac{\partial f}{\partial \mathbf{v}} = 0 . \quad (1.44)$$

$$\frac{\partial f(\mathbf{r}, \mathbf{v})}{\partial t} + \mathbf{v} \cdot \frac{\partial f(\mathbf{r}, \mathbf{v})}{\partial \mathbf{r}} + \frac{\mathbf{F}(\mathbf{r})}{m} \cdot \frac{\partial f(\mathbf{r}, \mathbf{v})}{\partial \mathbf{v}} = 0 \quad (1.45)$$

## 1.5 Fluid equations

$$\begin{pmatrix} \rho(\mathbf{r}; t) \\ \rho(\mathbf{r}; t) \mathbf{u}(\mathbf{r}; t) \\ \rho(\mathbf{r}; t) \mathcal{E}(\mathbf{r}; t) \end{pmatrix} \equiv \int \begin{pmatrix} m \\ m \mathbf{v} \\ \frac{1}{2} m |\mathbf{v} - \mathbf{u}|^2 \end{pmatrix} f(\mathbf{r}; \mathbf{t}, \mathbf{v}) d^3 \mathbf{v} . \quad (1.46)$$

$$\int m \frac{\partial f}{\partial t} d^3 \mathbf{v} + \int m \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} d^3 \mathbf{v} - \int \nabla [V(\mathbf{r}) + \Phi(\mathbf{r})] \cdot \frac{\partial f}{\partial \mathbf{v}} d^3 \mathbf{v} = 0 . \quad (1.47)$$

$$\frac{\partial}{\partial t} \int m f d^3 \mathbf{v} + \frac{\partial}{\partial \mathbf{r}} \cdot \int m \mathbf{v} f d^3 \mathbf{v} - \nabla [V(\mathbf{r}) + \Phi(\mathbf{r})] \cdot \int \frac{\partial f}{\partial \mathbf{v}} d^3 \mathbf{v} = 0 . \quad (1.48)$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0 . \quad (1.49)$$

$$\int m v_i \frac{\partial f}{\partial t} d^3 \mathbf{v} + \int m v_i v_j \frac{\partial f}{\partial r_j} d^3 \mathbf{v} - \int v_i \frac{\partial}{\partial r_j} [V(\mathbf{r}) + \Phi(\mathbf{r})] \frac{\partial f}{\partial v_j} d^3 \mathbf{v} = 0 . \quad (1.50)$$

$$\frac{\partial \rho u_i}{\partial t} + \frac{\partial}{\partial r_j} (\rho \langle v_i v_j \rangle) - \frac{\partial}{\partial r_j} [V(\mathbf{r}) + \Phi(\mathbf{r})] \int v_i \frac{\partial f}{\partial v_j} d^3 \mathbf{v} = 0 , \quad (1.51)$$

$$\int v_i \frac{\partial f}{\partial v_j} d^3 \mathbf{v} = -\frac{\rho}{m} \delta_{ij} \quad (1.52)$$

$$\langle v_i v_j \rangle = u_i u_j + \langle w_i w_j \rangle ,^1 \quad (1.53)$$

$$\frac{\partial}{\partial t} (\rho u_i) + \frac{\partial}{\partial r_j} (\rho u_i u_j + P \delta_{ij} - \pi_{ij}) = -\frac{\rho}{m} \frac{\partial}{\partial r_i} [V(\mathbf{r}) + \Phi(\mathbf{r})] \quad (1.54)$$

$$u_i \frac{\partial \rho}{\partial t} + u_i \frac{\partial}{\partial r_j} (\rho u_j) + \rho \frac{\partial u_i}{\partial t} + \rho u_j \frac{\partial u_i}{\partial r_j} = -\frac{\rho}{m} \frac{\partial}{\partial r_i} [V(\mathbf{r}) + \Phi(\mathbf{r})] - \frac{\partial P}{\partial r_i} + \frac{\partial \pi_{ij}}{\partial r_j} . \quad (1.55)$$

$$\frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial r_j} = -\frac{1}{m} \frac{\partial}{\partial r_i} [V(\mathbf{r}) + \Phi(\mathbf{r})] - \frac{1}{\rho} \frac{\partial P}{\partial r_i} + \frac{1}{\rho} \frac{\partial \pi_{ij}}{\partial r_j} . \quad (1.56)$$

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = \mathbf{f} - \frac{1}{\rho} \nabla P + \frac{1}{\rho} \nabla \cdot \boldsymbol{\pi} , \quad (1.57)$$

$$\frac{d\mathbf{u}}{dt} = \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} , \quad (1.58)$$

$$\frac{\partial}{\partial t} [\tfrac{1}{2} \rho (u^2 + \langle w^2 \rangle)] + \frac{\partial}{\partial r_i} [\tfrac{1}{2} \rho \langle (u_i + w_i) |\mathbf{u} + \mathbf{w}|^2 \rangle] = -\frac{\rho}{m} u_i \frac{\partial}{\partial r_i} (V + \Phi) , \quad (1.59)$$

$$\frac{\partial}{\partial t} (\tfrac{1}{2} \rho u^2 + \rho \mathcal{E}) + \frac{\partial}{\partial r_i} [\tfrac{1}{2} \rho u^2 u_i + u_j (P \delta_{ij} - \pi_{ij}) + \rho \mathcal{E} u_i + F_i] = -\frac{\rho}{m} u_i \frac{\partial}{\partial r_i} (V + \Phi) , \quad (1.60)$$

$$\frac{\partial}{\partial t} (\tfrac{1}{2} \rho u^2) + \frac{\partial}{\partial r_i} (\tfrac{1}{2} \rho u^2 u_i) = -\frac{\rho}{m} u_i \frac{\partial}{\partial r_i} (V + \Phi) - u_i \frac{\partial P}{\partial r_i} + u_i \frac{\partial \pi_{ij}}{\partial r_j} , \quad (1.61)$$

$$\frac{\partial}{\partial t} (\rho \mathcal{E}) + \frac{\partial}{\partial r_i} (\rho \mathcal{E} u_i) = -P \frac{\partial u_i}{\partial r_i} - \frac{\partial F_i}{\partial r_i} + \Psi , \quad (1.62)$$

$$\rho \frac{d\mathcal{E}}{dt} = -P \nabla \cdot \mathbf{u} - \nabla \cdot \mathbf{F} + \Psi . \quad (1.63)$$

$$\pi_{ij} = \mu D_{ij} = \mu \left[ \frac{\partial u_i}{\partial r_j} + \frac{\partial u_j}{\partial r_i} - \frac{3}{2} \frac{\partial u_k}{\partial r_k} \delta_{ij} \right] , \quad (1.64)$$

$$\mathbf{F} = -\kappa \nabla T . \quad (1.65)$$

$$\Delta \Phi = 4\pi G \rho , \quad (1.66)$$

$$\Delta V = 4\pi G \rho_{\text{ext}} , \quad (1.67)$$

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<sup>1</sup>  $\langle v_i v_j \rangle = \frac{m}{\rho} \int d^3 \mathbf{v} f v_i v_j = \frac{m}{\rho} \int d^3 \mathbf{v} f (u_i + w_i)(u_j + w_j) = \frac{m}{\rho} \int d^3 \mathbf{v} [f u_i u_j + f u_i w_j + f u_j w_i + w_i w_j]$ . Considering that  $u_i$  may be taken out of the integrals and  $\int w_i f d^3 \mathbf{v} = \int (v_i - u_i) f d^3 \mathbf{v} = \int v_i f d^3 \mathbf{v} - u_i \int f d^3 \mathbf{v} = \frac{\rho}{m} u_i - u_i \frac{\rho}{m} = 0$  and using the definition of the mean value, we have  $\langle v_i v_j \rangle = u_i u_j + \langle w_i w_j \rangle$ .

## 1.6 Thermodynamical equilibria

$$S(D) \equiv -K_B \text{Tr} (D \ln D) , \quad (1.68)$$

$$\delta \left\{ \frac{1}{K_B} S(D) - \sum_{i=0}^n \lambda_i [\text{Tr}(A_i D) - \langle A_i \rangle] \right\} = 0 , \quad (1.69)$$

$$-\text{Tr} [(\ln D + \mathbb{1}) \delta D] - \lambda_0 \text{Tr}(\delta D) - \sum_{i=1}^n \lambda_i \text{Tr}(A_i \delta D) = 0 , \quad (1.70)$$

$$-\text{Tr} [\ln(D) \delta D] - \lambda'_0 \text{Tr}(\delta D) - \sum_{i=1}^n \lambda_i \text{Tr}(A_i \delta D) = 0 , \quad (1.71)$$

$$\text{Tr} \left[ \left( -\ln D - \sum_{i=0}^n \lambda_i A_i \right) \delta D \right] = 0 . \quad (1.72)$$

$$D = \exp \left[ - \sum_{i=0}^n \lambda_i A_i \right] . \quad (1.73)$$

$$D = e^{-\lambda_0} \exp \left[ - \sum_{i=1}^n \lambda_i A_i \right] . \quad (1.74)$$

$$e^{-\lambda_0} \text{Tr} \left\{ \exp \left[ - \sum_{i=1}^n \lambda_i A_i \right] \right\} = 1 . \quad (1.75)$$

$$Z \equiv e^{\lambda_0} = \text{Tr} \left\{ \exp \left[ - \sum_{i=1}^n \lambda_i A_i \right] \right\} . \quad (1.76)$$

$$D = \frac{1}{Z} \exp \left[ - \sum_{i=1}^n \lambda_i A_i \right] . \quad (1.77)$$

$$\frac{\partial \ln Z}{\partial \lambda_i} = -\langle A_i \rangle \quad \text{and} \quad \frac{\partial^2 \ln Z}{\partial \lambda_i \partial \lambda_j} = \langle A_i A_j \rangle - \langle A_i \rangle \langle A_j \rangle . \quad (1.78)$$

$$\frac{\partial S}{\partial \langle A_i \rangle} = K_B \lambda_i . \quad (1.79)$$

$$D_N = \frac{1}{Z_M} . \quad (1.80)$$

$$D_N = \frac{1}{Z_C} \exp [-\beta H] , \quad (1.81)$$

$$D = \frac{1}{Z_G} \exp [-\beta H + \alpha N] , \quad (1.82)$$

$$D = \frac{1}{Z_G} \exp [-\beta H + \alpha N] , \quad (1.83)$$

$$Z_G = \sum_{N=0}^{\infty} e^{\alpha N} Z_C(\beta, N), \quad (1.84)$$

$$H_N = \sum_i H_1(\mathbf{r}_i, \mathbf{p}_i), \quad (1.85)$$

$$H_1(\mathbf{r}, \mathbf{p}) = p^2/(2m) + V(\mathbf{r}) \quad (1.86)$$

$$Z_C = \int_{-\infty}^{+\infty} \exp(-\beta H_N) d\tau_N, \quad (1.87)$$

$$d\tau_N \equiv \frac{d^3\mathbf{r}_1 d^3\mathbf{p}_1 \dots d^3\mathbf{r}_N d^3\mathbf{p}_N}{\hbar^{3N} N!}. \quad (1.88)$$

$$\begin{aligned} Z_C(\beta, N) &= \int \exp[-\beta H_N] d\tau_N = \int_{-\infty}^{+\infty} \prod_{i=1}^N \left( e^{-\beta p_i^2/(2m)} e^{-\beta V(\mathbf{r}_i)} d^3\mathbf{r}_i d^3\mathbf{p}_i \right) \frac{1}{\hbar^{3N} N!} = \\ &= \frac{1}{\hbar^{3N} N!} \left[ \int_{-\infty}^{+\infty} e^{-\beta p^2/(2m)} e^{-\beta V(\mathbf{r})} d^3\mathbf{r} d^3\mathbf{p} \right]^N. \end{aligned} \quad (1.89)$$

$$\Omega \Xi \equiv \int_{-\infty}^{+\infty} e^{-\beta V(\mathbf{r})} d^3\mathbf{r}. \quad (1.90)$$

$$Z_C(\beta, N) = \frac{1}{\hbar^{3N} N!} (2\pi m K_B T)^{3N/2} \Omega^N \Xi^N \equiv \frac{Z_1^N}{N!}, \quad (1.91)$$

$$Z_1 \equiv \frac{(2\pi m K_B T)^{3/2}}{\hbar^3} \Omega \Xi. \quad (1.92)$$

$$Z_G = \sum_{N=0}^{\infty} \frac{1}{N!} (e^{\alpha} Z_1)^N = \exp[e^{\alpha} Z_1], \quad (1.93)$$

$$f(\mathbf{r}, \mathbf{p}) \equiv \sum_{N=0}^{\infty} \left[ \mathcal{P}(N) \sum_{i=1}^N \mathcal{P}_N(i, \mathbf{r}, \mathbf{p}) \right], \quad (1.94)$$

$$D \sim \begin{bmatrix} \frac{1}{Z_G} & & & \\ & \frac{1}{Z_G} e^{\alpha} \exp(-\beta H_1) & & \\ & & \frac{1}{Z_G} e^{2\alpha} \exp(-\beta [H_1(\mathbf{r}_1, \mathbf{p}_1) + H_1(\mathbf{r}_2, \mathbf{p}_2)]) & \\ & & & \ddots \end{bmatrix} \quad (1.95)$$

$$D_{N,N} = \frac{1}{Z_G} e^{\alpha N} \exp \left[ -\beta \sum_{i=1}^N H_1(\mathbf{r}_i, \mathbf{p}_i) \right]. \quad (1.96)$$

$$\mathcal{P}(N) = \int_{-\infty}^{+\infty} D_{N,N} d\tau_N = \frac{e^{\alpha N} Z_C(\beta, N)}{Z_G(\alpha, \beta)}. \quad (1.97)$$

$$\mathcal{P}_N(1, \mathbf{r}, \mathbf{p}) = \int_{-\infty}^{+\infty} D_N \frac{d^3 \mathbf{r}_2 d^3 \mathbf{p}_2 \dots d^3 \mathbf{r}_N d^3 \mathbf{p}_N}{\hbar^{3N} N!} = \frac{1}{N \hbar^3} \int_{-\infty}^{+\infty} D_N d\tau_{N-1} \quad (1.98)$$

$$\begin{aligned} \mathcal{P}_N(1, \mathbf{r}, \mathbf{p}) &= \frac{1}{N \hbar^3} \int_{-\infty}^{+\infty} \frac{e^{-\beta H_N} d\tau_{N-1}}{\frac{1}{N!} Z_1^N(\beta)} = \frac{1}{N \hbar^3} \int_{-\infty}^{+\infty} \frac{N(N-1)! e^{-\beta H_N} d\tau_{N-1}}{Z_1^N(\beta)} = \\ &= \frac{(N-1)!}{\hbar^3} e^{-\beta H_1(\mathbf{r}, \mathbf{p})} \frac{1}{Z_1^N} \int e^{-\beta H_{N-1}} d\tau_{N-1} = \frac{1}{\hbar^3} \frac{e^{-\beta H_1(\mathbf{r}, \mathbf{p})}}{Z_1(\beta)}. \end{aligned} \quad (1.99)$$

$$\int e^{-\beta H_{N-1}} d\tau_{N-1} = \frac{1}{(N-1)!} Z_1^{(N-1)}, \quad (1.100)$$

$$f(\mathbf{r}, \mathbf{p}) = \sum_{N=0}^{\infty} \frac{e^{\alpha N} Z_C(\beta, N)}{Z_G(\beta, \alpha)} N \frac{1}{\hbar^3} \frac{e^{-\beta H_1}}{Z_1}, \quad (1.101)$$

$$f(\mathbf{r}, \mathbf{p}) = \frac{1}{\hbar^3 Z_G} \sum_{N=0}^{\infty} \frac{e^{\alpha N} Z_1^N}{N!} N \frac{e^{-\beta H_1}}{Z_1} = \frac{e^{\alpha} e^{-\beta H_1}}{\hbar^3} \frac{\sum_{N=1}^{\infty} \frac{e^{\alpha(N-1)} Z_1^{N-1}}{(N-1)!}}{\exp[e^{\alpha} Z_1]}, \quad (1.102)$$

$$f(\mathbf{r}, \mathbf{p}) = \frac{e^{\alpha - \beta H_1(\mathbf{r}, \mathbf{p})}}{\hbar^3}. \quad (1.103)$$

$$\langle N \rangle = \frac{\partial \ln Z_G}{\partial \alpha} = \frac{\partial}{\partial \alpha} e^{\alpha} Z_1 = e^{\alpha} Z_1 \quad (1.104)$$

$$\begin{aligned} f'(\mathbf{p}) &= \int_{-\infty}^{+\infty} f(\mathbf{r}, \mathbf{p}) d^3 \mathbf{r} = \int_{-\infty}^{+\infty} d^3 \mathbf{r} \frac{1}{\hbar^3} e^{\alpha} \exp \{ -\beta [p^2/(2m) + V(\mathbf{r})] \} = \\ &= \frac{\Omega}{\hbar^3} \exp \left[ \alpha - \beta \frac{p^2}{2m} \right] = \langle N \rangle g(\mathbf{p}). \end{aligned} \quad (1.105)$$

$$g(\mathbf{p}) = \frac{\Omega}{\hbar^3} e^{-\frac{\beta p^2}{2m}} \frac{1}{Z_1} = \frac{e^{-\beta p^2/2m}}{(2\pi K_B m T)^{3/2}} = \frac{e^{-\frac{p^2}{2K_B T m}}}{(2\pi K_B m T)^{3/2}}. \quad (1.106)$$

$$H = \sum_{i=1}^{N^e} H_1^e(\mathbf{r}_i \mathbf{p}_i) + \sum_{i=1}^{N^p} H_1^p(\mathbf{r}_i \mathbf{p}_i) + \sum_{i=1}^{N^H} H_1^H(\mathbf{r}_i \mathbf{p}_i) \quad (1.107)$$

$$H_1^a(\mathbf{r}, \mathbf{p}) = \frac{p^2}{2m^a} + V(\mathbf{r}) + h_1^a, \quad a = \{e, p, H\}. \quad (1.108)$$

$$Z_G(\beta, \mu^e, \mu^p, \mu^H) = \prod_a \{ \text{Tr} [\exp(\beta \mu^a N^a)] \text{Tr} [\exp(-\beta H^a)] \} . \quad (1.109)$$

$$\Omega \equiv \int \exp[-\beta V(\mathbf{r})] d^3\mathbf{r} , \quad (1.110)$$

$$\zeta^a(T) \equiv \text{Tr} \{ \exp(-\beta h_1^a) \} = \sum d_i^a \exp[-\beta \mathcal{E}_i^a] , \quad (1.111)$$

$$Z_G(\beta, \mu^e, \mu^p, \mu^H) = \prod_a \sum_{N^a} \frac{1}{N^a!} \left[ e^{\beta \mu^a} Z_1^a \zeta^a(T) \right]^{N^a} = \prod_a \exp[e^{\beta \mu^a} Z_1^a \zeta^a(T)] . \quad (1.112)$$

$$\langle N^a \rangle = \frac{1}{\beta} \frac{\partial \ln Z_G}{\partial \mu^a} = \frac{(2\pi m^a K_B T)^{3/2}}{\hbar^3} \Omega \zeta^a(T) \exp\left(\frac{\mu^a}{K_B T}\right) . \quad (1.113)$$

$$\mu^e + \mu^p = \mu^H . \quad (1.114)$$

$$\prod_a \left[ \exp\left(\frac{\mu^a}{K_B T}\right) \right]^{\nu^a} = 1 , \text{ with } \nu^e = \nu^p = 1, \nu^H = -1 . \quad (1.115)$$

$$\prod_a \left[ n^a \frac{\hbar^3}{(2\pi m^a K_B T)^{3/2}} \frac{1}{\zeta^a(T)} \right]^{\nu^a} = 1 \quad (1.116)$$

$$\frac{n_e}{n_H} = \frac{4}{d_0} \frac{1}{n_e} \frac{(2\pi m^e K_B T)^{3/2}}{\hbar^3} \exp\left(-\frac{U_i}{K_B T}\right) . \quad (1.117)$$

## Chapter 2

# Plasma

$$\nabla \cdot \mathbf{E}(\mathbf{r}) = \frac{\rho_e(\mathbf{r})}{\varepsilon_0}, \quad (2.1)$$

$$\mathbf{E}(\mathbf{r}) = -\nabla\phi(\mathbf{r}), \quad (2.2)$$

$$\Delta\phi(\mathbf{r}) = -\frac{\rho_e(\mathbf{r})}{\varepsilon_0}. \quad (2.3)$$

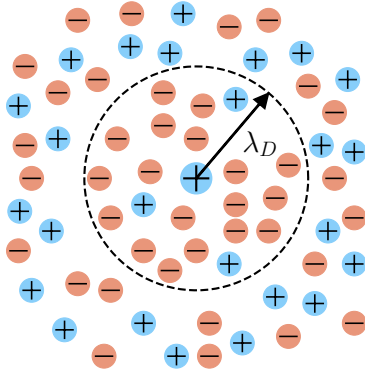
$$f_\alpha(\mathbf{r}, \mathbf{v}) = A_\alpha \exp \left[ -\frac{\frac{1}{2}m_\alpha v^2 + q_\alpha\phi(\mathbf{r})}{K_B T_\alpha} \right], \quad (2.4)$$

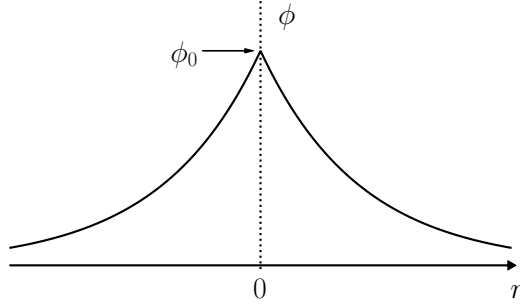
$$n_\alpha(\phi) = \int_0^\infty f_\alpha(\mathbf{v}) d^3\mathbf{v} = \exp \left[ -\frac{q_\alpha\phi(\mathbf{r})}{K_B T_\alpha} \right] \int_0^\infty A_\alpha \exp \left[ -\frac{\frac{1}{2}m_\alpha v^2}{K_B T_\alpha} \right] d^3\mathbf{v}. \quad (2.5)$$

$$n_\alpha(\phi) = n_{\alpha,0} \exp \left[ -\frac{q_\alpha\phi(\mathbf{r})}{K_B T_\alpha} \right]. \quad (2.6)$$

$$\rho_e(\phi) = \sum_\alpha q_\alpha n_\alpha(\phi(\mathbf{r})). \quad (2.7)$$

$$\Delta\phi(\mathbf{r}) = -\frac{1}{\varepsilon_0} \sum_\alpha q_\alpha n_{\alpha,0} \exp \left[ -\frac{q_\alpha\phi(\mathbf{r})}{K_B T_\alpha} \right]. \quad (2.8)$$





$$\Delta\phi(\mathbf{r}) = -\frac{1}{\varepsilon_0} \left[ \sum_{\alpha} q_{\alpha} n_{\alpha,0} - \sum_{\alpha} \frac{q_{\alpha}^2 n_{\alpha,0} \phi(\mathbf{r})}{K_{\text{B}} T_{\alpha}} + \dots \right]. \quad (2.9)$$

$$\Delta\phi(\mathbf{r}) = \left[ \sum_{\alpha} \frac{q_{\alpha}^2 n_{\alpha,0}}{\varepsilon_0 K_{\text{B}} T_{\alpha}} \right] \phi(\mathbf{r}). \quad (2.10)$$

$$\frac{1}{r} \frac{d^2}{dr^2} [r\phi(r)] = a\phi(r), \quad (2.11)$$

$$a = \sum_{\alpha} \frac{q_{\alpha}^2 n_{\alpha,0}}{\varepsilon_0 K_{\text{B}} T_{\alpha}} \quad (2.12)$$

$$\frac{d^2}{dr^2} \psi(r) = a\psi(r), \quad (2.13)$$

$$\psi(r) = c_1 e^{\sqrt{a}r} + c_2 e^{-\sqrt{a}r}. \quad (2.14)$$

$$\phi(r) = \frac{c_2}{r} e^{-\sqrt{a}r}. \quad (2.15)$$

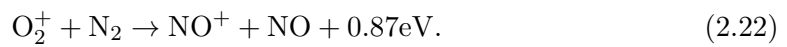
$$\frac{c_2}{r} \rightarrow \frac{Q}{4\pi\varepsilon_0 r}. \quad (2.16)$$

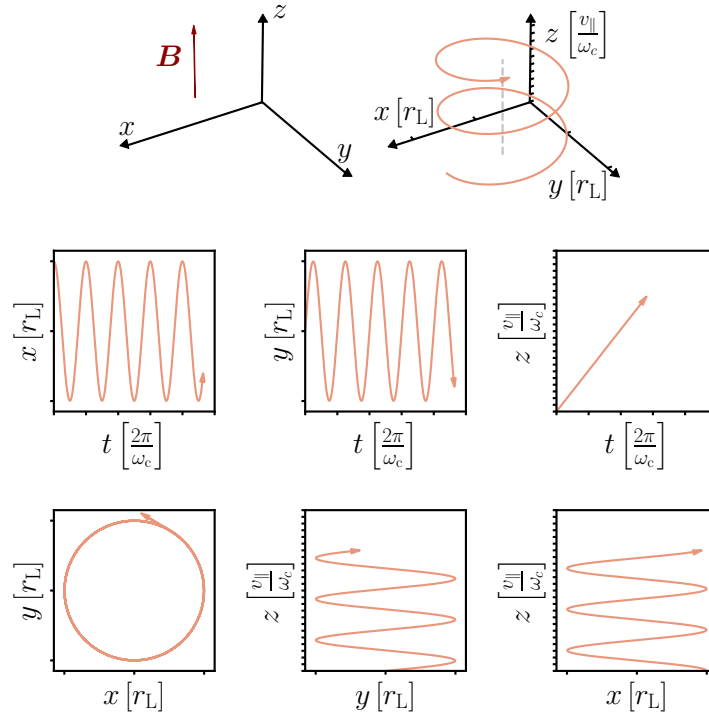
$$\phi(r) = \frac{Q}{4\pi\varepsilon_0 r} \exp\left[-\frac{r}{\lambda_{\text{D}}}\right], \quad (2.17)$$

$$\lambda_{\text{D}} \equiv \left[ \sum_{\alpha} \frac{q_{\alpha}^2 n_{\alpha,0}}{\varepsilon_0 K_{\text{B}} T_{\alpha}} \right]^{-1/2} \quad (2.18)$$

$$\lambda_{\text{D}} \equiv \sqrt{\frac{\varepsilon_0 K_{\text{B}} T_{\text{e}}}{n_0 e^2}}, \quad (2.19)$$

$$E_{\text{k}} = \frac{1}{2} m \bar{v}^2 = K_{\text{B}} T \quad (2.20)$$





## Chapter 3

# Charged particle motion in electromagnetic fields

$$m \frac{d\mathbf{v}(\mathbf{r}, t)}{dt} = q [\mathbf{E}(\mathbf{r}, t) + \mathbf{v}(\mathbf{r}, t) \times \mathbf{B}(\mathbf{r}, t)], \quad (3.1)$$

### 3.1 Homogeneous magnetic field

$$m \frac{d}{dt} \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix} = q \left[ \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix} \times \begin{pmatrix} 0 \\ 0 \\ B \end{pmatrix} \right], \quad (3.2)$$

$$m \frac{dv_x}{dt} = qBv_y, \quad (3.3)$$

$$m \frac{dv_y}{dt} = -qBv_x, \quad (3.4)$$

$$m \frac{dv_z}{dt} = 0. \quad (3.5)$$

$$\frac{d^2v_x}{dt^2} = \frac{q}{m}B \frac{dv_y}{dt} \quad \text{and} \quad \frac{d^2v_y}{dt^2} = -\frac{q}{m}B \frac{dv_x}{dt}. \quad (3.6)$$

$$\frac{d^2v_x}{dt^2} = -\left(\frac{q}{m}B\right)^2 v_x \quad \text{and} \quad \frac{d^2v_y}{dt^2} = -\left(\frac{q}{m}B\right)^2 v_y, \quad (3.7)$$

$$\omega_c \equiv \frac{qB \operatorname{sgn} q}{m} = \frac{|q|B}{m} \quad (3.8)$$

$$v_{x,y} = v_{\perp} \exp(i\omega_c t + i\delta_{x,y}), \quad (3.9)$$

$$v_x = v_{\perp} \exp i\omega_c t, \quad (3.10)$$

$$v_y = \frac{m}{qB} \frac{dv_x}{dt} = \frac{m}{qB} v_{\perp} i\omega_c \exp i\omega_c t = \frac{m}{qB} v_{\perp} i \frac{qB \operatorname{sgn} q}{m} \exp i\omega_c t = v_{\perp} i \operatorname{sgn} q \exp i\omega_c t. \quad (3.11)$$

$$\frac{dx}{dt} = \Re[v_x] = \Re[v_{\perp} \exp i\omega_c t] \rightarrow x - x_0 = \Re\left[-i \frac{v_{\perp}}{\omega_c} \exp i\omega_c t\right] = \frac{v_{\perp}}{\omega_c} \sin \omega_c t \quad (3.12)$$

$$\frac{dy}{dt} = \Re[v_y] = \Re[v_{\perp} i \exp i\omega_c t] \rightarrow y - y_0 = \Re\left[\frac{v_{\perp}}{\omega_c} \operatorname{sgn} q \exp i\omega_c t\right] = \frac{v_{\perp}}{\omega_c} \operatorname{sgn} q \cos \omega_c t \quad (3.13)$$

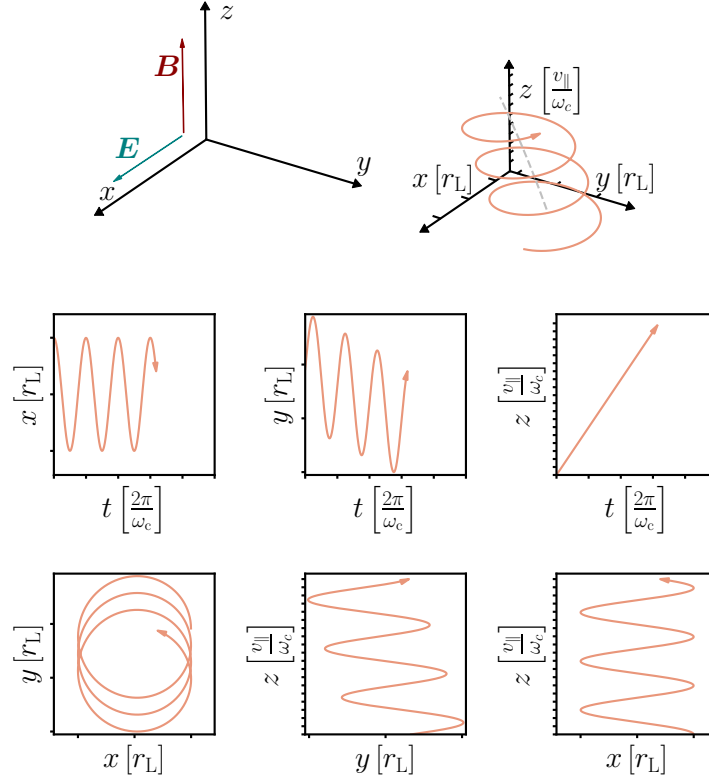
$$\frac{dz}{dt} = \Re[v_z] = 0 \rightarrow z - z_0 = v_{\parallel} t. \quad (3.14)$$

$$r_L = \frac{v_{\perp} m}{qB \operatorname{sgn} q} = \frac{v_{\perp} m}{|q|B} \quad (3.15)$$

### 3.2 Homogeneous electric field

$$m \frac{d}{dt} \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix} = q \left[ \begin{pmatrix} E_x \\ 0 \\ E_z \end{pmatrix} + \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix} \times \begin{pmatrix} 0 \\ 0 \\ B \end{pmatrix} \right]. \quad (3.16)$$

$$\frac{dv_z}{dt} = \frac{q}{m} E_z \rightarrow z - z_0 = \frac{q}{2m} E_z t^2 + v_{z,0} t. \quad (3.17)$$



$$\frac{dv_x}{dt} = \frac{q}{m} E_x + \omega_c \operatorname{sgn} q v_y, \quad (3.18)$$

$$\frac{dv_y}{dt} = -\omega_c \operatorname{sgn} q v_x. \quad (3.19)$$

$$\frac{d^2 v_x}{dt^2} = \omega_c \operatorname{sgn} q \frac{dv_y}{dt} = -\operatorname{sgn}^2 q \omega_c^2 v_x = -\omega_c^2 v_x, \quad (3.20)$$

$$\frac{d^2 v_y}{dt^2} = -\operatorname{sgn} q \omega_c \frac{dv_x}{dt} = -\operatorname{sgn} q \omega_c \left( \frac{q}{m} E_x + \omega_c \operatorname{sgn} q v_y \right) = -\omega_c^2 \left( \frac{E_x}{B} + v_y \right). \quad (3.21)$$

$$v_x = v_{\perp} \exp i \omega_c t. \quad (3.22)$$

$$\frac{d^2}{dt^2} \left( v_y + \frac{E_x}{B} \right) = -\omega_c^2 \left( v_y + \frac{E_x}{B} \right), \quad (3.23)$$

$$v_y = v_{\perp} i \operatorname{sgn} q \exp i \omega_c t - \frac{E_x}{B}. \quad (3.24)$$

$$m \frac{d\mathbf{v}}{dt} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B}) \quad (3.25)$$

$$\begin{aligned}
 m \frac{d\mathbf{v}}{dt} &= m \frac{d\mathbf{v}'}{dt} = q \left( \mathbf{E} + \mathbf{v}' \times \mathbf{B} + \frac{\mathbf{E} \times \mathbf{B}}{B^2} \times \mathbf{B} \right) = \\
 &= q \left[ \mathbf{E} + \mathbf{v}' \times \mathbf{B} - \frac{1}{B^2} \mathbf{E} B^2 + \frac{\mathbf{B}(\mathbf{E} \cdot \mathbf{B})}{B^2} \right] = \\
 &= q \left[ \mathbf{v}' \times \mathbf{B} + \mathbf{E}_{\parallel} \right], \tag{3.26}
 \end{aligned}$$

$$\mathbf{E}_{\parallel} \equiv \frac{\mathbf{B}(\mathbf{E} \cdot \mathbf{B})}{B^2}. \tag{3.27}$$

$$\mathbf{v} = \mathbf{v}' + \frac{\mathbf{E} \times \mathbf{B}}{B^2}. \tag{3.28}$$

$$\mathbf{v}_E = \frac{\mathbf{E} \times \mathbf{B}}{B^2}, \tag{3.29}$$

$$\mathbf{F}_E = q\mathbf{E}. \tag{3.30}$$

$$\mathbf{v}_E = \frac{1}{q} \frac{\mathbf{F}_E \times \mathbf{B}}{B^2}. \tag{3.31}$$

$$\mathbf{v}_{gs} = \frac{1}{q} \frac{\mathbf{F} \times \mathbf{B}}{B^2}. \tag{3.32}$$

$$\mathbf{v}_g = \frac{m}{q} \frac{\mathbf{g} \times \mathbf{B}}{B^2}, \tag{3.33}$$

$$\mathbf{v}_c = \frac{mv_{\parallel}^2}{qR^2B^2} \mathbf{R} \times \mathbf{B} \tag{3.34}$$

### 3.3 Inhomogeneous magnetic field

$$L \sim \frac{|\mathbf{B}|}{|\nabla B|}. \tag{3.35}$$

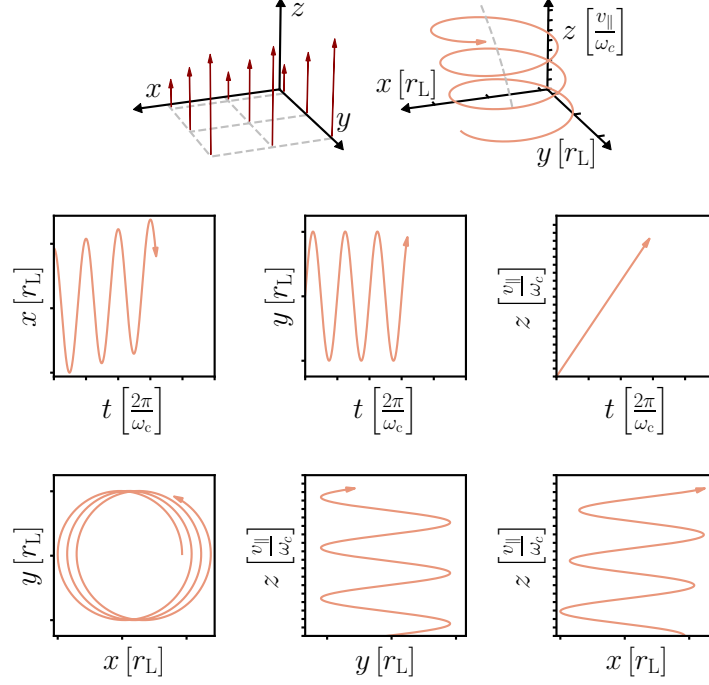
$$\mathbf{B} = \mathbf{B}_0 + (\mathbf{r} \cdot \nabla) \mathbf{B} \tag{3.36}$$

$$B_z = B_0 + y \frac{\partial B_z}{\partial y}. \tag{3.37}$$

$$F_y = -qv_x B_z = -qv_x \left( B_0 + y \frac{\partial B_z}{\partial y} \right). \tag{3.38}$$

$$F_y = -qv_{\perp} \cos \omega_c t \left( B_0 + r_L \operatorname{sgn} q \cos \omega_c t \frac{\partial B_z}{\partial y} \right). \tag{3.39}$$

$$\begin{aligned}
 \langle F_y \rangle &= \frac{1}{2\pi} \int_0^{2\pi} F_y d(\omega_c t) = -\frac{1}{2\pi} \int_0^{2\pi} qv_{\perp} \cos \omega_c t \left( B_0 + r_L \operatorname{sgn} q \cos \omega_c t \frac{\partial B_z}{\partial y} \right) d(\omega_c t) = \\
 &= -qv_{\perp} B_0 \frac{1}{2\pi} \int_0^{2\pi} \cos \omega_c t d(\omega_c t) - qv_{\perp} r_L \operatorname{sgn} q \frac{\partial B_z}{\partial y} \frac{1}{2\pi} \int_0^{2\pi} \cos^2 \omega_c t d(\omega_c t) = \\
 &= 0 - \frac{1}{2} v_{\perp} r_L q \operatorname{sgn} q \frac{\partial B_z}{\partial y}. \tag{3.40}
 \end{aligned}$$



$$\langle \mathbf{F} \rangle = -\frac{1}{2} v_{\perp} r_L q \operatorname{sgn} q \nabla |\mathbf{B}|, \quad (3.41)$$

$$\mathbf{v}_{\nabla B} = \frac{1}{2} \operatorname{sgn} q v_{\perp} r_L \frac{\mathbf{B} \times \nabla |\mathbf{B}|}{B^2}. \quad (3.42)$$

$$\Delta = \tau |\mathbf{v}_{\nabla B}| = \frac{2\pi}{\omega_c} \frac{1}{2} v_{\perp} r_L \frac{|\nabla B|}{|B|} = \pi r_L \frac{r_L}{L} \ll r_L \quad (3.43)$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j} = 0, \quad (3.44)$$

$$\nabla \times \mathbf{B} = \frac{1}{R} \frac{\partial}{\partial R} (R B_{\theta}) = 0 \quad (3.45)$$

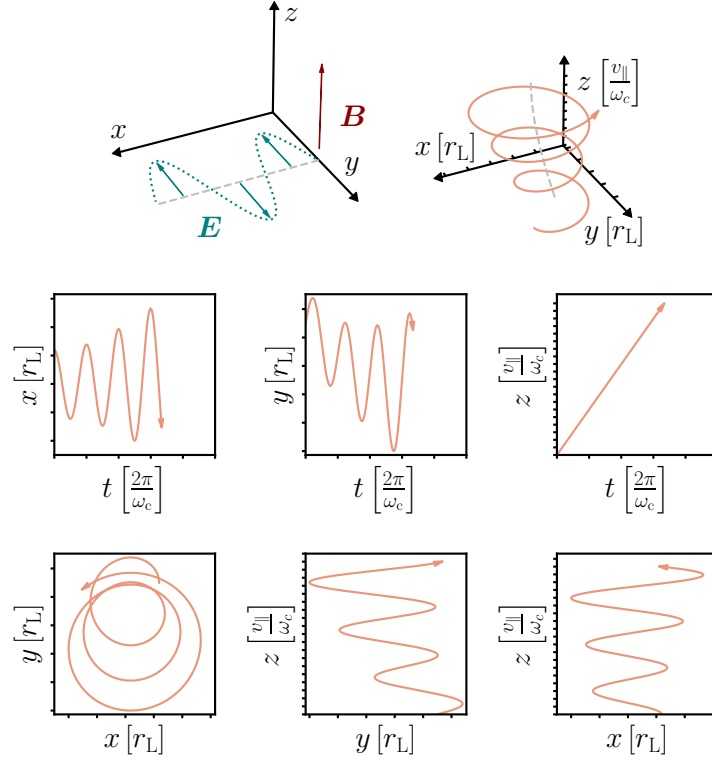
$$B_{\theta} = \frac{\mathcal{C}}{R}, \quad (3.46)$$

$$\nabla B = \frac{\partial}{\partial R} \frac{\mathcal{C}}{R} = -\mathcal{C} \frac{1}{R^2} \frac{\mathbf{R}}{R} \quad (3.47)$$

$$\nabla B = -\frac{|B|}{R^2} \mathbf{R}. \quad (3.48)$$

$$\mathbf{v}_{\nabla B} = -\frac{1}{2} \operatorname{sgn} q \frac{v_{\perp} r_L}{B^2} \mathbf{B} \times |\mathbf{B}| \frac{\mathbf{R}}{R^2} = \frac{1}{2} \frac{m v_{\perp}^2}{q B^2 R^2} \mathbf{R} \times \mathbf{B}. \quad (3.49)$$

$$\mathbf{v}_{gc} = \mathbf{v}_c + \mathbf{v}_{\nabla B} = \frac{m}{q} \frac{\mathbf{R} \times \mathbf{B}}{R^2 B^2} \left( v_{||}^2 + \frac{1}{2} v_{\perp}^2 \right). \quad (3.50)$$



### 3.4 Inhomogeneous electric field

$$\mathbf{E} = E_0 \cos ky \mathbf{e}_x, \quad (3.51)$$

$$\frac{dv_x}{dt} = \frac{qB}{m} v_y + \frac{q}{m} E_x(y) \quad \text{and} \quad \frac{dv_y}{dt} = -\frac{qB}{m} v_x. \quad (3.52)$$

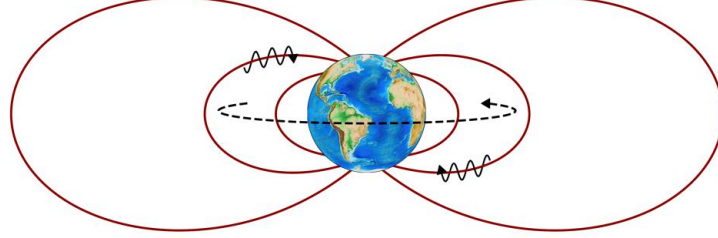
$$\frac{d^2 v_x}{dt^2} = -\omega_c^2 v_x + \omega_c \frac{\text{sgn } q}{B} \frac{dE_x}{dt} \quad \text{and} \quad \frac{d^2 v_y}{dt^2} = -\omega_c^2 v_y - \omega_c^2 \frac{E_x(y)}{B}. \quad (3.53)$$

$$\frac{d^2 v_y}{dt^2} = -\omega_c^2 v_y - \omega_c^2 \frac{E_0}{B} \cos[k(y_0 + r_L \text{sgn } q \cos \omega_c t)], \quad (3.54)$$

$$\left\langle \frac{d^2 v_y}{dt^2} \right\rangle = -\omega_c^2 \langle v_y \rangle - \omega_c^2 \frac{E_0}{B} \langle \cos[k(y_0 + r_L \text{sgn } q \cos \omega_c t)] \rangle. \quad (3.55)$$

$$\left\langle \frac{d^2 v_y}{dt^2} \right\rangle = 0. \quad (3.56)$$

$$\begin{aligned} \cos[k(y_0 + r_L \text{sgn } q \cos \omega_c t)] &= \\ &= \cos ky_0 \cos(kr_L \text{sgn } q \cos \omega_c t) - \sin ky_0 \sin(kr_L \text{sgn } q \cos \omega_c t) = \\ &= \cos ky_0 \cos(kr_L \cos \omega_c t) - \text{sgn } q \sin ky_0 \sin(kr_L \cos \omega_c t) \\ &= \cos ky_0 \left[ 1 - \frac{1}{2} k^2 r_L^2 \cos^2 \omega_c t \right] - \text{sgn } q \sin ky_0 (kr_L \cos \omega_c t) \end{aligned} \quad (3.57)$$



$$\langle v_y \rangle = -\frac{E_0}{B} \cos ky_0 \left( 1 - \frac{1}{4} k^2 r_L^2 \right) = -\frac{E_x(y)}{B} \left( 1 - \frac{1}{4} k^2 r_L^2 \right). \quad (3.58)$$

$$\mathbf{v}_E = \frac{\mathbf{E} \times \mathbf{B}}{B^2} \left( 1 - \frac{1}{4} k^2 r_L^2 \right) \quad (3.59)$$

$$\mathbf{v}_E = -\frac{\mathbf{B} \times}{B^2} \left( 1 + \frac{1}{4} r_L^2 \Delta \right) \mathbf{E}. \quad (3.60)$$

### 3.5 Drift currents

$$\mathbf{j}_{\text{drift}} = \sum_k n_k q_k \mathbf{v}_{\text{drift},k} = \frac{\sum_k n_k \mathbf{F}_k}{B^2} \times \mathbf{B}. \quad (3.61)$$

$$\mathbf{j}_{\nabla B} = \frac{1}{2} \sum_k \frac{m_k n_k v_{\perp,k}^2}{B^2} (\mathbf{B} \times \nabla B). \quad (3.62)$$

### 3.6 Guiding centre motion

$$\frac{r_L}{L} = \frac{v_{\perp}}{\omega_c L} \ll 1 \quad \text{and} \quad \frac{1}{\omega_c \tau} \ll 1, \quad (3.63)$$

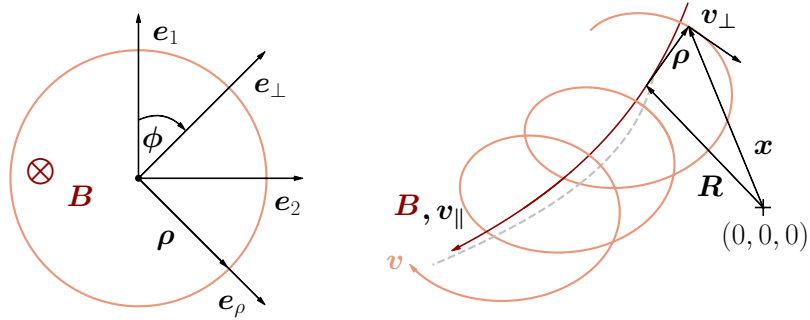
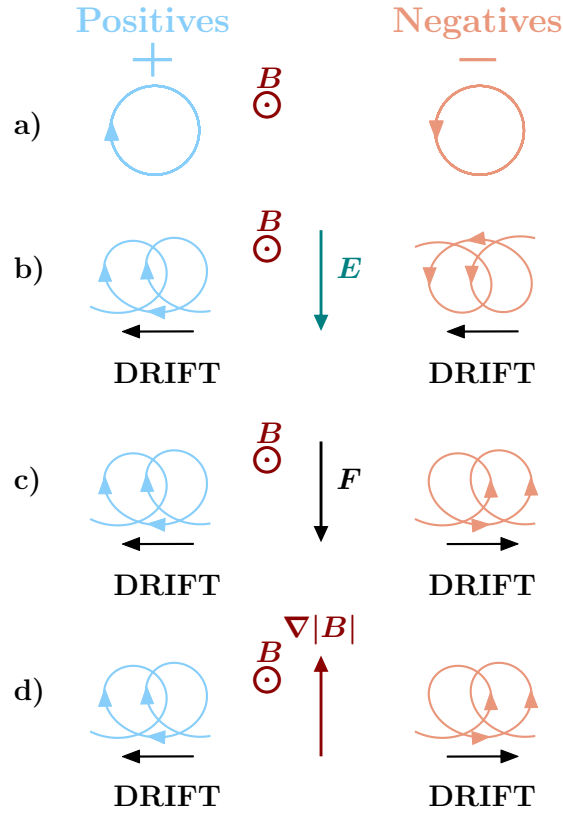
$$\frac{d\mathbf{x}}{dt} = \mathbf{v}, \quad (3.64)$$

$$\frac{d\mathbf{v}}{dt} = \frac{q}{m} (\mathbf{E} + \mathbf{v} \times \mathbf{B}). \quad (3.65)$$

$$\mathbf{R}(t) = \mathbf{x}(t) - \boldsymbol{\rho}(t), \quad (3.66)$$

$$\mathbf{u} = \mathbf{v} - \mathbf{v}_E. \quad (3.67)$$

$$\boldsymbol{\rho} \equiv \frac{\epsilon}{B} \mathbf{b} \times \mathbf{u}, \quad (3.68)$$



$$\omega_c(t) = \frac{B(\mathbf{x}, t)}{\epsilon} \quad (3.69)$$

$$\mathbf{u} = v_\parallel \mathbf{b} + u_\perp \mathbf{e}_\perp, \quad \text{where} \quad \mathbf{e}_\perp \equiv \mathbf{e}_1 \cos \phi + \mathbf{e}_2 \sin \phi, \quad \text{and} \quad \phi = \phi_0 - \omega_c t. \quad (3.70)$$

$$\mathbf{e}_\rho \equiv \mathbf{e}_2 \cos \phi - \mathbf{e}_1 \sin \phi \quad (3.71)$$

$$\mathbf{x} = \mathbf{R} + \boldsymbol{\rho}, \quad (3.72)$$

$$\boldsymbol{\rho} = -\frac{m}{qB^2} \mathbf{u} \times \mathbf{B} = \rho \operatorname{sgn} q \mathbf{e}_\rho, \quad (3.73)$$

$$\mathbf{u}_\perp = u_\perp \mathbf{e}_\perp \quad \text{and} \quad (3.74)$$

$$\mathbf{v} = v_\parallel \mathbf{b} + u_\perp \mathbf{e}_\perp + \mathbf{v}_E. \quad (3.75)$$

$$\begin{aligned} \frac{d\mathbf{R}}{dt} &= \frac{d\mathbf{x}}{dt} - \frac{d\boldsymbol{\rho}}{dt} = \mathbf{v} - \frac{d}{dt} \left( \frac{\epsilon}{B} \mathbf{b} \times \mathbf{u} \right) = \\ &= v_\parallel \mathbf{b} + u_\perp \mathbf{e}_\perp + \mathbf{v}_E + \epsilon \mathbf{u} \times \frac{d}{dt} \left( \frac{\mathbf{b}}{B} \right) - \frac{\epsilon}{B} \mathbf{b} \times \frac{d\mathbf{u}}{dt} = \\ &= v_\parallel \mathbf{b} + u_\perp \mathbf{e}_\perp + \mathbf{v}_E + \epsilon \mathbf{u} \times \frac{d}{dt} \left( \frac{\mathbf{b}}{B} \right) - \frac{\epsilon}{B} \mathbf{b} \times \left( \frac{d\mathbf{v}}{dt} - \frac{d\mathbf{v}_E}{dt} \right) = \\ &= v_\parallel \mathbf{b} + \mathbf{v}_E + \epsilon \mathbf{u} \times \frac{d}{dt} \left( \frac{\mathbf{b}}{B} \right) + \frac{\epsilon}{B} \mathbf{b} \times \frac{d\mathbf{v}_E}{dt}, \end{aligned} \quad (3.76)$$

$$\begin{aligned} -\frac{\epsilon}{B} \mathbf{b} \times \frac{d\mathbf{v}}{dt} &= -\frac{\epsilon}{B} \mathbf{b} \times \left[ \frac{1}{\epsilon} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \right] = \\ &= -\frac{1}{B} \mathbf{b} \times \mathbf{E} - \frac{1}{B} (\mathbf{v} \mathbf{b} \cdot \mathbf{B} - \mathbf{B} \mathbf{b} \cdot \mathbf{v}) = -\frac{1}{B} \mathbf{b} \times \mathbf{E} - \mathbf{v} \frac{\mathbf{B} \cdot \mathbf{B}}{B^2} + v_\parallel \frac{\mathbf{B}}{B} = \\ &= \mathbf{v}_E - \mathbf{v} + v_\parallel \mathbf{b} = -u_\perp \mathbf{e}_\perp. \end{aligned} \quad (3.77)$$

$$\begin{aligned} \frac{d(v_\parallel \mathbf{b})}{dt} + \frac{d(u_\perp \mathbf{e}_\perp)}{dt} + \frac{d\mathbf{v}_E}{dt} &= \frac{1}{\epsilon} [\mathbf{E} + (v_\parallel \mathbf{b} + u_\perp \mathbf{e}_\perp + \mathbf{v}_E) \times \mathbf{B}] = \\ &= \frac{1}{\epsilon} [\mathbf{E} + v_\parallel \mathbf{b} \times \mathbf{B} + u_\perp \mathbf{e}_\perp \times \mathbf{B} + \mathbf{v}_E \times \mathbf{B}] = \\ &= \frac{1}{\epsilon} \left[ \mathbf{E} + u_\perp \mathbf{e}_\perp \times \mathbf{B} - \cancel{\frac{E}{B^2} \mathbf{B} \cdot \mathbf{B}} + \frac{\mathbf{E} \cdot \mathbf{B}}{B^2} \mathbf{B} \right] = \\ &= \frac{1}{\epsilon} [E_\parallel \mathbf{b} + u_\perp \mathbf{e}_\perp \times \mathbf{B}], \end{aligned} \quad (3.78)$$

$$\frac{1}{\epsilon} [E_\parallel \mathbf{b} \cdot \mathbf{e}_\perp + u_\perp (\mathbf{e}_\perp \times \mathbf{B}) \cdot \mathbf{e}_\perp] = \left[ \frac{d(v_\parallel \mathbf{b})}{dt} + \frac{d(u_\perp \mathbf{e}_\perp)}{dt} + \frac{d\mathbf{v}_E}{dt} \right] \cdot \mathbf{e}_\perp. \quad (3.79)$$

$$0 = \mathbf{e}_\perp \cdot v_\parallel \frac{d\mathbf{b}}{dt} + \mathbf{e}_\perp \cdot \mathbf{b} \frac{dv_\parallel}{dt} + \mathbf{e}_\perp \cdot \mathbf{e}_\perp \frac{du_\perp}{dt} + u_\perp \mathbf{e}_\perp \cdot \frac{d\mathbf{e}_\perp}{dt} + \mathbf{e}_\perp \cdot \frac{d\mathbf{v}_E}{dt}. \quad (3.80)$$

$$\mathbf{e}_\perp \cdot \frac{d\mathbf{e}_\perp}{dt} = \frac{1}{2} \frac{d}{dt} (\mathbf{e}_\perp \cdot \mathbf{e}_\perp) = 0, \quad (3.81)$$

$$\frac{du_\perp}{dt} = -\mathbf{e}_\perp \cdot \left( v_\parallel \frac{d\mathbf{b}}{dt} + \frac{d\mathbf{v}_E}{dt} \right). \quad (3.82)$$

$$\frac{1}{\epsilon} [E_\parallel \mathbf{b} \cdot \mathbf{b} + u_\perp (\mathbf{e}_\perp \times \mathbf{B}) \cdot \mathbf{b}] = \left[ \frac{d(v_\parallel \mathbf{b})}{dt} + \frac{d(u_\perp \mathbf{e}_\perp)}{dt} + \frac{d\mathbf{v}_E}{dt} \right] \cdot \mathbf{b}. \quad (3.83)$$

$$\frac{E_{\parallel}}{\epsilon} = \mathbf{b} \cdot \mathbf{b} \frac{dv_{\parallel}}{dt} + v_{\parallel} \frac{d\mathbf{b}}{dt} \cdot \mathbf{b} + \mathbf{b} \cdot \mathbf{e}_{\perp} \frac{du_{\perp}}{dt} + u_{\perp} \mathbf{b} \cdot \frac{d\mathbf{e}_{\perp}}{dt} + \frac{d\mathbf{v}_E \cdot \mathbf{b}}{dt} - \mathbf{v}_E \cdot \frac{d\mathbf{b}}{dt}. \quad (3.84)$$

$$u_{\perp} \mathbf{b} \cdot \frac{d\mathbf{e}_{\perp}}{dt} = u_{\perp} \frac{d\mathbf{e}_{\perp} \cdot \mathbf{b}}{dt} - u_{\perp} \mathbf{e}_{\perp} \cdot \frac{d\mathbf{b}}{dt} = -u_{\perp} \mathbf{e}_{\perp} \cdot \frac{d\mathbf{b}}{dt}. \quad (3.85)$$

$$\frac{dv_{\parallel}}{dt} = \frac{E_{\parallel}}{\epsilon} + (u_{\perp} \mathbf{e}_{\perp} + \mathbf{v}_E) \cdot \frac{d\mathbf{b}}{dt}. \quad (3.86)$$

$$\frac{1}{\epsilon} \left[ E_{\parallel} \overset{\textcircled{1}}{\mathbf{b} \cdot \mathbf{e}_{\rho}} + u_{\perp} (\mathbf{e}_{\perp} \times \mathbf{B}) \cdot \mathbf{e}_{\rho} \right] = \overset{\textcircled{3}}{\mathbf{e}_{\rho} \cdot \frac{d(v_{\parallel} \mathbf{b})}{dt}} + \overset{\textcircled{4}}{\mathbf{e}_{\rho} \cdot \frac{d(u_{\perp} \mathbf{e}_{\perp})}{dt}} + \overset{\textcircled{5}}{\mathbf{e}_{\rho} \cdot \frac{d\mathbf{v}_E}{dt}}. \quad (3.87)$$

$$\mathbf{b} \perp \mathbf{e}_{\rho} \Rightarrow \mathbf{b} \cdot \mathbf{e}_{\rho} = 0 \quad (3.88)$$

$$\begin{aligned} u_{\perp} (\mathbf{e}_{\perp} \times \mathbf{B}) \cdot \mathbf{e}_{\rho} &= u_{\perp} (\mathbf{e}_{\rho} \times \mathbf{e}_{\perp}) \cdot \mathbf{B} = u_{\perp} \left[ \begin{pmatrix} -\sin \phi \\ \cos \phi \\ 0 \end{pmatrix} \times \begin{pmatrix} \cos \phi \\ \sin \phi \\ 0 \end{pmatrix} \right] \cdot \mathbf{B} = \\ &= u_{\perp} \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 0 \\ B \end{pmatrix} = -u_{\perp} B \end{aligned} \quad (3.89)$$

$$\mathbf{e}_{\rho} \cdot \frac{d(v_{\parallel} \mathbf{b})}{dt} = \mathbf{e}_{\rho} \cdot v_{\parallel} \frac{d\mathbf{b}}{dt} + \cancel{\mathbf{e}_{\rho} \cdot \mathbf{b} \frac{dv_{\parallel}}{dt}} \quad (3.90)$$

$$\begin{aligned} \mathbf{e}_{\rho} \cdot \frac{d(u_{\perp} \mathbf{e}_{\perp})}{dt} &= u_{\perp} (\mathbf{e}_2 \cos \phi - \mathbf{e}_1 \sin \phi) \cdot \frac{d(\mathbf{e}_1 \cos \phi + \mathbf{e}_2 \sin \phi)}{dt} + \cancel{\mathbf{e}_{\rho} \cdot \mathbf{e}_{\perp} \frac{du_{\perp}}{dt}} = \\ &= u_{\perp} \left[ (\mathbf{e}_2 \cos \phi - \mathbf{e}_1 \sin \phi) \cdot \left( \cos \phi \frac{d\mathbf{e}_1}{dt} - \mathbf{e}_1 \sin \phi \frac{d\phi}{dt} + \sin \phi \frac{d\mathbf{e}_2}{dt} + \mathbf{e}_2 \cos \phi \frac{d\phi}{dt} \right) \right] = \\ &= u_{\perp} \left[ \cos^2 \phi \mathbf{e}_2 \cdot \frac{d\mathbf{e}_1}{dt} - \cancel{\sin \phi \cos \phi \mathbf{e}_1 \cdot \frac{d\mathbf{e}_1}{dt}} + \cancel{\sin \phi \cos \phi \frac{d\phi}{dt} \mathbf{e}_1 \cdot \mathbf{e}_2} + \sin^2 \phi \frac{d\phi}{dt} \mathbf{e}_1 \cdot \mathbf{e}_1 + \right. \\ &\quad \left. + \cancel{\sin \phi \cos \phi \mathbf{e}_2 \cdot \frac{d\mathbf{e}_2}{dt}} - \sin^2 \phi \mathbf{e}_1 \cdot \frac{d\mathbf{e}_2}{dt} + \cos^2 \phi \frac{d\phi}{dt} \mathbf{e}_2 \cdot \mathbf{e}_2 - \cancel{\sin \phi \cos \phi \frac{d\phi}{dt} \mathbf{e}_1 \cdot \mathbf{e}_2} \right] = \\ &= u_{\perp} \left[ (\sin^2 \phi + \cos^2 \phi) \left( \frac{d\phi}{dt} + \mathbf{e}_2 \cdot \frac{d\mathbf{e}_1}{dt} \right) - \sin^2 \phi \left( \mathbf{e}_1 \cdot \frac{d\mathbf{e}_2}{dt} + \mathbf{e}_2 \cdot \frac{d\mathbf{e}_1}{dt} \right) \right] = \\ &= u_{\perp} \left[ \frac{d\phi}{dt} + \mathbf{e}_2 \cdot \frac{d\mathbf{e}_1}{dt} - \sin^2 \phi \left( \frac{d(\mathbf{e}_1 \cdot \mathbf{e}_2)}{dt} - \cancel{\mathbf{e}_2 \cdot \frac{d\mathbf{e}_1}{dt}} + \cancel{\mathbf{e}_2 \cdot \frac{d\mathbf{e}_1}{dt}} \right) \right], \end{aligned} \quad (3.91)$$

$$\frac{d\phi}{dt} = -\frac{B}{\epsilon} - \mathbf{e}_2 \cdot \frac{d\mathbf{e}_1}{dt} - \frac{1}{u_{\perp}} \mathbf{e}_{\rho} \cdot \left( v_{\parallel} \frac{d\mathbf{b}}{dt} + \frac{d\mathbf{v}_E}{dt} \right). \quad (3.92)$$

$$\frac{d\mathbf{R}}{dt} = v_{\parallel} \mathbf{b} + \mathbf{v}_E + \epsilon \mathbf{u} \times \frac{d}{dt} \left( \frac{\mathbf{b}}{B} \right) + \frac{\epsilon}{B} \mathbf{b} \times \frac{d\mathbf{v}_E}{dt}, \quad (3.93)$$

$$\frac{du_{\perp}}{dt} = -\mathbf{e}_{\perp} \cdot \left( v_{\parallel} \frac{d\mathbf{b}}{dt} + \frac{d\mathbf{v}_E}{dt} \right), \quad (3.94)$$

$$\frac{dv_{\parallel}}{dt} = \frac{E_{\parallel}}{\epsilon} + (u_{\perp} \mathbf{e}_{\perp} + \mathbf{v}_E) \cdot \frac{d\mathbf{b}}{dt}, \quad (3.95)$$

$$\frac{d\phi}{dt} = -\frac{B}{\epsilon} - \mathbf{e}_2 \cdot \frac{d\mathbf{e}_1}{dt} - \frac{1}{u_{\perp}} \mathbf{e}_{\rho} \cdot \left( v_{\parallel} \frac{d\mathbf{b}}{dt} + \frac{d\mathbf{v}_E}{dt} \right). \quad (3.96)$$

$$\frac{d}{dt} = \frac{\partial}{\partial t} + (v_{\parallel} \mathbf{b} + \mathbf{v}_E + u_{\perp} \mathbf{e}_{\perp}) \cdot \nabla \quad (3.97)$$

$$\begin{aligned} \frac{d\mathbf{R}}{dt} = & v_{\parallel} \mathbf{b} + \mathbf{v}_E + \frac{\epsilon}{B} (v_{\parallel} \mathbf{b} + u_{\perp} \mathbf{e}_{\perp}) \times \left[ \frac{\partial}{\partial t} + (v_{\parallel} \mathbf{b} + \mathbf{v}_E + u_{\perp} \mathbf{e}_{\perp}) \cdot \nabla \right] \mathbf{b} + \\ & \epsilon (v_{\parallel} \mathbf{b} + u_{\perp} \mathbf{e}_{\perp}) \times \mathbf{b} \left[ \frac{\partial}{\partial t} + (v_{\parallel} \mathbf{b} + \mathbf{v}_E + u_{\perp} \mathbf{e}_{\perp}) \cdot \nabla \right] \left( \frac{1}{B} \right) + \\ & \frac{\epsilon}{B} \mathbf{b} \times \left[ \frac{\partial}{\partial t} + (v_{\parallel} \mathbf{b} + \mathbf{v}_E) \cdot \nabla \right] \mathbf{v}_E + \frac{\epsilon}{B} \mathbf{b} \times (u_{\perp} \mathbf{e}_{\perp} \cdot \nabla) \mathbf{v}_E. \end{aligned} \quad (3.98)$$

$$\frac{\epsilon}{B} \mathbf{b} \times \mathbf{b} \left[ \frac{\partial}{\partial t} + (v_{\parallel} \mathbf{b} + \mathbf{v}_E) \cdot \nabla \right] v_{\parallel} = 0 \quad (3.99)$$

$$\begin{aligned} \frac{d\mathbf{R}}{dt} = & v_{\parallel} \mathbf{b} + \mathbf{v}_E + \frac{\epsilon}{B} \mathbf{b} \times D_t (v_{\parallel} \mathbf{b} + \mathbf{v}_E) + \frac{\epsilon}{B} v_{\parallel} \mathbf{b} \times [u_{\perp} \mathbf{e}_{\perp} \cdot \nabla] \mathbf{b} + \\ & \frac{\epsilon u_{\perp}}{B} \mathbf{e}_{\perp} \times \left[ \frac{\partial}{\partial t} + (v_{\parallel} \mathbf{b} + \mathbf{v}_E + u_{\perp} \mathbf{e}_{\perp}) \cdot \nabla \right] \mathbf{b} - \\ & \frac{\epsilon u_{\perp}}{B^2} \mathbf{e}_{\perp} \times \mathbf{b} \left[ \frac{\partial}{\partial t} + (v_{\parallel} \mathbf{b} + \mathbf{v}_E + u_{\perp} \mathbf{e}_{\perp}) \cdot \nabla \right] B + \frac{\epsilon}{B} \mathbf{b} \times (u_{\perp} \mathbf{e}_{\perp} \cdot \nabla) \mathbf{v}_E, \end{aligned} \quad (3.100)$$

$$D_t \equiv \frac{\partial}{\partial t} + (v_{\parallel} \mathbf{b} + \mathbf{v}_E) \cdot \nabla. \quad (3.101)$$

$$\frac{dz}{dt} = f_z(z) = \langle f_z \rangle + \tilde{f}_z, \quad (3.102)$$

$$\bar{z} = z + \frac{\epsilon}{B} \int_0^{\phi} \tilde{f}_z(\phi') d\phi'. \quad (3.103)$$

$$\frac{d\phi}{dt} = -\frac{B}{\epsilon} - \mathbf{e}_2 \cdot \frac{d\mathbf{e}_1}{dt} - \frac{1}{u_{\perp}} \mathbf{e}_{\rho} \cdot \left( v_{\parallel} \frac{d\mathbf{b}}{dt} + \frac{d\mathbf{v}_E}{dt} \right) = -\frac{B}{\epsilon} + \mathcal{O}(\epsilon^0). \quad (3.104)$$

$$\frac{d}{dt} = \frac{\partial}{\partial t} + \frac{d\mathbf{R}}{dt} \cdot \frac{\partial}{\partial \mathbf{R}} + \frac{dv_{\parallel}}{dt} \frac{\partial}{\partial v_{\parallel}} + \frac{du_{\perp}}{dt} \frac{\partial}{\partial u_{\perp}} + \frac{d\phi}{dt} \frac{\partial}{\partial \phi}. \quad (3.105)$$

$$\frac{d}{dt} = -\frac{B}{\epsilon} \frac{\partial}{\partial \phi} + \mathcal{O}(\epsilon^0). \quad (3.106)$$

$$\begin{aligned}\frac{d\bar{z}}{dt} &= \frac{dz}{dt} - \frac{B}{\epsilon} \frac{\partial}{\partial \phi} \frac{\epsilon}{B} \int_{\phi} \tilde{f}_z(\phi') d\phi' = \\ &= \frac{dz}{dt} - \tilde{f}_z + \mathcal{O}(\epsilon) = \langle f_z(z) \rangle + \tilde{f}_z - \tilde{f}_z + \mathcal{O}(\epsilon) = \langle f_z(z) \rangle + \mathcal{O}(\epsilon).\end{aligned}\quad (3.107)$$

$$\mathbf{b}(\mathbf{x}, t) = \mathbf{b}(\mathbf{R}, t) + \boldsymbol{\rho} \cdot \nabla \mathbf{b}(\mathbf{R}, t) + \dots = \mathbf{b}(\mathbf{R}, t) + \epsilon \frac{u_{\perp}}{B} \mathbf{e}_{\rho} \cdot \nabla \mathbf{b}(\mathbf{R}, t) + \dots \quad (3.108)$$

$$\langle \mathbf{b}(\mathbf{x}, t) \rangle = \mathbf{b}(\mathbf{R}, t) + \mathcal{O}(\epsilon). \quad (3.109)$$

$$\frac{d\bar{z}}{dt} = \langle f_z(\bar{z}) \rangle + \mathcal{O}(\epsilon), \quad \text{Q.E.D.} \quad (3.110)$$

$$\left\langle \frac{d\mathbf{R}}{dt} \right\rangle = \left\langle v_{\parallel} \mathbf{b} + \mathbf{v}_E + \frac{\epsilon}{B} \mathbf{b} \times D_t (v_{\parallel} \mathbf{b} + \mathbf{v}_E) + \frac{\epsilon u_{\perp}^2}{B} \mathbf{e}_{\perp} \times (\mathbf{e}_{\perp} \cdot \nabla) \mathbf{b} - \frac{\epsilon u_{\perp}^2}{B^2} \mathbf{e}_{\perp} \times \mathbf{b} (\mathbf{e}_{\perp} \cdot \nabla) B \right\rangle. \quad (3.111)$$

$$(\mathbf{e}_{\perp} \times \mathbf{b}) (\mathbf{e}_{\perp} \cdot \nabla) B = (\mathbf{e}_1 \sin \phi - \mathbf{e}_2 \cos \phi) (\cos \phi \nabla_1 + \sin \phi \nabla_2) B \quad (3.112)$$

$$\langle (\mathbf{e}_{\perp} \times \mathbf{b}) (\mathbf{e}_{\perp} \cdot \nabla) B \rangle = \frac{1}{2} (\mathbf{e}_1 \nabla_2 B - \mathbf{e}_2 \nabla_1 B) = -\frac{1}{2} (\mathbf{b} \times \nabla B). \quad (3.113)$$

$$\langle \mathbf{e}_{\perp} \times (\mathbf{e}_{\perp} \cdot \nabla) \mathbf{b} \rangle = \frac{1}{2} (\mathbf{e}_1 \nabla_2 - \mathbf{e}_2 \nabla_1) b_3 + \frac{1}{2} \mathbf{e}_3 (\nabla_1 b_2 - \nabla_2 b_1). \quad (3.114)$$

$$\mathbf{b}(\mathbf{R} + \delta \mathbf{R}) = \mathbf{b}(\mathbf{R}) + \delta \mathbf{R} \cdot \nabla \mathbf{b}(\mathbf{R}) + \mathcal{O}(\delta R^2). \quad (3.115)$$

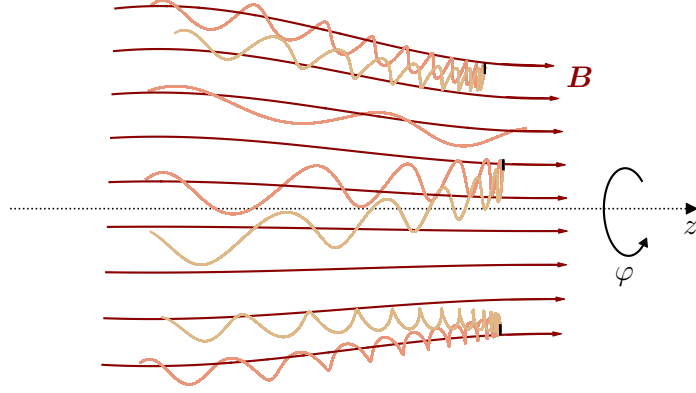
$$|\mathbf{b}'|^2 = 1 = \left| \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} + \begin{pmatrix} \delta b_1 \\ \delta b_2 \\ \delta b_3 \end{pmatrix} \right|^2 = \delta b_1^2 + \delta b_2^2 + (1 + \delta b_3)^2 \sim \delta b_1^2 + \delta b_2^2 + \delta b_3^2 + 2\delta b_3 + 1. \quad (3.116)$$

$$\langle \mathbf{e}_{\perp} \times (\mathbf{e}_{\perp} \cdot \nabla) \mathbf{b} \rangle = \frac{1}{2} \mathbf{e}_3 (\mathbf{e}_3 \cdot \nabla \times \mathbf{b}) = \frac{1}{2} \mathbf{b} (\mathbf{b} \cdot \nabla \times \mathbf{b}), \quad (3.117)$$

$$\left\langle \frac{d\mathbf{R}}{dt} \right\rangle = \overset{(1)}{v_{\parallel} \mathbf{b}} + \overset{(2)}{\mathbf{v}_E} + \frac{\epsilon}{B} \mathbf{b} \times D_t (\overset{(3)}{v_{\parallel} \mathbf{b} + \mathbf{v}_E}) + \frac{\epsilon u_{\perp}^2}{2B} (\overset{(4)}{\mathbf{b} \cdot \nabla \times \mathbf{b}}) \mathbf{b} + \frac{\epsilon u_{\perp}^2}{2B^2} \overset{(5)}{\mathbf{b} \times \nabla B} + \mathcal{O}(\epsilon^2). \quad (3.118)$$

$$\begin{aligned}\mathbf{v}_c &= \frac{mv_{\parallel}^2}{qB} [\mathbf{b} \times (\mathbf{b} \cdot \nabla) \mathbf{b}] = \frac{mv_{\parallel}^2}{qB^2} \mathbf{B} \times [\mathbf{b} \cdot (\mathbf{e}_1 \nabla b_1 + \mathbf{e}_2 \nabla b_2 + \mathbf{b} \nabla b_3)] = \\ &= \frac{mv_{\parallel}^2}{qB^2} \mathbf{B} \times \nabla b_3 = -\frac{mv_{\parallel}^2}{qB^2 R^2} \mathbf{B} \times \mathbf{R},\end{aligned}\quad (3.119)$$

$$\begin{aligned}\frac{\epsilon}{B} \mathbf{b} \times D_t (v_{\parallel} \mathbf{b} + \mathbf{v}_E) &= \frac{\epsilon}{B} \mathbf{b} \times \frac{d\mathbf{v}_E}{dt} + \mathcal{A} = \frac{m}{qB^2} \mathbf{B} \times \frac{d}{dt} \frac{\mathbf{E} \times \mathbf{B}}{B^2} + \mathcal{A} = \frac{m}{qB^4} \mathbf{B} \times \left( \frac{\partial \mathbf{E}}{\partial t} \times \mathbf{B} \right) + \mathcal{B} = \\ &= \frac{m}{qB^4} \left( \mathbf{B} \cdot \mathbf{B} \frac{\partial \mathbf{E}}{\partial t} - \mathbf{B} \cdot \frac{\partial \mathbf{E}}{\partial t} \mathbf{B} \right) + \mathcal{B} = \frac{m}{qB^2} \frac{\partial \mathbf{E}}{\partial t} + \mathcal{C} = \frac{\text{sgn } q}{\omega_c B} \frac{\partial \mathbf{E}}{\partial t} + \mathcal{C},\end{aligned}\quad (3.120)$$



$$\left\langle \frac{dv_{\parallel}}{dt} \right\rangle = \frac{E_{\parallel}}{\epsilon} - \frac{u_{\perp}^2}{2B} \nabla_{\parallel} B + \mathbf{v}_E \cdot D_t \mathbf{b} + \mathcal{O}(\epsilon) \quad (3.121)$$

$$\left\langle \frac{du_{\perp}}{dt} \right\rangle = \frac{v_{\parallel} u_{\perp}}{2B} \nabla_{\parallel} B - \frac{u_{\perp}}{2} (\nabla \cdot \mathbf{v}_E - \mathbf{b} \cdot \nabla_{\parallel} \mathbf{v}_E) + \mathcal{O}(\epsilon) \quad (3.122)$$

$$\left\langle \frac{d\phi}{dt} \right\rangle = -\frac{B}{\epsilon} - \mathbf{e}_2 \cdot D_t \mathbf{e}_1 - \frac{v_{\parallel}}{2} \mathbf{b} \cdot \nabla \times (v_{\parallel} \mathbf{b} + \mathbf{v}_E) + \mathcal{O}(\epsilon) \quad (3.123)$$

### 3.7 Magnetic mirrors

$$\nabla \cdot \mathbf{B} = \frac{1}{r} \frac{\partial}{\partial r} r B_r + \frac{1}{r} \frac{\partial B_{\varphi}}{\partial \varphi} + \frac{\partial B_z}{\partial z} = 0 \quad (3.124)$$

$$r B_r = - \int_0^r r' \frac{\partial B_z}{\partial z} dr'. \quad (3.125)$$

$$r B_r \sim - \left[ \frac{\partial B_z}{\partial z} \right]_{r=0} \int_0^r r' dr' = - \frac{1}{2} r^2 \left[ \frac{\partial B_z}{\partial z} \right]_{r=0}, \quad (3.126)$$

$$B_r \sim - \frac{1}{2} r \left[ \frac{\partial B_z}{\partial z} \right]_{r=0}. \quad (3.127)$$

$$\begin{aligned} F &= q \mathbf{v} \times \mathbf{B} = q \begin{pmatrix} v_r \\ v_{\varphi} \\ v_z \end{pmatrix} \times \begin{pmatrix} B_r \\ B_{\varphi} \\ B_z \end{pmatrix} = q \begin{pmatrix} v_{\varphi} B_z - v_z B_{\varphi} \\ v_z B_r - v_r B_z \\ v_r B_{\varphi} - v_{\varphi} B_r \end{pmatrix} = \\ &= q \begin{pmatrix} v_{\varphi} B_z^{(1)} \\ v_z B_r^{(2)} - v_r B_z^{(3)} \\ -v_{\varphi} B_r^{(4)} \end{pmatrix}. \end{aligned} \quad (3.128)$$

$$F_z = -qv_\varphi B_r = \frac{1}{2}qv_\varphi r \frac{\partial B_z}{\partial z}, \quad (3.129)$$

$$F_z = -\frac{1}{2}qv_\perp \operatorname{sgn} q \frac{v_\perp m}{q \operatorname{sgn} q B} \frac{\partial B_z}{\partial z} = -\frac{1}{2} \frac{mv_\perp^2}{B} \frac{\partial B_z}{\partial z} = -\mu \frac{\partial B_z}{\partial z}, \quad (3.130)$$

$$F_\parallel = -\mu \frac{\partial B}{\partial s} = -\mu \nabla_\parallel B, \quad (3.131)$$

$$m \frac{dv_\parallel}{dt} = -\mu \frac{\partial B}{\partial s} \quad (3.132)$$

$$m \mathbf{v}_\parallel \cdot \frac{d\mathbf{v}_\parallel}{dt} = \frac{d}{dt} \left( \frac{1}{2} m v_\parallel^2 \right) = -\mu \frac{\partial B}{\partial s} \cdot \mathbf{v}_\parallel = -\mu \frac{\partial B}{\partial s} \cdot \frac{d\mathbf{s}}{dt} = -\mu \frac{dB}{dt}. \quad (3.133)$$

$$\frac{d}{dt} \left( \frac{1}{2} m v_\parallel^2 + \frac{1}{2} m v_\perp^2 \right) = \frac{d}{dt} \left( \frac{1}{2} m v_\parallel^2 + \mu B \right) = 0, \quad (3.134)$$

$$0 = -\mu \frac{dB}{dt} + \frac{d}{dt} \mu B = -\mu \frac{dB}{dt} + \mu \frac{dB}{dt} + B \frac{d\mu}{dt} \quad (3.135)$$

$$\frac{d\mu}{dt} = 0. \quad (3.136)$$

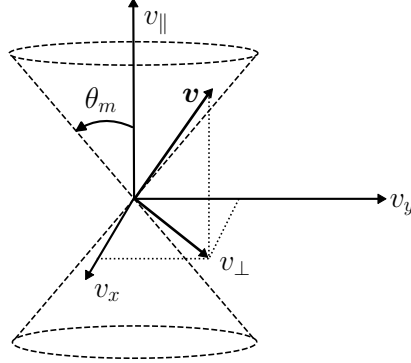
$$\frac{du_\perp}{dt} = \frac{v_\parallel u_\perp}{2B} \mathbf{b} \cdot \nabla_\parallel B - \frac{u_\perp}{2} (\nabla \cdot \overset{\textcircled{A}}{\mathbf{v}_E} - \mathbf{b} \cdot \overset{\textcircled{B}}{\nabla_\parallel \mathbf{v}_E}) + \mathcal{O}(\epsilon), \quad (3.137)$$

$$\begin{aligned} \nabla \cdot \mathbf{v}_E &= \nabla \cdot \frac{\mathbf{E} \times \mathbf{B}}{B^2} = \frac{1}{B^2} [\mathbf{B} \cdot (\nabla \times \mathbf{E}) - \mathbf{E} \cdot (\nabla \times \mathbf{B})] + (\mathbf{E} \times \mathbf{B}) \cdot \nabla \frac{1}{B^2} = \\ &= -\frac{1}{B^2} \mathbf{B} \cdot \frac{\partial \mathbf{B}}{\partial t} - \frac{\mathbf{E}_\parallel + \mathbf{E}_\perp}{B^2} \cdot (\nabla \times \mathbf{B}) - \frac{2(\mathbf{E} \times \mathbf{B})}{B^2} \cdot \frac{\nabla B}{B} = \\ &= -\frac{B}{B^2} \mathbf{b} \cdot \left[ B \frac{\partial \mathbf{b}}{\partial t} + \mathbf{b} \frac{\partial B}{\partial t} \right] - \frac{\mathbf{E}_\parallel}{B^2} \cdot [(\nabla B) \times \mathbf{b} + B(\nabla \times \mathbf{b})] - \\ &- \frac{\mathbf{E}_\perp}{B^2} \cdot (\nabla \times \mathbf{B}) - 2\mathbf{v}_E \cdot \frac{\nabla B}{B} = -\mathbf{b} \cdot \frac{\partial \mathbf{b}}{\partial t} - \frac{1}{B} \frac{\partial B}{\partial t} - \\ &- \frac{E_\parallel}{B^2} \mathbf{b} \cdot [(\nabla B) \times \mathbf{b}] - \frac{E_\parallel}{B} \mathbf{b} \cdot (\nabla \times \mathbf{b}), \end{aligned} \quad (3.138)$$

$$\begin{aligned} \mathbf{b} \cdot \nabla_\parallel \mathbf{v}_E &= \nabla_\parallel (\mathbf{b} \cdot \mathbf{v}_E) - \mathbf{v}_E \cdot (\nabla_\parallel \mathbf{b}) = -\mathbf{v}_E \cdot [(\mathbf{b} \cdot \nabla) \mathbf{b}] = \\ &- \frac{\mathbf{v}_E}{B} \cdot [(\mathbf{b} \cdot \nabla) \mathbf{B}] - (\mathbf{v}_E \cdot \mathbf{B}) (\mathbf{b} \cdot \nabla) \frac{1}{B} = -\frac{\mathbf{v}_E}{B} \cdot [(\mathbf{b} \cdot \nabla) \mathbf{B}], \end{aligned} \quad (3.139)$$

$$\nabla B^2 = \nabla (\mathbf{B} \cdot \mathbf{B}) = 2\mathbf{B} \cdot \nabla \mathbf{B} = 2B \mathbf{b} \cdot \nabla \mathbf{B} \quad \text{and} \quad \nabla B^2 = 2B \nabla B. \quad (3.140)$$

$$\begin{aligned} \frac{du_\perp}{dt} &= \frac{v_\parallel u_\perp}{2B} \mathbf{b} \cdot \nabla_\parallel B + \frac{u_\perp}{2} \frac{1}{B} \frac{\partial B}{\partial t} + \frac{u_\perp}{2} \mathbf{v}_E \cdot \frac{\nabla B}{B} + \mathcal{O}(\epsilon) = \\ &= \frac{u_\perp}{2} \left[ \frac{\partial}{\partial t} + (v_\parallel \mathbf{b} + \mathbf{v}_E) \cdot \nabla \right] B + \mathcal{O}(\epsilon) = \frac{u_\perp}{2B} D_t B + \mathcal{O}(\epsilon), \end{aligned} \quad (3.141)$$



$$\frac{d\mu}{dt} = \frac{u_{\perp} \frac{d}{dt} u_{\perp}}{B} + \frac{u_{\perp}^2}{2B} \frac{d}{dt} \frac{1}{B} = \frac{u_{\perp}^2}{2B^2} D_t B + \mathcal{O}(\epsilon) - \frac{u_{\perp}^2}{2B^2} \frac{dB}{dt} = \mathcal{O}(\epsilon). \quad (3.142)$$

$$\mu = \text{const} = \frac{1}{2} m v_{\perp,0}^2 / B_0 = \frac{1}{2} m v_{\perp}'^2 / B' \quad \text{and} \quad v_{\perp}'^2 = v_{\perp,0}^2 + v_{\parallel,0}^2 \equiv v_0^2, \quad (3.143)$$

$$\frac{B_0}{B'} = \frac{v_{\perp,0}^2}{v_{\perp}'^2} = \frac{v_{\perp,0}^2}{v_0^2} = \sin^2 \vartheta, \quad (3.144)$$

$$\sin^2 \vartheta_{\max} = \frac{B_0}{B_{\max}}. \quad (3.145)$$

### 3.8 Adiabatic invariants

$$\frac{d\mathbf{q}}{dt} = \frac{\partial \mathcal{H}}{\partial \mathbf{p}} \quad \text{and} \quad \frac{d\mathbf{p}}{dt} = -\frac{\partial \mathcal{H}}{\partial \mathbf{q}}. \quad (3.146)$$

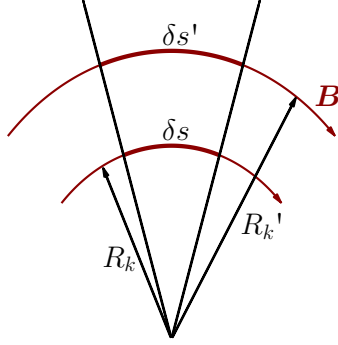
$$q_i(t) = q_i(t + T), \quad (3.147)$$

$$\oint p_i dq_i = \text{const} \quad (3.148)$$

$$\begin{aligned} \int_0^T p dq &= \int_0^T (-ml^2 \varphi_0 \omega \sin \omega t) (-\varphi_0 \omega \sin \omega t) d(\omega t) = \\ &= ml^2 \varphi_0^2 \omega \int_0^T \sin^2 \omega t d(\omega t) = \pi ml^2 \varphi_0^2 \omega = \text{const}. \end{aligned} \quad (3.149)$$

$$J_1 = \oint p_i dq_i = \int_0^{2\pi} m v_{\perp} r_L d\phi = 2\pi v_{\perp} \frac{m^2 v_{\perp}}{q \operatorname{sgn} q B} = \frac{4\pi m}{q \operatorname{sgn} q} \mu = \text{const}. \quad (3.150)$$

$$J_2 = \int_a^b v_{\parallel} ds \quad (3.151)$$



$$\frac{\delta s}{R_k} = \frac{\delta s'}{R_k'}. \quad (3.152)$$

$$\frac{\delta s' - \delta s}{\Delta t \delta s} = \frac{R_k' - R_k}{\Delta t R_k} = \mathbf{v}_{gc} \cdot \frac{\mathbf{R}_k}{R_k^2} = \frac{1}{\delta s} \frac{\Delta \delta s}{\Delta t} = \frac{1}{\delta s} \frac{d\delta s}{dt}, \quad (3.153)$$

$$\left| \frac{R_k' - R_k}{\Delta t} \right| = \mathbf{v}_{gc} \cdot \frac{\mathbf{R}_k}{R_k}. \quad (3.154)$$

$$\mathbf{v}_{gc} = \frac{1}{2} \operatorname{sgn} q v_{\perp} r_L \frac{\mathbf{B} \times \nabla B}{B^2} + \frac{m v_{\parallel}^2}{q} \frac{\mathbf{R}_k \times \mathbf{B}}{R_k^2 B^2}. \quad (3.155)$$

$$\frac{1}{\delta s} \frac{d\delta s}{dt} = \frac{1}{2} \frac{m}{q} \frac{v_{\perp}^2}{B^3} (\mathbf{B} \times \nabla B) \cdot \frac{\mathbf{R}_k}{R_k^2}, \quad (3.156)$$

$$W = \frac{1}{2} m v_{\parallel}^2 + \frac{1}{2} m v_{\perp}^2 = \frac{1}{2} m v_{\parallel}^2 + \mu B = W_{\parallel} + W_{\perp}. \quad (3.157)$$

$$\frac{dv_{\parallel}}{dt} = \frac{1}{2} \frac{1}{\sqrt{\frac{2}{m}(W - \mu B)}} \left( -\mu \frac{dB}{dt} \right) \frac{2}{m} \quad (3.158)$$

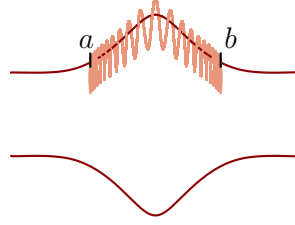
$$\frac{\frac{dv_{\parallel}}{dt}}{v_{\parallel}} = -\frac{1}{2} \frac{\mu \frac{dB}{dt}}{W - \mu B} = -\frac{\mu \frac{dB}{dt}}{m v_{\parallel}^2}, \quad (3.159)$$

$$\frac{dB}{dt} = \frac{dB}{dR} \frac{dR}{dt} = \mathbf{v}_{gc} \cdot \nabla B = \frac{m v_{\parallel}^2}{q} \frac{\mathbf{R}_k \times \mathbf{B}}{R_k^2 B^2} \cdot \nabla B. \quad (3.160)$$

$$\frac{\frac{dv_{\parallel}}{dt}}{v_{\parallel}} = -\frac{\mu}{q} \frac{(\mathbf{R}_k \times \mathbf{B}) \cdot \nabla B}{R_k^2 B^2} = -\frac{1}{2} \frac{m v_{\perp}^2}{B q} \frac{(\mathbf{B} \times \nabla B) \cdot \mathbf{R}_k}{R_k^2 B^2}. \quad (3.161)$$

$$\frac{1}{v_{\parallel} \delta s} \frac{d(v_{\parallel} \delta s)}{dt} = \frac{1}{\delta s} \frac{d\delta s}{dt} + \frac{1}{v_{\parallel}} \frac{dv_{\parallel}}{dt}. \quad (3.162)$$

$$\frac{1}{v_{\parallel} \delta s} \frac{d(v_{\parallel} \delta s)}{dt} = 0 \quad (3.163)$$



$$J_2 = \int_a^{a'} v_{\parallel} \delta s + \int_{a'}^{b'} v_{\parallel} \delta s + \int_{b'}^b v_{\parallel} \delta s, \quad (3.164)$$

$$J_2 = \langle v_{\parallel} \rangle L, \quad (3.165)$$

$$J_3 = \Phi = \int \mathbf{B} \cdot d\mathbf{S} \quad (3.166)$$

### 3.9 van Allen radiation belts

$$\mathbf{B} = (B_R(R, z), 0, B_z(R, z)), \quad (3.167)$$

$$\nabla \cdot \mathbf{B} = 0 = \frac{1}{R} \frac{\partial}{\partial R} R B_R + \frac{\partial B_z}{\partial z}. \quad (3.168)$$

$$B_z = \frac{1}{R} \frac{\partial F}{\partial R} \quad \text{and} \quad B_R = -\frac{1}{R} \frac{\partial F}{\partial z}. \quad (3.169)$$

$$F = \mathcal{M} \frac{R^2}{(R^2 + z^2)^{3/2}}, \quad (3.170)$$

$$\left( m \frac{d\mathbf{v}}{dt} \right)_{\varphi} = q (v_z B_R - v_R B_z) = -\frac{q}{R} \left( v_z \frac{\partial F}{\partial z} + v_R \frac{\partial F}{\partial R} \right) = -\frac{q}{R} \frac{dF}{dt}, \quad (3.171)$$

$$\mathbf{e}_{\varphi} \propto \mathbf{e}_z \times \mathbf{r}. \quad (3.172)$$

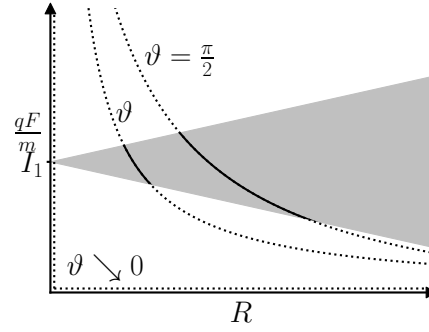
$$|z||r| \sin \vartheta = R, \quad (3.173)$$

$$R \mathbf{e}_{\varphi} = \mathbf{e}_z \times \mathbf{r}. \quad (3.174)$$

$$R v_{\varphi} = R \mathbf{e}_{\varphi} \cdot \mathbf{v} = (\mathbf{e}_z \times \mathbf{r}) \cdot \mathbf{v}. \quad (3.175)$$

$$\frac{d}{dt} R v_{\varphi} = (\mathbf{e}_z \times \mathbf{r}) \cdot \frac{d\mathbf{v}}{dt} + (\mathbf{e}_z \times \frac{d\mathbf{r}}{dt}) \cdot \mathbf{v}. \quad (3.176)$$

$$\frac{d}{dt} R v_{\varphi} = R \mathbf{e}_{\varphi} \cdot \frac{d\mathbf{v}}{dt} = R \left( \frac{d\mathbf{v}}{dt} \right)_{\varphi}. \quad (3.177)$$



$$\frac{d}{dt} \left( Rv_\varphi + \frac{qF}{m} \right) = R \frac{dv_\varphi}{dt} + \frac{q}{m} \frac{dF}{dt} = 0, \quad (3.178)$$

$$Rv_\varphi + \frac{qF}{m} = I_1 \quad (3.179)$$

$$v_R^2 + v_\varphi^2 + v_z^2 = v^2 = I_2^2 \quad (3.180)$$

$$R|v_\varphi| \leq |v|R \quad \rightarrow \quad \left| I_1 - \frac{qF}{m} \right| \leq RI_2. \quad (3.181)$$

## Chapter 4

# Plasma as fluid

### 4.1 Double-fluid model in physics of plasmas

$$\frac{\partial n_\alpha}{\partial t} + \nabla \cdot (n_\alpha \mathbf{u}_\alpha) = 0 \quad (4.1)$$

$$m_\alpha n_\alpha \left[ \frac{\partial \mathbf{u}_\alpha}{\partial t} + (\mathbf{u}_\alpha \cdot \nabla) \mathbf{u}_\alpha \right] = q_\alpha n_\alpha (\mathbf{E} + \mathbf{u}_\alpha \times \mathbf{B}) - \nabla p_\alpha \quad (4.2)$$

$$p_\alpha = C_\alpha (m_\alpha n_\alpha)^{\gamma_\alpha}, \quad (4.3)$$

$$\varepsilon_0 \nabla \cdot \mathbf{E} = n_i q_i + n_e q_e \quad (4.4)$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad (4.5)$$

$$\nabla \cdot \mathbf{B} = 0 \quad (4.6)$$

$$\frac{1}{\mu_0} \nabla \times \mathbf{B} = n_i q_i \mathbf{u}_i + n_e q_e \mathbf{u}_e + \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t}. \quad (4.7)$$

$$\nabla \cdot (\nabla \times \mathbf{E}) = -\nabla \cdot \frac{\partial \mathbf{B}}{\partial t} = -\frac{\partial}{\partial t} \nabla \cdot \mathbf{B}. \quad (4.8)$$

$$\nabla \cdot \mathbf{B} = \text{const}, \quad (4.9)$$

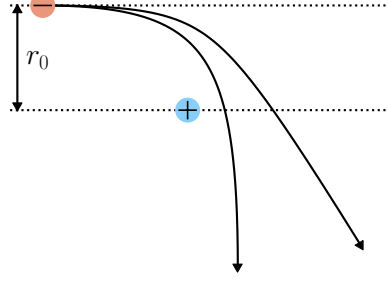
$$p = C \rho^\gamma \quad \text{or} \quad \frac{\nabla p}{p} = \gamma \frac{\nabla \rho}{\rho}. \quad (4.10)$$

### 4.2 Highly ionised plasmas

$$m_e n_e \frac{d\mathbf{u}_e}{dt} = -en_e (\mathbf{E} + \mathbf{u}_e \times \mathbf{B}) - \nabla p_e + \mathbf{P}_{ei}. \quad (4.11)$$

$$\mathbf{P}_{ei} = -\mathbf{P}_{ie}. \quad (4.12)$$

$$\mathbf{P}_{ei} = \underbrace{\eta}_{(5)} \underbrace{e^2}_{(1)} \underbrace{n_e}_{(2)} \underbrace{n_i (\mathbf{u}_i - \mathbf{u}_e)}_{(3)} \sim \underbrace{\eta e^2 n^2 (\mathbf{u}_i - \mathbf{u}_e)}_{(4)}, \quad (4.13)$$



$$\mathbf{P}_{\text{ei}} = m_e n (\mathbf{u}_i - \mathbf{u}_e) \nu_{\text{ei}}, \quad (4.14)$$

$$\nu_{\text{ei}} = \frac{ne^2}{m_e} \eta = \omega_p^2 \varepsilon_0 \eta, \quad (4.15)$$

$$F = -\frac{e^2}{4\pi\varepsilon_0 r^2}, \quad (4.16)$$

$$\Delta t = \frac{r_0}{v}, \quad (4.17)$$

$$\Delta(m_e v) \sim |F \Delta t| = \frac{e^2}{4\pi\varepsilon_0 r_0 v}. \quad (4.18)$$

$$\Delta(m_e v) = m_e v = \frac{e^2}{4\pi\varepsilon_0 r_0 v}, \quad (4.19)$$

$$r_{90} = \frac{e^2}{4\pi\varepsilon_0 m_e v^2}. \quad (4.20)$$

$$\sigma = \pi r_{90}^2 = \frac{e^4}{16\pi\varepsilon_0^2 m_e^2 v^4}. \quad (4.21)$$

$$\nu_{\text{ei}} = \sigma n v = \frac{ne^4}{16\pi\varepsilon_0^2 m_e^2 v^3}. \quad (4.22)$$

$$\eta = \frac{\nu_{\text{ei}} m_e}{ne^2} = \frac{e^2}{16\pi\varepsilon_0^2 m_e v^3} = \frac{e^2 m_e^{1/2}}{16\pi\varepsilon_0^2 (K_B T_e)^{3/2}}. \quad (4.23)$$

$$\int_{r_{90}}^{\lambda_D} \frac{1}{r} dr = [\ln r]_{r_{90}}^{\lambda_D} = \ln \frac{\lambda_D}{r_{90}} = \ln \Lambda, \quad (4.24)$$

$$\eta = \frac{e^2 m_e^{1/2}}{16\pi\varepsilon_0^2 (K_B T_e)^{3/2}} \ln \Lambda. \quad (4.25)$$

$$\mathbf{P}_{\text{ei}} = \eta e^2 n^2 (\mathbf{u}_i - \mathbf{u}_e) = en_e \mathbf{E}. \quad (4.26)$$

$$\mathbf{E} = \eta en (\mathbf{u}_i - \mathbf{u}_e) = \eta \mathbf{j}, \quad (4.27)$$

### 4.3 Single-fluid model

$$\rho \equiv n_i M + n_e m \sim n(M + m) \sim nM, \quad (4.28)$$

$$\mathbf{u} \equiv \frac{1}{\rho}(n_i M \mathbf{u}_i + n_e m \mathbf{u}_e) \sim \frac{n(M \mathbf{u}_i + m \mathbf{u}_e)}{n(M + m)} = \frac{M \mathbf{u}_i + m \mathbf{u}_e}{M + m}, \quad (4.29)$$

$$\mathbf{j} \equiv e(n_i \mathbf{u}_i - n_e \mathbf{u}_e) \sim ne(\mathbf{u}_i - \mathbf{u}_e), \quad (4.30)$$

$$p = p_i + p_e. \quad (4.31)$$

$$Mn \frac{\partial \mathbf{u}_i}{\partial t} = en(\mathbf{E} + \mathbf{u}_i \times \mathbf{B}) - \nabla p_i + Mn\mathbf{g} + \mathbf{P}_{ie}, \quad (4.32)$$

$$mn \frac{\partial \mathbf{u}_e}{\partial t} = -en(\mathbf{E} + \mathbf{u}_e \times \mathbf{B}) - \nabla p_e + mn\mathbf{g} + \mathbf{P}_{ei} \quad (4.33)$$

$$n(M + m) \frac{\partial}{\partial t} \left( \frac{M \mathbf{u}_i + m \mathbf{u}_e}{M + m} \right) = n \frac{\partial}{\partial t} (M \mathbf{u}_i + m \mathbf{u}_e) = en(\mathbf{u}_i - \mathbf{u}_e) \times \mathbf{B} - \nabla p + n(M + m)\mathbf{g}. \quad (4.34)$$

$$\rho \frac{\partial \mathbf{u}}{\partial t} = \mathbf{j} \times \mathbf{B} - \nabla p + \rho \mathbf{g}, \quad (4.35)$$

$$mMn \frac{\partial \mathbf{u}_i}{\partial t} = emn(\mathbf{E} + \mathbf{u}_i \times \mathbf{B}) - m \nabla p_i + mMn\mathbf{g} + m\mathbf{P}_{ie}, \quad (4.36)$$

$$mMn \frac{\partial \mathbf{u}_e}{\partial t} = -eMn(\mathbf{E} + \mathbf{u}_e \times \mathbf{B}) - M \nabla p_e + mMn\mathbf{g} + M\mathbf{P}_{ei}. \quad (4.37)$$

$$nmM \frac{\partial}{\partial t} (\mathbf{u}_i - \mathbf{u}_e) = en(M + m)\mathbf{E} + en(m\mathbf{u}_i + M\mathbf{u}_e) \times \mathbf{B} - m \nabla p_i + M \nabla p_e - (M + m)\mathbf{P}_{ei}. \quad (4.38)$$

$$m\mathbf{u}_i + M\mathbf{u}_e = M\mathbf{u}_i + m\mathbf{u}_e + M(\mathbf{u}_e - \mathbf{u}_i) + m(\mathbf{u}_i - \mathbf{u}_e) = \frac{\rho}{n}\mathbf{u} - (M - m)\frac{\mathbf{j}}{ne}, \quad (4.39)$$

$$P_{ei} = \eta e^2 n^2 (\mathbf{u}_i - \mathbf{u}_e) = \eta en\mathbf{j}. \quad (4.40)$$

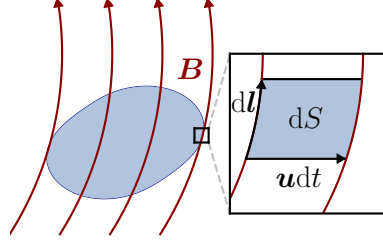
$$\mathbf{E} + \mathbf{u} \times \mathbf{B} - \eta \mathbf{j} = \frac{1}{e\rho} \left[ \frac{Mmn}{e} \frac{\partial \mathbf{j}}{\partial t n} + (M - m)\mathbf{j} \times \mathbf{B} + m \nabla p_i - M \nabla p_e \right]. \quad (4.41)$$

$$\frac{Mmn}{e} \frac{\partial \mathbf{j}}{\partial t n} = MnB \frac{m}{eB} \frac{\partial \mathbf{j}}{\partial t n} = MnB\omega_c^{-1} \frac{\partial \mathbf{j}}{\partial t n} \ll (M - m)\mathbf{j} \times \mathbf{B}, \quad (4.42)$$

$$\frac{1}{\omega_c} \frac{d}{dt} \left( \frac{\mathbf{j}}{n} \right) \ll \frac{\mathbf{j}}{n} \quad (4.43)$$

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} - \eta \mathbf{j} = \frac{1}{en} (\mathbf{j} \times \mathbf{B} - \nabla p_e), \quad (4.44)$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{u} = 0 \quad (4.45)$$



$$\frac{\partial \rho_e}{\partial t} + \nabla \cdot \mathbf{j} = 0. \quad (4.46)$$

$$\rho \frac{\partial \mathbf{u}}{\partial t} = \mathbf{j} \times \mathbf{B} - \nabla p + \rho \mathbf{g}, \quad (4.47)$$

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} = \eta \mathbf{j}, \quad (4.48)$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{u} = 0, \quad (4.49)$$

$$\frac{\partial \rho_e}{\partial t} + \nabla \cdot \mathbf{j} = 0. \quad (4.50)$$

#### 4.4 Approximations to plasma fluid

$$f_{\text{Lorentz}} = \rho_e \mathbf{E} + \mathbf{j} \times \mathbf{B} = 0. \quad (4.51)$$

$$\mathbf{j} \times \mathbf{B} = 0. \quad (4.52)$$

$$\nabla \times \mathbf{B} = \mu \mathbf{j}. \quad (4.53)$$

$$(\nabla \times \mathbf{B}) \times \mathbf{B} = 0, \quad (4.54)$$

$$\nabla \times \mathbf{B} = \alpha \mathbf{B}. \quad (4.55)$$

$$\mathbf{E}' = \mathbf{E} + \mathbf{u} \times \mathbf{B} = \mathbf{0}. \quad (4.56)$$

$$\mathbf{E}' = \eta \mathbf{j}, \quad (4.57)$$

$$\nabla \times (\nabla \times \mathbf{B}) = \nabla \nabla \cdot \mathbf{B} - \Delta \mathbf{B} = \nabla \times \left( \mu_0 \mathbf{j} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right). \quad (4.58)$$

$$\sigma \mu_0 \nabla \times (\mathbf{E} + \mathbf{u} \times \mathbf{B}) = -\epsilon_0 \mu_0 \nabla \times \frac{\partial \mathbf{E}}{\partial t} - \Delta \mathbf{B}. \quad (4.59)$$

$$-\frac{\partial \mathbf{B}}{\partial t} + \nabla \times (\mathbf{u} \times \mathbf{B}) = 0 \quad \rightarrow \quad \nabla \times (\mathbf{u} \times \mathbf{B}) = \frac{\partial \mathbf{B}}{\partial t}. \quad (4.60)$$

$$\frac{d\Psi}{dt} = \frac{d}{dt} \int \mathbf{B} \cdot d\mathbf{S} = \int \mathbf{B} \cdot \frac{d\mathbf{S}}{dt} + \int \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S}. \quad (4.61)$$

$$\oint_l \mathbf{B} \cdot (\mathbf{u} \times d\mathbf{l}) + \int \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} = - \oint_l (\mathbf{u} \times \mathbf{B}) \cdot d\mathbf{l} + \int \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S}. \quad (4.62)$$

$$- \int \nabla \times (\mathbf{u} \times \mathbf{B}) \cdot d\mathbf{S} + \int \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} = - \int \left[ \nabla \times (\mathbf{u} \times \mathbf{B}) - \frac{\partial \mathbf{B}}{\partial t} \right] \cdot d\mathbf{S} = \mathbf{0}, \quad (4.63)$$

$$\nabla \times \mathbf{E} = \left( \frac{1}{R} \frac{\partial E_z}{\partial \varphi} - \frac{\partial E_R}{\partial z}, \frac{\partial E_R}{\partial z} - \frac{\partial E_z}{\partial R}, \frac{1}{R} \frac{\partial}{\partial R} R E_\varphi - \frac{1}{R} \frac{\partial E_R}{\partial \varphi} \right) = \mathbf{0}. \quad (4.64)$$

$$E_\varphi = -(\mathbf{u} \times \mathbf{B})_\varphi = -u_R B_z + u_z B_R = 0, \quad (4.65)$$

$$\frac{u_R}{u_z} = \frac{B_R}{B_z}. \quad (4.66)$$

$$\mathbf{u}_p = \xi(R, z) \mathbf{B}_p, \quad (4.67)$$

$$\Psi(R, z) = R A_\varphi, \quad (4.68)$$

$$B_R = -\frac{1}{R} \frac{\partial \Psi}{\partial z} \quad \text{and} \quad B_z = \frac{1}{R} \frac{\partial \Psi}{\partial R}. \quad (4.69)$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}_p) + \frac{\partial \rho v_\varphi}{\partial \varphi} = \nabla \cdot (\rho \mathbf{u}_p) = 0, \quad (4.70)$$

$$\nabla \cdot (\rho \mathbf{u}_p) = \nabla \cdot (\rho \xi \mathbf{B}_p) = \nabla \rho \xi \cdot \mathbf{B}_p + \rho \xi \nabla \cdot \mathbf{B}_p = 0. \quad (4.71)$$

$$\nabla \cdot (\rho \xi \mathbf{B}_p) = \nabla_R (\rho \xi) B_R + \nabla_z (\rho \xi) B_z = -\frac{1}{R} \nabla_R (\rho \xi) \nabla_z \Psi + \frac{1}{R} \nabla_z (\rho \xi) \nabla_R \Psi = 0, \quad (4.72)$$

$$\nabla (\rho \xi) \times \nabla \Psi = \mathbf{0}. \quad (4.73)$$

$$4\pi \rho \xi = F_1(\Psi). \quad (4.74)$$

$$\begin{aligned} \mathbf{u} \times \mathbf{B} &= u_\varphi \mathbf{e}_\varphi \times \mathbf{B}_p + \mathbf{u}_p \times B_\varphi \mathbf{e}_\varphi = u_\varphi \mathbf{e}_\varphi \times \mathbf{B}_p + \xi \mathbf{B}_p \times B_\varphi \mathbf{e}_\varphi \\ &= (u_\varphi - \xi B_\varphi) (B_z - B_R) \mathbf{e}_z = (u_\varphi - \xi B_\varphi) \frac{1}{R} \left( \frac{\partial \Psi}{\partial R} - \frac{\partial \Psi}{\partial z} \right) \mathbf{e}_z \\ &= \frac{u_\varphi - \xi B_\varphi}{R} \nabla \Psi. \end{aligned} \quad (4.75)$$

$$\nabla \times (\mathbf{u} \times \mathbf{B}) = \frac{\partial \mathbf{B}}{\partial t} = \mathbf{0}, \quad (4.76)$$

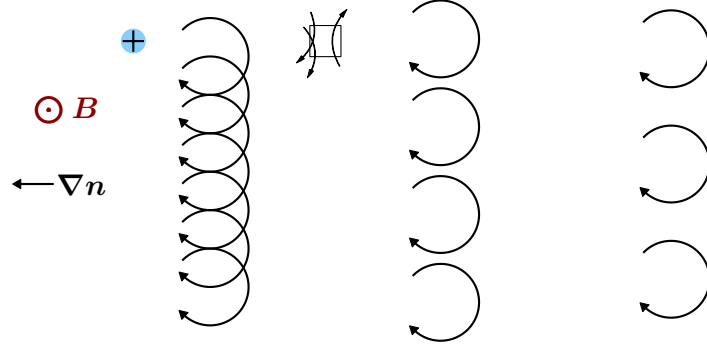
$$\nabla \times \left( \frac{u_\varphi - \xi B_\varphi}{R} \nabla \Psi \right) = \nabla \left( \frac{u_\varphi - \xi B_\varphi}{R} \right) \times \nabla \Psi = \mathbf{0}, \quad (4.77)$$

$$\frac{u_\varphi - \xi B_\varphi}{R} = F_2(\Psi). \quad (4.78)$$

$$B_\varphi \mathbf{e}_\varphi \cdot \nabla (R B_\varphi - F_1 R v_\varphi) = 0, \quad (4.79)$$

$$R B_\varphi - F_1(\Psi) R v_\varphi = F_3(\Psi). \quad (4.80)$$

$$\frac{1}{2} u^2 + \int_{\Psi=\text{const}} \frac{dP}{\rho} + \Phi - R u_\varphi F_3(\Psi) = F_4(\Psi). \quad (4.81)$$



### 4.5 Drifts in the plasma fluid

$$mn \left( \frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = qn(\mathbf{E} + \mathbf{u} \times \mathbf{B}) - \nabla p \quad (4.82)$$

$$0 = qn(\mathbf{E} + \mathbf{u} \times \mathbf{B}) - \nabla p \quad (4.83)$$

$$\begin{aligned} 0 &= qn[\mathbf{E} \times \mathbf{B} + (\mathbf{u} \times \mathbf{B}) \times \mathbf{B}] - \nabla p \times \mathbf{B} = \\ &= qn(\mathbf{E} \times \mathbf{B} - \mathbf{u}B^2 + \mathbf{B}\mathbf{u} \cdot \mathbf{B}) - \nabla p \times \mathbf{B}. \end{aligned} \quad (4.84)$$

$$qn(\mathbf{E} \times \mathbf{B} - \mathbf{u}_\perp B^2) - \nabla p \times \mathbf{B} = 0 \quad (4.85)$$

$$\mathbf{u}_\perp = \frac{\mathbf{E} \times \mathbf{B}}{B^2} - \frac{\nabla p \times \mathbf{B}}{qnB^2} = \mathbf{v}_E + \mathbf{v}_D. \quad (4.86)$$

$$mn \left( \frac{\partial u_z}{\partial t} + (\mathbf{u} \cdot \nabla)u_z \right) = qnE_z - \frac{\partial p}{\partial z} \quad (4.87)$$

$$\frac{\nabla p}{p} = \gamma \frac{\nabla n}{n}, \quad (4.88)$$

$$mn \frac{\partial u_z}{\partial t} = qnE_z - K_B T \frac{\partial n}{\partial z}. \quad (4.89)$$

$$e \frac{\partial \phi}{\partial z} = \frac{K_B T_e}{n} \frac{\partial n}{\partial z} \quad (4.90)$$

$$e\phi = K_B T_e \ln n + \mathcal{C} \quad (4.91)$$

$$n = n_0 \exp \left[ \frac{e\phi}{K_B T_e} \right], \quad (4.92)$$

## Chapter 5

# Waves in plasmas

### 5.1 Linear waves

$$\mathbf{v} = \mathbf{u} + \mathbf{w}, \quad (5.1)$$

$$\mathbf{u} = \mathbf{u}_0 + \mathbf{u}_1. \quad (5.2)$$

### 5.2 Example: Acoustic waves in fluids

$$\rho \left[ \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right] = -\nabla p = -\frac{\gamma p}{\rho} \nabla \rho, \quad (5.3)$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0. \quad (5.4)$$

$$(\rho_0 + \rho_1) \left\{ \frac{\partial \mathbf{u}_0}{\partial t} + \frac{\partial \mathbf{u}_1}{\partial t} + [(\mathbf{u}_0 + \mathbf{u}_1) \cdot \nabla] (\mathbf{u}_0 + \mathbf{u}_1) \right\} = -\frac{\gamma(p_0 + p_1)}{(\rho_0 + \rho_1)} \nabla (\rho_0 + \rho_1) \quad (5.5)$$

$$\frac{\partial(\rho_0 + \rho_1)}{\partial t} + \nabla \cdot [(\rho_0 + \rho_1)(\mathbf{u}_0 + \mathbf{u}_1)] = 0. \quad (5.6)$$

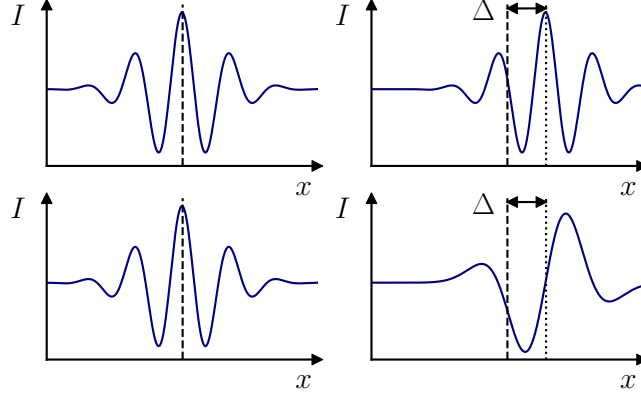
$$\frac{\overset{(1)}{\partial \rho_0}}{\partial t} + \frac{\overset{(2)}{\partial \rho_1}}{\partial t} + \nabla \cdot (\overset{(3)}{\rho_0 \mathbf{u}_0}) + \rho_0 \overset{(4)}{\nabla \cdot \mathbf{u}_1} + \mathbf{u}_1 \cdot \overset{(5)}{\nabla \rho_0} + \nabla \cdot (\overset{(6)}{\rho_1 \mathbf{u}_0}) + \nabla \cdot (\overset{(7)}{\rho_1 \mathbf{u}_1}) = 0. \quad (5.7)$$

$$\frac{\partial \rho_1}{\partial t} + \rho_0 \nabla \cdot \mathbf{u}_1 = 0 \quad (5.8)$$

$$\frac{p_0 + p_1}{\rho_0 + \rho_1} = \frac{\frac{p_0 + p_1}{\rho_0}}{\frac{\rho_0 + \rho_1}{\rho_0}} = \frac{\frac{p_0}{\rho_0} + \frac{p_1}{\rho_0}}{1 + \frac{\rho_1}{\rho_0}} = \frac{p_0}{\rho_0},$$

$$\rho_0 \frac{\partial \mathbf{u}_1}{\partial t} = -\gamma \frac{p_0}{\rho_0} \nabla \rho_1. \quad (5.9)$$

$$A_1 = \tilde{A}_1 \exp [i(\mathbf{k} \cdot \mathbf{r} - \omega t)], \quad (5.10)$$



$$\frac{\partial \bullet}{\partial t} \rightarrow -i\omega \bullet \quad (5.11)$$

$$\nabla \bullet \rightarrow i\mathbf{k} \bullet \quad (5.12)$$

$$\Delta \bullet \rightarrow -k^2 \bullet \quad (5.13)$$

$$-i\omega \rho_1 + i\rho_0 \mathbf{k} \cdot \mathbf{u}_1 = 0, \quad (5.14)$$

$$-i\omega \rho_0 \mathbf{u}_1 + i\gamma \frac{p_0}{\rho_0} \mathbf{k} \rho_1 = 0. \quad (5.15)$$

$$\frac{\omega}{k} = \sqrt{\frac{\gamma p_0}{\rho_0}} \equiv c_s, \quad (5.16)$$

$$-i\omega \rho_1 + i\rho_1 \mathbf{k} \cdot \mathbf{u}_0 + i\rho_0 \mathbf{k} \cdot \mathbf{u}_1 = 0, \quad (5.17)$$

$$-i\rho_1(\omega - \mathbf{k} \cdot \mathbf{u}_0) + i\rho_0 \mathbf{k} \cdot \mathbf{u}_1 = 0. \quad (5.18)$$

$$\rho_0 \frac{\partial \mathbf{u}_1}{\partial t} + \rho_0(\mathbf{u}_0 \cdot \nabla) \mathbf{u}_1 = -\gamma \frac{p_0}{\rho_0} \nabla \rho_1, \quad (5.19)$$

$$-i\rho_0(\omega - \mathbf{k} \cdot \mathbf{u}_0) \mathbf{u}_1 + i\mathbf{k} \gamma \frac{p_0}{\rho_0} \rho_1 = 0, \quad (5.20)$$

$$\omega - \mathbf{k} \cdot \mathbf{u}_0 = kc_s, \quad (5.21)$$

### 5.3 Types of plasma waves

### 5.4 Electrostatic waves without a background magnetic field

$$mn_0 \frac{\partial \mathbf{u}_{e1}}{\partial t} = -en_0 \mathbf{E}_1, \quad (5.22)$$

$$\frac{\partial n_{e1}}{\partial t} + \nabla \cdot [n_0 \mathbf{u}_{e1}] = 0, \quad (5.23)$$

$$\mathbf{E}_1 = -\nabla \phi_1, \quad (5.24)$$

$$\varepsilon_0 \nabla \cdot \mathbf{E}_1 = -en_{e1}. \quad (5.25)$$

$$\varepsilon_0 \Delta \phi_1 = n_{e1} e, \quad (5.26)$$

$$mn_0 \frac{\partial \nabla \cdot \mathbf{u}_{e1}}{\partial t} = en_0 \Delta \phi_1, \quad (5.27)$$

$$\frac{\partial n_{e1}}{\partial t} + n_0 \nabla \cdot \mathbf{u}_{e1} = 0. \quad (5.28)$$

$$mn_0 \frac{\partial \nabla \cdot \mathbf{u}_{e1}}{\partial t} = \frac{e^2 n_0}{\varepsilon_0} n_{e1}, \quad (5.29)$$

$$\nabla \cdot \mathbf{u}_{e1} = -\frac{1}{n_0} \frac{\partial n_{e1}}{\partial t}. \quad (5.30)$$

$$-m \frac{\partial^2 n_{e1}}{\partial t^2} = \frac{e^2 n_0}{\varepsilon_0} n_{e1}. \quad (5.31)$$

$$\omega^2 n_{e1} = \frac{e^2 n_0}{\varepsilon_0 m} n_{e1}. \quad (5.32)$$

$$\omega_p^2 = \frac{n_0 e^2}{\varepsilon_0 m}. \quad (5.33)$$

$$\nabla p_e = \nabla (3K_B T_e n_e) = 3K_B T_e \nabla (n_0 + n_{e1}), \quad (5.34)$$

$$mn_0 \frac{\partial \mathbf{u}_{e1}}{\partial t} = -en_0 \mathbf{E}_1 - 3K_B T_e \nabla n_{e1}, \quad (5.35)$$

$$\frac{\partial n_{e1}}{\partial t} + n_0 \nabla \cdot \mathbf{u}_{e1} = 0, \quad (5.36)$$

$$\mathbf{E}_1 = -\nabla \phi_1, \quad (5.37)$$

$$\varepsilon_0 \nabla \cdot \mathbf{E}_1 = -en_{e1}. \quad (5.38)$$

$$-m \frac{\partial^2 n_{e1}}{\partial t^2} = \frac{e^2 n_0}{\varepsilon_0} n_{e1} - 3K_B T_e \Delta n_{e1}. \quad (5.39)$$

$$\omega^2 n_{e1} - \omega_p^2 n_{e1} - 3k^2 \frac{K_B T_e}{m} n_{e1} = 0 \quad (5.40)$$

$$\omega^2 = \omega_p^2 + \frac{3}{2} k^2 v_T^2, \quad (5.41)$$

$$n_{e1} = \overline{n_{e1}} \cos(\mathbf{k} \cdot \mathbf{r} - \omega t) \quad (5.42)$$

$$\mathbf{E}_1 = \overline{\mathbf{E}_1} \cos(\mathbf{k} \cdot \mathbf{r} - \omega t + \delta_E). \quad (5.43)$$

$$-\varepsilon_0 \mathbf{k} \cdot \overline{\mathbf{E}_1} \sin(\mathbf{k} \cdot \mathbf{r} - \omega t + \delta_E) = -e \overline{n_{e1}} \cos(\mathbf{k} \cdot \mathbf{r} - \omega t). \quad (5.44)$$

$$\sin(\mathbf{k} \cdot \mathbf{r} - \omega t) \cos \delta_E + \cos(\mathbf{k} \cdot \mathbf{r} - \omega t) \sin \delta_E = \frac{e}{\varepsilon_0} \frac{\overline{n_{e1}}}{\overline{\mathbf{E}_1} \cdot \mathbf{k}} \cos(\mathbf{k} \cdot \mathbf{r} - \omega t). \quad (5.45)$$

$$-i\omega m \mathbf{u}_{e1} = e(\mathbf{E}_1 + \mathbf{u}_{e1} \times \mathbf{B}_0). \quad (5.46)$$

$$\mathbf{u}_{e1} = -\frac{e}{\omega} \mathbf{E}_1, \quad (5.47)$$

$$\mathbf{k} \times \mathbf{E}_1 = \omega \mathbf{B}_1, \quad (5.48)$$

$$m n_0 \frac{\partial \mathbf{u}_{e1}}{\partial t} = -e n_0 \mathbf{E}_1 - 3K_B T_e \nabla n_{e1}, \quad (5.49)$$

$$\frac{\partial n_{e1}}{\partial t} + n_0 \nabla \cdot \mathbf{u}_{e1} = 0, \quad (5.50)$$

$$M n_0 \frac{\partial \mathbf{u}_{i1}}{\partial t} = e n_0 \mathbf{E}_1 - 3K_B T_i \nabla n_{i1}, \quad (5.51)$$

$$\frac{\partial n_{i1}}{\partial t} + n_0 \nabla \cdot \mathbf{u}_{i1} = 0, \quad (5.52)$$

$$\mathbf{E}_1 = -\nabla \phi_1, \quad (5.53)$$

$$\varepsilon_0 \nabla \cdot \mathbf{E}_1 = e(n_{i1} - n_{e1}). \quad (5.54)$$

$$n_{e1} = n_e - n_{e0} = n_0 \left( \exp \frac{e\phi_1}{K_B T_e} - 1 \right) \sim n_0 \frac{e\phi_1}{K_B T_e} \quad (5.55)$$

$$M n_0 \frac{\partial \nabla \cdot \mathbf{u}_{i1}}{\partial t} = -e n_0 \Delta \phi_1 - 3K_B T_i \Delta n_{i1}, \quad (5.56)$$

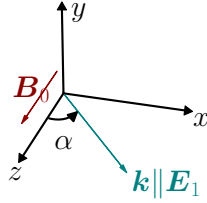
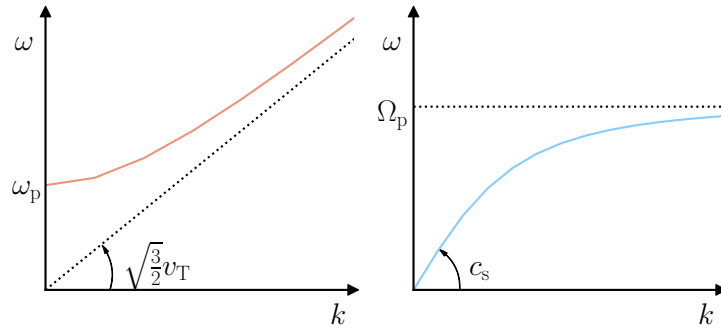
$$\frac{\partial n_{i1}}{\partial t} + n_0 \nabla \cdot \mathbf{u}_{i1} = 0, \quad (5.57)$$

$$-\varepsilon_0 \Delta \phi_1 = e n_{i1} - e n_0 \frac{e\phi_1}{K_B T_e}. \quad (5.58)$$

$$\omega^2 M n_{i1} = e n_0 k^2 \phi_1 + 3K_B T_i k^2 n_{i1}, \quad (5.59)$$

$$\varepsilon_0 k^2 \phi_1 = e n_{i1} - e^2 n_0 \frac{\phi_1}{K_B T_e}. \quad (5.60)$$

$$\omega^2 = \frac{e^2 n_0 k^2}{M \left( \varepsilon_0 k^2 + \frac{e^2 n_0}{K_B T_e} \right)} + \frac{3K_B T_i k^2}{M}. \quad (5.61)$$



$$\omega^2 = \frac{\Omega_p^2 \lambda_D^2 k^2}{\lambda_D^2 k^2 + 1} + \frac{3}{2} v_{Ti}^2 k^2. \quad (5.62)$$

$$\omega^2 = k^2 \left( \frac{K_B T_e}{M} + 3 \frac{K_B T_i}{M} \right). \quad (5.63)$$

## 5.5 Electrostatic waves in background magnetic field

$$-i\omega m \mathbf{u}_{e1} = -e(\mathbf{E}_1 + \mathbf{u}_{e1} \times \mathbf{B}_0), \quad (5.64)$$

$$-i\omega n_{e1} + in_0 \mathbf{k} \cdot \mathbf{u}_{e1} = 0, \quad (5.65)$$

$$i\mathbf{k} \cdot \mathbf{E}_1 = -\frac{e}{\varepsilon_0} n_{e1}, \quad (5.66)$$

$$\mathbf{E}_1 = \frac{i\omega m}{e} \mathbf{u}_{e1} - \mathbf{u}_{e1} \times \mathbf{B}_0, \quad (5.67)$$

$$\mathbf{k} \cdot \mathbf{E}_1 = \frac{ien_0}{\varepsilon_0 \omega} \mathbf{k} \cdot \mathbf{u}_{e1}. \quad (5.68)$$

$$E_1 = \frac{ien_0}{\varepsilon_0 \omega k} \mathbf{k} \cdot \mathbf{u}_{e1}. \quad (5.69)$$

$$\begin{pmatrix} E_1 \sin \alpha \\ 0 \\ E_1 \cos \alpha \end{pmatrix} = \frac{ien_0}{\varepsilon_0 \omega} (ku_{e1,x} \sin \alpha + ku_{e1,z} \cos \alpha) \begin{pmatrix} \sin \alpha \\ 0 \\ \cos \alpha \end{pmatrix} = \quad (5.70)$$

$$= \frac{i\omega m}{e} \begin{pmatrix} u_{e1,x} \\ u_{e1,y} \\ u_{e1,z} \end{pmatrix} - B_0 \begin{pmatrix} u_{e1,y} \\ -u_{e1,x} \\ 0 \end{pmatrix}. \quad (5.71)$$

$$\begin{pmatrix} u_{e1,x} \\ u_{e1,y} \\ u_{e1,z} \end{pmatrix} - \frac{\omega_p^2}{\omega^2} (u_{e1,x} \sin \alpha + u_{e1,z} \cos \alpha) \begin{pmatrix} \sin \alpha \\ 0 \\ \cos \alpha \end{pmatrix} + i \frac{\omega_c}{\omega} \begin{pmatrix} u_{e1,y} \\ -u_{e1,x} \\ 0 \end{pmatrix} = 0. \quad (5.72)$$

$$M_u \cdot \mathbf{u}_{e1} = 0, \quad (5.73)$$

$$M_u = \begin{bmatrix} 1 - \frac{\omega_p^2}{\omega^2} \sin^2 \alpha & i \frac{\omega_c}{\omega} & -\frac{\omega_p^2}{\omega^2} \sin \alpha \cos \alpha \\ -i \frac{\omega_c}{\omega} & 1 & 0 \\ -\frac{\omega_p^2}{\omega^2} \sin \alpha \cos \alpha & 0 & 1 - \frac{\omega_p^2}{\omega^2} \cos^2 \alpha \end{bmatrix} \quad (5.74)$$

$$\det M_u = -\frac{\omega_p^4}{\omega^4} \sin^2 \alpha \cos^2 \alpha + \left(1 - \frac{\omega_p^2}{\omega^2} \cos^2 \alpha\right) \left(1 - \frac{\omega_p^2}{\omega^2} \sin^2 \alpha - \frac{\omega_c^2}{\omega^2}\right) = 0, \quad (5.75)$$

$$\left(1 - \frac{\omega_p^2}{\omega^2}\right) \left(1 - \frac{\omega_c^2}{\omega^2}\right) = 0, \quad (5.76)$$

$$\omega^2 = \omega_p^2, \quad \text{and} \quad \omega^2 = \omega_c^2. \quad (5.77)$$

$$1 - \frac{\omega_p^2}{\omega^2} - \frac{\omega_c^2}{\omega^2} = 0, \quad (5.78)$$

$$\omega^2 = \omega_p^2 + \omega_c^2 \equiv \omega_h^2. \quad (5.79)$$

$$-i\omega M \mathbf{u}_{i1} = -eik\phi_1 + e\mathbf{u}_{i1} \times \mathbf{B}_0 \quad (5.80)$$

$$-i\omega M u_{i1,x} = -eik\phi_1 + e u_{i1,y} B_0, \quad (5.81)$$

$$-i\omega M u_{i1,y} = -e u_{i1,x} B_0, \quad (5.82)$$

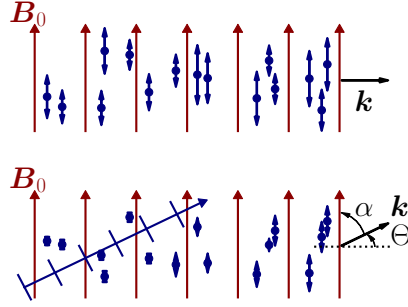
$$u_{i1,x} = \frac{ek\phi_1}{\omega M} \left(1 - \frac{\Omega_c^2}{\omega^2}\right)^{-1}. \quad (5.83)$$

$$n_{i1} = n_0 \frac{k}{\omega} u_{i1,x}, \quad (5.84)$$

$$n_{e1} = n_0 \frac{e\phi_1}{K_B T_e}. \quad (5.85)$$

$$u_{i1,x} = \frac{\omega e \phi_1}{K_B T_e k}, \quad (5.86)$$

$$\omega^2 - \Omega_c^2 = \frac{K_B T_e}{M} k^2. \quad (5.87)$$



$$c_s^2 = \left( \frac{K_B T_e}{M} + \frac{3K_B T_i}{M} \right) \quad (5.88)$$

$$\omega^2 = \Omega_c^2 + k^2 c_s^2. \quad (5.89)$$

$$u_{e1,x} = -\frac{ek\phi_1}{\omega m} \left( 1 - \frac{\omega_c^2}{\omega^2} \right)^{-1}. \quad (5.90)$$

$$n_{i1} = n_0 \frac{k}{\omega} u_{i1,x}, \quad (5.91)$$

$$n_{e1} = n_0 \frac{k}{\omega} u_{e1,x}. \quad (5.92)$$

$$M \left( 1 - \frac{\Omega_c^2}{\omega^2} \right) = -m \left( 1 - \frac{\omega_c^2}{\omega^2} \right), \quad (5.93)$$

$$\omega^2(M + m) = m\omega_c^2 + M\Omega_c^2 = e^2 B_0^2 \left( \frac{1}{m} + \frac{1}{M} \right) = e^2 B_0^2 \left( \frac{m + M}{mM} \right) \quad (5.94)$$

$$\omega^2 = \frac{eB_0}{m} \frac{eB_0}{M} = \omega_c \Omega_c \equiv \omega_d^2. \quad (5.95)$$

## 5.6 Electromagnetic waves

$$\nabla \times \mathbf{E}_1 = -\frac{\partial \mathbf{B}_1}{\partial t}, \quad (5.96)$$

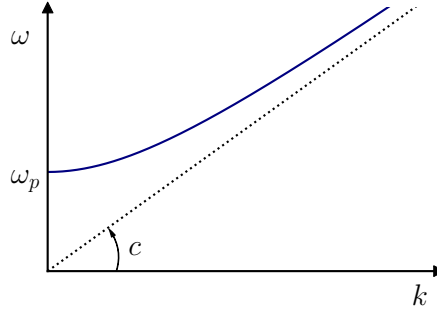
$$\frac{1}{\mu_0} \nabla \times \mathbf{B}_1 = \varepsilon_0 \frac{\partial \mathbf{E}_1}{\partial t}, \quad (5.97)$$

$$\nabla \cdot \mathbf{B}_1 = 0. \quad (5.98)$$

$$\nabla \times \frac{\partial \mathbf{E}_1}{\partial t} = -\frac{\partial^2 \mathbf{B}_1}{\partial t^2}, \quad (5.99)$$

$$\frac{1}{\mu_0 \varepsilon_0} \nabla \times \nabla \times \mathbf{B}_1 = \nabla \times \frac{\partial \mathbf{E}_1}{\partial t}, \quad (5.100)$$

$$\nabla \cdot \mathbf{B}_1 = 0. \quad (5.101)$$



$$-c^2 \mathbf{k} \times (\mathbf{k} \times \mathbf{B}_1) = \omega^2 \mathbf{B}_1, \quad (5.102)$$

$$\mathbf{k} \cdot \mathbf{B}_1 = 0, \quad (5.103)$$

$$-c^2 [(\mathbf{k} \cdot \mathbf{B}_1) \mathbf{k} - k^2 \mathbf{B}_1] = \omega^2 \mathbf{B}_1, \quad (5.104)$$

$$\mathbf{k} \cdot \mathbf{B}_1 = 0. \quad (5.105)$$

$$c^2 k^2 \mathbf{B}_1 = \omega^2 \mathbf{B}_1, \quad (5.106)$$

$$\omega^2 = c^2 k^2. \quad (5.107)$$

$$\nabla \times \mathbf{E}_1 = -\frac{\partial \mathbf{B}_1}{\partial t}, \quad (5.108)$$

$$\frac{1}{\mu_0} \nabla \times \mathbf{B}_1 = \varepsilon_0 \frac{\partial \mathbf{E}_1}{\partial t} + \mathbf{j}_1. \quad (5.109)$$

$$\nabla \times \nabla \times \mathbf{E}_1 = -\frac{1}{c^2} \frac{\partial^2 \mathbf{E}_1}{\partial t^2} - \frac{1}{\varepsilon_0 c^2} \frac{\partial \mathbf{j}_1}{\partial t}. \quad (5.110)$$

$$-(\mathbf{k} \cdot \mathbf{E}_1) \mathbf{k} + k^2 \mathbf{E}_1 = \frac{i\omega}{\varepsilon_0 c^2} \mathbf{j}_1 + \frac{\omega^2}{c^2} \mathbf{E}_1. \quad (5.111)$$

$$(\omega^2 - c^2 k^2) \mathbf{E}_1 = -\frac{i\omega}{\varepsilon_0} \mathbf{j}_1. \quad (5.112)$$

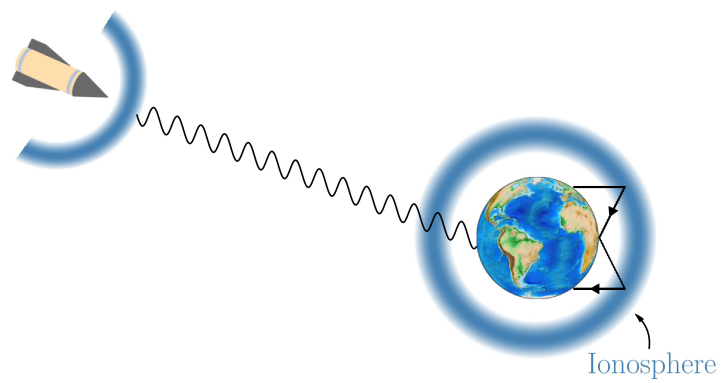
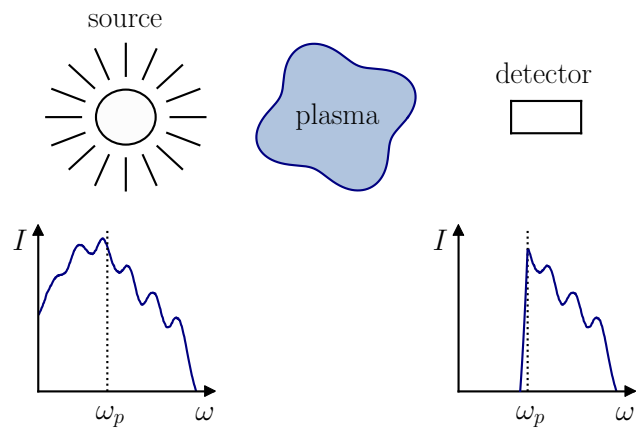
$$\mathbf{j}_1 = -n_0 e \mathbf{u}_{e1}. \quad (5.113)$$

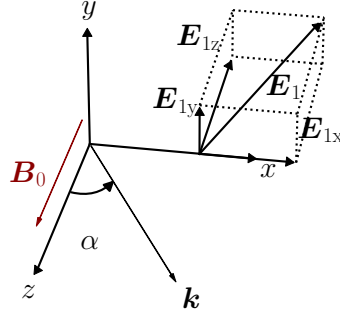
$$m \frac{\partial \mathbf{u}_{e1}}{\partial t} = -e \mathbf{E}_1. \quad (5.114)$$

$$\mathbf{j}_1 = -\frac{n_0 e^2}{i\omega m} \mathbf{E}_1, \quad (5.115)$$

$$(\omega^2 - c^2 k^2) \mathbf{E}_1 = \frac{i\omega}{\varepsilon_0} \frac{n_0 e^2}{i\omega m} \mathbf{E}_1 = \frac{n_0 e^2}{\varepsilon_0 m} \mathbf{E}_1 = \omega_p^2 \mathbf{E}_1. \quad (5.116)$$

$$\omega^2 = \omega_p^2 + c^2 k^2. \quad (5.117)$$





$$-i\omega m \mathbf{u}_{e1} = -e(\mathbf{E}_1 + \mathbf{u}_{e1} \times \mathbf{B}_0), \quad (5.118)$$

$$i\mathbf{k} \times \mathbf{E}_1 = i\omega \mathbf{B}_1, \quad (5.119)$$

$$\frac{i}{\mu_0} \mathbf{k} \times \mathbf{B}_1 = -\varepsilon_0 i\omega \mathbf{E}_1 - en_0 \mathbf{u}_{e1}. \quad (5.120)$$

$$\mathbf{u}_{e1} = -i \frac{e}{m\omega} \mathbf{E}_1 - i \frac{eB_0}{m\omega} \mathbf{u}_{e1} \times \mathbf{b}_0, \quad (5.121)$$

$$\mathbf{u}_{e1} = -i \frac{\varepsilon_0 \omega}{en_0} \mathbf{E}_1 - i \frac{c^2 \varepsilon_0 \omega}{\omega^2 en_0} \mathbf{k} \times (\mathbf{k} \times \mathbf{E}_1), \quad (5.122)$$

$$\begin{aligned} -i \frac{\varepsilon_0 \omega}{en_0} \mathbf{E}_1 - i \frac{c^2 \varepsilon_0}{\omega en_0} \mathbf{k} \times (\mathbf{k} \times \mathbf{E}_1) = \\ -i \frac{e}{m\omega} \mathbf{E}_1 - i \frac{\omega_c}{\omega} \left[ -i \frac{\varepsilon_0 \omega}{en_0} \mathbf{E}_1 - i \frac{c^2 \varepsilon_0}{\omega en_0} \mathbf{k} \times (\mathbf{k} \times \mathbf{E}_1) \right] \times \mathbf{b}_0. \end{aligned} \quad (5.123)$$

$$\begin{aligned} \omega^2 \mathbf{E}_1 + c^2 (\mathbf{k} \cdot \mathbf{E}_1) \mathbf{k} - c^2 k^2 \mathbf{E}_1 = \\ = \frac{e^2 n_0}{\varepsilon_0 m} \mathbf{E}_1 - i\omega_c \omega \mathbf{E}_1 \times \mathbf{b}_0 - i \frac{\omega_c}{\omega} c^2 (\mathbf{k} \cdot \mathbf{E}_1) (\mathbf{k} \times \mathbf{b}_0) + i \frac{\omega_c}{\omega} c^2 k^2 (\mathbf{E}_1 \times \mathbf{b}_0). \end{aligned} \quad (5.124)$$

$$(\omega^2 - \omega_p^2 - c^2 k^2) \mathbf{E}_1 + i \frac{\omega_c}{\omega} (\omega^2 - c^2 k^2) (\mathbf{E}_1 \times \mathbf{b}_0) + c^2 (\mathbf{k} \cdot \mathbf{E}_1) \mathbf{k} + i \frac{\omega_c}{\omega} c^2 (\mathbf{k} \cdot \mathbf{E}_1) (\mathbf{k} \times \mathbf{b}_0) = 0. \quad (5.125)$$

$$\begin{aligned}
& (\omega^2 - \omega_p^2 - c^2 k^2) \begin{pmatrix} E_{1x} \\ E_{1y} \\ E_{1z} \end{pmatrix} + i \frac{\omega_c}{\omega} (\omega^2 - c^2 k^2) \begin{pmatrix} E_{1y} \\ -E_{1x} \\ 0 \end{pmatrix} + \\
& + c^2 (k E_{1x} \sin \alpha + k E_{1z} \cos \alpha) \begin{pmatrix} k \sin \alpha \\ 0 \\ k \cos \alpha \end{pmatrix} + \\
& + i \frac{\omega_c}{\omega} c^2 (E_{1x} k \sin \alpha + k E_{1z} \cos \alpha) \begin{pmatrix} 0 \\ -k \sin \alpha \\ 0 \end{pmatrix} = 0. \tag{5.126}
\end{aligned}$$

$$M_{\mathbf{E}_1} \cdot \mathbf{E}_1 = 0 \tag{5.127}$$

$$M_{\mathbf{E}_1} = \begin{pmatrix} \omega^2 - \omega_p^2 - c^2 k^2 \cos^2 \alpha & i \frac{\omega_c}{\omega} (\omega^2 - c^2 k^2) & c^2 k^2 \sin \alpha \cos \alpha \\ -i \frac{\omega_c}{\omega} (\omega^2 - c^2 k^2 \cos^2 \alpha) & \omega^2 - \omega_p^2 - c^2 k^2 & -i \frac{\omega_c}{\omega} c^2 k^2 \cos \alpha \sin \alpha \\ c^2 k^2 \sin \alpha \cos \alpha & 0 & \omega^2 - \omega_p^2 - c^2 k^2 \sin^2 \alpha \end{pmatrix}. \tag{5.128}$$

$$\begin{aligned}
\det M_{\mathbf{E}_1} &= c^4 k^4 \sin^2 \alpha \cos^2 \alpha \left[ \frac{\omega_c^2}{\omega^2} (\omega^2 - c^2 k^2) - (\omega^2 - \omega_p^2 - c^2 k^2) \right] + \\
&+ (\omega^2 - \omega_p^2 - c^2 k^2 \sin^2 \alpha) \left[ (\omega^2 - \omega_p^2 - c^2 k^2 \cos^2 \alpha) (\omega^2 - \omega_p^2 - c^2 k^2) - \right. \\
&\left. - \frac{\omega_c^2}{\omega^2} (\omega^2 - c^2 k^2) (\omega^2 - c^2 k^2 \cos^2 \alpha) \right] = 0. \tag{5.129}
\end{aligned}$$

$$(\omega^2 - \omega_p^2) \left[ (\omega^2 - \omega_p^2 - c^2 k^2)^2 - \frac{\omega_c^2}{\omega^2} (\omega^2 - c^2 k^2)^2 \right] = 0. \tag{5.130}$$

$$\omega^2 = \omega_p^2, \tag{5.131}$$

$$\omega^2 - \omega_p^2 - c^2 k^2 = \pm \frac{\omega_c}{\omega} (\omega^2 - c^2 k^2) \tag{5.132}$$

$$\omega^2 - c^2 k^2 = \frac{\omega_p^2}{1 - \frac{\omega_c}{\omega}} \tag{5.133}$$

$$\omega^2 - c^2 k^2 = \frac{\omega_p^2}{1 + \frac{\omega_c}{\omega}}. \tag{5.134}$$

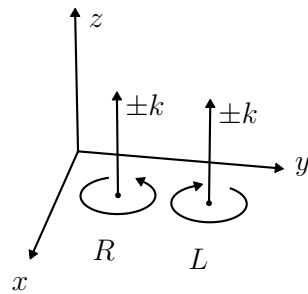
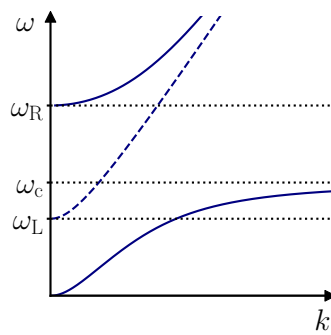
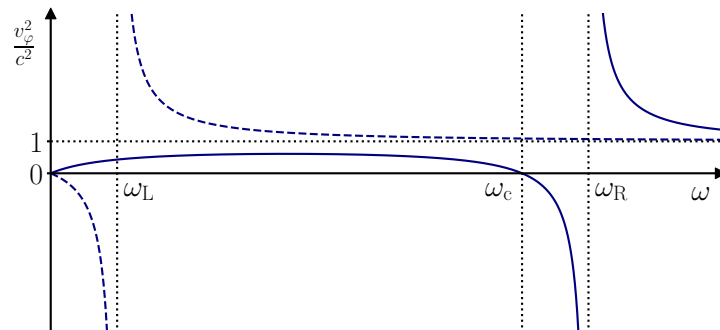
$$\begin{pmatrix} \omega^2 - \omega_p^2 - c^2 k^2 & i \frac{\omega_c}{\omega} (\omega^2 - c^2 k^2) & 0 \\ -i \frac{\omega_c}{\omega} (\omega^2 - c^2 k^2) & \omega^2 - \omega_p^2 - c^2 k^2 & 0 \\ 0 & 0 & \omega^2 - \omega_p^2 \end{pmatrix} \cdot \begin{pmatrix} E_{1x} \\ E_{1y} \\ E_{1z} \end{pmatrix} = 0, \tag{5.135}$$

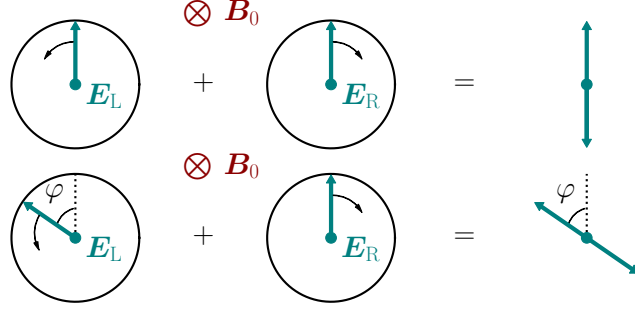
$$(\omega^2 - \omega_p^2 - c^2 k^2) E_{1x} + i \frac{\omega_c}{\omega} (\omega^2 - c^2 k^2) E_{1y} = 0. \tag{5.136}$$

$$E_{1y} = \pm i E_{1x}. \tag{5.137}$$

$$\mathbf{E}_1 = \overline{\mathbf{E}_1} \exp [i(\mathbf{k} \cdot \mathbf{r} - \omega t)] = \overline{\mathbf{E}_1} [\cos(\mathbf{k} \cdot \mathbf{r} - \omega t) + i \sin(\mathbf{k} \cdot \mathbf{r} - \omega t)]. \tag{5.138}$$

$$E_x = \Re(\mathbf{E}_1) = \overline{\mathbf{E}_1} \cos(\mathbf{k} \cdot \mathbf{r} - \omega t) \tag{5.139}$$





$$E_y = \Im(\mathbf{E}_1) = \overline{\mathbf{E}_1} \sin(\mathbf{k} \cdot \mathbf{r} - \omega t + \delta_E), \quad (5.140)$$

$$\mathbf{E}_R = E_0 \exp[i(\mathbf{k}_R \cdot \mathbf{r} - \omega t)] [\mathbf{e}_x + i\mathbf{e}_y], \quad (5.141)$$

$$\mathbf{E}_L = E_0 \exp[i(\mathbf{k}_L \cdot \mathbf{r} - \omega t)] [\mathbf{e}_x - i\mathbf{e}_y]. \quad (5.142)$$

$$\mathbf{E} = \frac{1}{2}(\mathbf{E}_R + \mathbf{E}_L) = \frac{1}{2}E_0 e^{-i\omega t} \left[ e^{i\mathbf{k}_R \cdot \mathbf{r}} (\mathbf{e}_x + i\mathbf{e}_y) + e^{i\mathbf{k}_L \cdot \mathbf{r}} (\mathbf{e}_x - i\mathbf{e}_y) \right]. \quad (5.143)$$

$$\mathbf{E} = \frac{1}{2}E_0 e^{-i\omega t} \left[ 2e^{i\mathbf{k} \cdot \mathbf{r}} \mathbf{e}_x \right] = E_0 \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)] \mathbf{e}_x, \quad (5.144)$$

$$k_{R,L} = \frac{\omega}{c} \sqrt{1 - \frac{\omega_p^2/\omega^2}{1 \mp \omega_c/\omega}}. \quad (5.145)$$

$$k_{R,L} \sim \frac{\omega}{c} \left( 1 - \frac{1}{2} \frac{\omega_p^2}{\omega^2} \frac{1}{1 \mp \omega_c/\omega} \right) \sim \frac{\omega}{c} \left[ 1 - \frac{1}{2} \frac{\omega_p^2}{\omega^2} \left( 1 \pm \frac{\omega_c}{\omega} \right) \right] \quad (5.146)$$

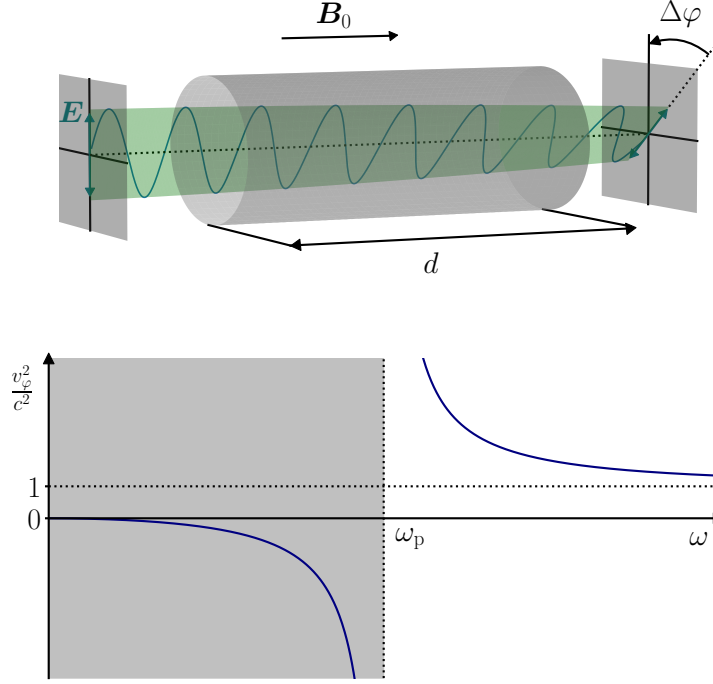
$$k_{R,L} = k \pm \Delta k, \quad (5.147)$$

$$k = \frac{\omega}{c} \left[ 1 - \frac{1}{2} \frac{\omega_p^2}{\omega^2} \right] \quad (5.148)$$

$$\Delta k = -\frac{1}{2} \frac{\omega_p^2 \omega_c}{c \omega^2}. \quad (5.149)$$

$$\begin{aligned} \mathbf{E} &= \frac{1}{2} E_0 e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \left[ e^{i\Delta \mathbf{k} \cdot \mathbf{r}} (\mathbf{e}_x + i\mathbf{e}_y) + e^{-i\Delta \mathbf{k} \cdot \mathbf{r}} (\mathbf{e}_x - i\mathbf{e}_y) \right] \\ &= E_0 e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} [\cos(-\Delta \mathbf{k} \cdot \mathbf{r}) \mathbf{e}_x + \sin(-\Delta \mathbf{k} \cdot \mathbf{r}) \mathbf{e}_y]. \end{aligned} \quad (5.150)$$

$$\frac{d\varphi}{dz} = -\Delta k. \quad (5.151)$$



$$\begin{aligned}
 \varphi &= \varphi_0 + \int_0^d \frac{d\varphi}{dz} dz = \varphi_0 + \int_0^d -\Delta k dz = \varphi_0 + \frac{1}{2c\omega^2} \int_0^d \omega_p^2(z) \omega_c(z) dz = \\
 &= \varphi_0 + \frac{e^3}{2m^2 \varepsilon_0 c \omega^2} \int_0^d n_0(z) B_0(z) dz.
 \end{aligned} \tag{5.152}$$

$$\varphi = \varphi_0 + \frac{e^3}{2m^2 \varepsilon_0 c \omega^2} n_0 B_0 d. \tag{5.153}$$

$$(\omega^2 - \omega_p^2 - c^2 k^2) [(\omega^2 - \omega_p^2)(\omega^2 - \omega_p^2 - c^2 k^2) - \omega_c^2(\omega^2 - c^2 k^2)] = 0. \tag{5.154}$$

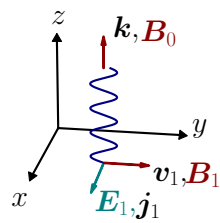
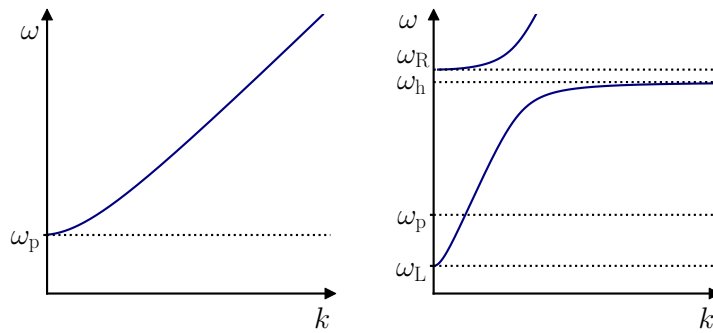
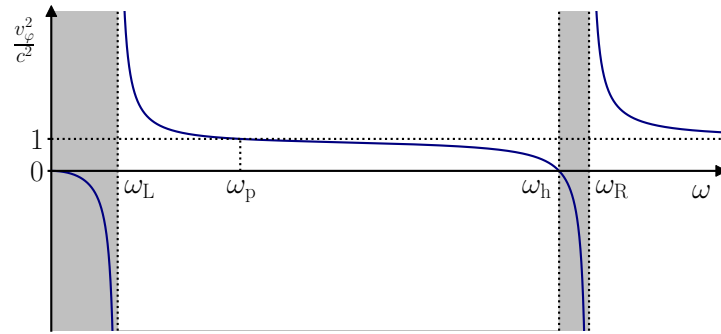
$$\omega^2 = \omega_p^2 + c^2 k^2, \tag{5.155}$$

$$(\omega^2 - \omega_p^2)(\omega^2 - \omega_p^2 - c^2 k^2) = \omega_c^2(\omega^2 - c^2 k^2), \tag{5.156}$$

$$\omega^2 - c^2 k^2 = \omega_p^2 \frac{\omega^2 - \omega_p^2}{\omega^2 - \omega_h^2}, \tag{5.157}$$

## 5.7 MHD waves

$$(\omega^2 - c^2 k^2) \mathbf{E}_1 = -\frac{i\omega}{\varepsilon_0} \mathbf{j}_1. \tag{5.158}$$



$$\mathbf{j}_1 = j_1 \mathbf{e}_x = n_0 e (u_{i1,x} - u_{e1,x}) \mathbf{e}_x. \quad (5.159)$$

$$M \frac{\partial \mathbf{u}_{i1}}{\partial t} = e (\mathbf{E}_1 + \mathbf{u}_{i1} \times \mathbf{B}_0), \quad (5.160)$$

$$-i\omega M u_{i1,x} = e E_1 + e u_{i1,y} B_0, \quad (5.161)$$

$$-i\omega M u_{i1,y} = -e u_{i1,x} B_0. \quad (5.162)$$

$$-i\omega M u_{i1,x} = e E_1 + e B_0 \frac{e B_0}{i\omega M} u_{i1,x} \quad (5.163)$$

$$u_{i1,x} = \frac{ie}{\omega M} \left(1 - \frac{\Omega_c^2}{\omega^2}\right)^{-1} E_1. \quad (5.164)$$

$$u_{i1,y} = \frac{e}{\omega M} \frac{\Omega_c}{\omega} \left(1 - \frac{\Omega_c^2}{\omega^2}\right)^{-1} E_1. \quad (5.165)$$

$$u_{e1,x} = -\frac{ie}{\omega m} \left(1 - \frac{\omega_c^2}{\omega^2}\right)^{-1} E_1, \quad (5.166)$$

$$u_{e1,y} = -\frac{e}{\omega m} \frac{-\omega_c}{\omega} \left(1 - \frac{\omega_c^2}{\omega^2}\right)^{-1} E_1. \quad (5.167)$$

$$u_{e1,x} = -\frac{ie}{\omega m} \left(\frac{\omega^2 - \omega_c^2}{\omega^2}\right)^{-1} E_1 \sim \frac{ie}{\omega m} \frac{\omega^2}{\omega_c^2} E_1 \rightarrow 0 \quad (5.168)$$

$$u_{e1,y} = \frac{e}{\omega m} \frac{\omega_c}{\omega} \left(\frac{\omega^2 - \omega_c^2}{\omega^2}\right)^{-1} E_1 \sim -\frac{e}{\omega m} \frac{\omega_c \omega^2}{\omega_c^2} E_1 = -\frac{e}{\omega_c m} E_1 = -\frac{E_1}{B_0}. \quad (5.169)$$

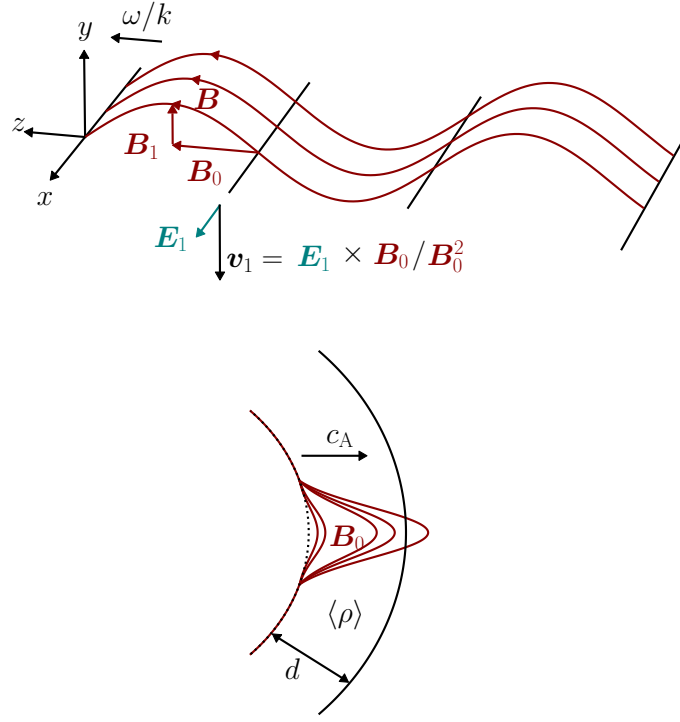
$$\begin{aligned} (\omega^2 - c^2 k^2) E_1 &= -\frac{i\omega}{\varepsilon_0} n_0 e \left[ \frac{ie}{\omega M} \left(1 - \frac{\Omega_c^2}{\omega^2}\right)^{-1} E_1 - 0 \right] = \\ &= \frac{n_0 e^2}{\varepsilon_0 M} \left(1 - \frac{\Omega_c^2}{\omega^2}\right)^{-1} E_1 = \Omega_p^2 \left(1 - \frac{\Omega_c^2}{\omega^2}\right)^{-1} E_1. \end{aligned} \quad (5.170)$$

$$\omega^2 - c^2 k^2 = \Omega_p^2 \left(1 - \frac{\Omega_c^2}{\omega^2}\right)^{-1}. \quad (5.171)$$

$$\omega^2 - c^2 k^2 = \Omega_p^2 \left(\frac{\omega^2 - \Omega_c^2}{\omega^2}\right)^{-1} \sim -\omega^2 \frac{\Omega_p^2}{\Omega_c^2} = -\omega^2 \frac{n_0 M}{\varepsilon_0 B_0^2} = -\omega^2 \frac{\rho_0}{\varepsilon_0 B_0^2}. \quad (5.172)$$

$$\frac{\omega^2}{k^2} = \frac{c^2}{1 + \frac{\rho_0}{\varepsilon_0 B_0^2}}. \quad (5.173)$$

$$\frac{\omega^2}{k^2} = \frac{c^2}{\frac{\rho \mu_0 c^2}{B_0^2}} = \frac{B_0^2}{\mu_0 \rho_0}. \quad (5.174)$$



$$\frac{\omega}{k} = \frac{B_0}{\sqrt{\mu_0 \rho_0}} \equiv c_A. \quad (5.175)$$

$$\frac{\partial \rho}{\partial t} + \rho \nabla \cdot \mathbf{u} = 0, \quad (5.176)$$

$$\rho \frac{d\mathbf{u}}{dt} = -\nabla p + (\nabla \times \mathbf{B}) \times \mathbf{B} / \mu_0, \quad (5.177)$$

$$\frac{d}{dt} \left( \frac{p}{\rho^\gamma} \right) = 0, \quad (5.178)$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{u} \times \mathbf{B}), \quad (5.179)$$

$$\nabla \cdot \mathbf{B} = 0. \quad (5.180)$$

$$\frac{\partial \rho_1}{\partial t} + \rho_0 (\nabla \cdot \mathbf{u}_1) = 0, \quad (5.181)$$

$$\rho_0 \frac{\partial \mathbf{u}_1}{\partial t} = -\nabla p_1 + (\nabla \times \mathbf{B}_1) \times \mathbf{B}_0 / \mu_0, \quad (5.182)$$

$$\frac{\partial \mathbf{B}_1}{\partial t} = \nabla \times (\mathbf{u}_1 \times \mathbf{B}_0), \quad (5.183)$$

$$\nabla \cdot \mathbf{B}_1 = 0. \quad (5.184)$$

$$\begin{aligned}
0 &= \left\{ \frac{\partial}{\partial t} + \mathbf{u}_1 \cdot \nabla \right\} \left[ \frac{p_0 + p_1}{(\rho_0 + \rho_1)^\gamma} \right] = \\
&= \frac{1}{(\rho_0 + \rho_1)^\gamma} \left[ \frac{\partial p_1}{\partial t} + (\mathbf{u}_1 \cdot \nabla) p_0 \right] + (p_0 + p_1) \left\{ \frac{-\gamma}{(\rho_0 + \rho_1)^{\gamma+1}} \left[ \frac{\partial \rho_1}{\partial t} + (\mathbf{u}_1 \cdot \nabla) \rho_0 \right] \right\}.
\end{aligned} \tag{5.185}$$

$$0 = \frac{1}{(\rho_0 + \rho_1)^\gamma} \left[ \frac{\partial p_1}{\partial t} - \gamma \frac{p_0 + p_1}{\rho_0 + \rho_1} \frac{\partial \rho_1}{\partial t} \right] = \frac{1}{(\rho_0 + \rho_1)^\gamma} \left\{ \frac{\partial p_1}{\partial t} - c_s^2 \frac{\partial \rho_1}{\partial t} \right\}, \tag{5.186}$$

$$\frac{\partial p_1}{\partial t} - c_s^2 \frac{\partial \rho_1}{\partial t} = 0. \tag{5.187}$$

$$\frac{\partial^2 \mathbf{u}_1}{\partial t^2} = -\frac{1}{\rho_0} \nabla \frac{\partial p_1}{\partial t} + (\nabla \times \frac{\partial \mathbf{B}_1}{\partial t}) \times \frac{\mathbf{B}_0}{\rho_0 \mu_0}. \tag{5.188}$$

$$\frac{\partial^2 \mathbf{u}_1}{\partial t^2} = -\frac{1}{\rho_0} \nabla \left( c_s^2 \frac{\partial \rho_1}{\partial t} \right) + (\nabla \times \frac{\partial \mathbf{B}_1}{\partial t}) \times \frac{\mathbf{B}_0}{\rho_0 \mu_0} \tag{5.189}$$

$$\frac{\partial^2 \mathbf{u}_1}{\partial t^2} = c_s^2 \nabla (\nabla \cdot \mathbf{u}_1) + \{ \nabla \times [\nabla \times (\mathbf{u}_1 \times \mathbf{B}_0)] \} \times \frac{\mathbf{B}_0}{\rho_0 \mu_0}. \tag{5.190}$$

$$\begin{aligned}
\omega^2 \mathbf{u}_1 &= c_s^2 \mathbf{k} (\mathbf{k} \cdot \mathbf{u}_1) + \{ \mathbf{k} \times [\mathbf{k} \times (\mathbf{u}_1 \times \mathbf{B}_0)] \} \times \frac{\mathbf{B}_0}{\mu_0 \rho_0} \\
&= c_s^2 \mathbf{k} (\mathbf{k} \cdot \mathbf{u}_1) + \{ \mathbf{k} \times [\mathbf{k} \times (\mathbf{u}_1 \times \mathbf{b}_0)] \} \times \mathbf{b}_0 c_A^2,
\end{aligned} \tag{5.191}$$

$$\begin{aligned}
\{ \mathbf{k} \times [\mathbf{k} \times (\mathbf{u}_1 \times \mathbf{b}_0)] \} \times \mathbf{b}_0 &= \{ \mathbf{k} \times [\mathbf{u}_1 (\mathbf{k} \cdot \mathbf{b}_0) - \mathbf{b}_0 (\mathbf{k} \cdot \mathbf{u}_1)] \} \times \mathbf{b}_0 = \\
&= (\mathbf{k} \cdot \mathbf{b}_0) [(\mathbf{k} \times \mathbf{u}_1) \times \mathbf{b}_0] - (\mathbf{k} \cdot \mathbf{u}_1) [(\mathbf{k} \times \mathbf{b}_0) \times \mathbf{b}_0] = \\
&= -(\mathbf{k} \cdot \mathbf{b}_0) [\mathbf{k} (\mathbf{u}_1 \cdot \mathbf{b}_0) - \mathbf{u}_1 (\mathbf{k} \cdot \mathbf{b}_0)] + (\mathbf{k} \cdot \mathbf{u}_1) [\mathbf{k} b_0^2 - \mathbf{b}_0 (\mathbf{k} \cdot \mathbf{b}_0)] = \\
&= (\mathbf{k} \cdot \mathbf{b}_0)^2 \mathbf{u}_1 - (\mathbf{k} \cdot \mathbf{u}_1) (\mathbf{k} \cdot \mathbf{b}_0) \mathbf{b}_0 + \mathbf{k} [\mathbf{k} \cdot \mathbf{u}_1 - (\mathbf{k} \cdot \mathbf{b}_0) (\mathbf{u}_1 \cdot \mathbf{b}_0)].
\end{aligned} \tag{5.192}$$

$$\omega^2 \mathbf{u}_1 = c_s^2 \mathbf{k} (\mathbf{k} \cdot \mathbf{u}_1) + c_A^2 \{ (\mathbf{k} \cdot \mathbf{b}_0)^2 \mathbf{u}_1 - (\mathbf{k} \cdot \mathbf{u}_1) (\mathbf{k} \cdot \mathbf{b}_0) \mathbf{b}_0 + \mathbf{k} [\mathbf{k} \cdot \mathbf{u}_1 - (\mathbf{k} \cdot \mathbf{b}_0) (\mathbf{u}_1 \cdot \mathbf{b}_0)] \}, \tag{5.193}$$

$$\frac{\omega^2 \mathbf{u}_1}{c_A^2} = k^2 \cos^2 \alpha \mathbf{u}_1 - (\mathbf{k} \cdot \mathbf{u}_1) k \cos \alpha \mathbf{b}_0 + \mathbf{k} \left[ \left( 1 + \frac{c_s^2}{c_A^2} \right) (\mathbf{k} \cdot \mathbf{u}_1) - k \cos \alpha (\mathbf{u}_1 \cdot \mathbf{b}_0) \right]. \tag{5.194}$$

$$\begin{aligned}
\omega^2 (\mathbf{k} \cdot \mathbf{u}_1) &= k^2 c_A^2 \cos^2 \alpha (\mathbf{u}_1 \cdot \mathbf{k}) - c_A^2 k^2 \cos^2 \alpha (\mathbf{k} \cdot \mathbf{u}_1) + \\
&+ c_A^2 (\mathbf{k} \cdot \mathbf{u}_1) k^2 + c_s^2 (\mathbf{k} \cdot \mathbf{u}_1) k^2 - c_A^2 k^3 \cos \alpha (\mathbf{u}_1 \cdot \mathbf{b}_0).
\end{aligned} \tag{5.195}$$

$$[-\omega^2 + c_A^2 k^2 + c_s^2 k^2] (\mathbf{u}_1 \cdot \mathbf{k}) = c_A^2 k^3 \cos \alpha (\mathbf{b}_0 \cdot \mathbf{u}_1). \tag{5.196}$$

$$\begin{aligned}
\omega^2 (\mathbf{b}_0 \cdot \mathbf{u}_1) &= c_A^2 k^2 \cos^2 \alpha (\mathbf{u}_1 \cdot \mathbf{b}_0) - c_A^2 (\mathbf{k} \cdot \mathbf{u}_1) k \cos \alpha + \\
&+ c_A^2 (\mathbf{k} \cdot \mathbf{u}_1) k \cos \alpha + c_s^2 (\mathbf{k} \cdot \mathbf{u}_1) k \cos \alpha - c_A^2 k^2 \cos^2 \alpha (\mathbf{u}_1 \cdot \mathbf{b}_0),
\end{aligned} \tag{5.197}$$

$$\omega^2(\mathbf{b}_0 \cdot \mathbf{u}_1) = c_s^2(\mathbf{k} \cdot \mathbf{u}_1)k \cos \alpha. \quad (5.198)$$

$$\frac{\mathbf{b}_0 \cdot \mathbf{u}_1}{\mathbf{u}_1 \cdot \mathbf{k}} = \frac{-\omega^2 + c_A^2 k^2 + c_s^2 k^2}{c_A^2 k^3 \cos \alpha} \quad (5.199)$$

$$\frac{\mathbf{b}_0 \cdot \mathbf{u}_1}{\mathbf{u}_1 \cdot \mathbf{k}} = \frac{c_s^2 k \cos \alpha}{\omega^2}. \quad (5.200)$$

$$\frac{-\omega^2 + c_A^2 k^2 + c_s^2 k^2}{c_A^2 k^3 \cos \alpha} = \frac{c_s^2 k \cos \alpha}{\omega^2} \quad (5.201)$$

$$\omega^4 - \omega^2 k^2 (c_A^2 + c_s^2) + c_s^2 c_A^2 k^4 \cos^2 \alpha = 0. \quad (5.202)$$

$$\left(\frac{\omega}{k}\right)^2 = \left[ \frac{1}{2}(c_s^2 + c_A^2) \pm \frac{1}{2} \sqrt{c_s^4 + c_A^4 - 4c_s^2 c_A^2 \cos^2 \alpha + 2c_A^2 c_s^2} \right]. \quad (5.203)$$

$$\left(\frac{\omega}{k}\right)^2 = \frac{1}{2}(c_s^2 + c_A^2) \pm \frac{1}{2} \sqrt{c_s^4 + c_A^4 - 2c_A^2 c_s^2} = \frac{1}{2}(c_s^2 + c_A^2) \pm \frac{1}{2} |c_s^2 - c_A^2| = \begin{cases} c_s^2 \\ c_A^2 \end{cases}, \quad (5.204)$$

$$\left(\frac{\omega}{k}\right)^2 = \frac{1}{2}(c_s^2 + c_A^2) \pm \frac{1}{2} \sqrt{c_s^4 + c_A^4 + 2c_A^2 c_s^2} = \frac{1}{2}(c_s^2 + c_A^2) \pm \frac{1}{2}(c_s^2 + c_A^2) = \begin{cases} c_s^2 + c_A^2 \\ 0 \end{cases}, \quad (5.205)$$

$$\left(\frac{\omega}{k}\right)^2 = \begin{cases} 0 \\ c_s^2 \end{cases}, \quad (5.206)$$

# Chapter 6

## Diffusion

### 6.1 Diffusion in weakly ionised plasmas

In Section ?? we introduced an additional term to the Euler equation accounting for collisions between particles. We defined the particle flux  $\Gamma$

$$\Gamma = n\mathbf{u}, \quad (6.1)$$

where  $n$  was the particle density and  $\mathbf{u}$  their speed.

The formalism used here is similar. For simplicity, let us assume that the particles flow along the  $x$ -axis of the Cartesian coordinate system. Due to collisions occurring in the spatial  $dx$ , a fraction of the flux  $d\Gamma$  is lost (see Fig. 6.1). We may write

$$d\Gamma = -n_n\sigma\Gamma dx, \quad (6.2)$$

where  $n_n$  is the density of neutral particles because we expect the collisions to be dominated by neutral particles.  $\sigma$  is the effective cross-section of the collisions. From the effective cross-section, we may define useful quantities  $\lambda_s = 1/n_n\sigma$  to be the mean free path between consecutive collisions and similarly the mean collisional frequency

$$\nu = n_n\langle\sigma v\rangle \approx n_n\sigma\mathbf{u}, \quad (6.3)$$

where the averaging is performed over the velocity space. The averaging is performed because quantum mechanics teaches us that the effective cross-section may be a function of velocity. The mean period between consecutive collisions is then  $\tau = 1/\nu$ .

We take the time derivative of (6.1) and (6.2), where for simplicity we neglect the possible time dependence of all variables apart from the coordinates and velocity. Then using the definition of  $\Gamma$  we have

$$\frac{d\Gamma}{dt} = n\frac{d\mathbf{u}}{dt} = -n_n\sigma n\mathbf{u}\frac{dx}{dt} = -n_n\sigma n\mathbf{u}\mathbf{u} = -n\nu\mathbf{u}, \quad (6.4)$$

Hence

$$\frac{d\mathbf{u}}{dt} = -\nu\mathbf{u}, \quad (6.5)$$

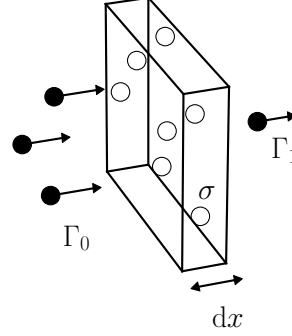


Figure 6.1: Illustration of the model of the change of the particle flux via collisions.

which is the (average) velocity change due to the collisions.

Hence the Euler equation with the magnetic field set to zero and with the collisional term added reads

$$mn \left[ \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right] = qn \mathbf{E} - \nabla p - mn\nu \mathbf{u} \quad (6.6)$$

Let's investigate a system where  $\frac{d}{dt} \mathbf{u} = 0$  (in such a system we only find a constant motion due to the diffusion). Then

$$\mathbf{u} = \frac{q\mathbf{E}}{m\nu} - \frac{K_B T \nabla n}{mn\nu} = \frac{\mathbf{\Gamma}}{n}, \quad (6.7)$$

where we used the definition of the particle flux. By introducing the mobility  $\mu = \frac{q \operatorname{sgn} q}{m\nu}$  and the diffusion coefficient  $D = \frac{K_B T}{m\nu}$  we have for the particle species  $j$

$$\mathbf{\Gamma}_j = \operatorname{sgn} q_j n \mu_j \mathbf{E} - D_j \nabla n. \quad (6.8)$$

For  $\mathbf{E} = 0$  we have

$$\mathbf{\Gamma} = -D \nabla n, \quad (6.9)$$

which is a form of *Fick's law* for diffusion. A simple interpretation of this useful relation is that the particles flow from locations with a higher density to lower density regions (as illustrated in Fig. 6.2), which is to be expected. In plasmas, Fick's law is generalised to (6.8) noting the effect of the electric field. One should also note that the coefficients  $\mu_j$  and  $D_j$  may generally be different for various particle species (most often electrons and ions) and hence the diffusion in plasmas may cause a separation of charges and thus drive an electric field. However, to fulfill the quasineutrality of plasma,  $\mathbf{\Gamma}_i = \mathbf{\Gamma}_e = \mathbf{\Gamma}$  on scales  $L \gg \lambda_D$ , which allows us to derive the necessary electric field that will keep the plasma quasineutral.

By assuming the plasma approximation  $n_e = n_i = n$  we have

$$n \mu_i \mathbf{E} - D_i \nabla n = -n \mu_e \mathbf{E} - D_e \nabla n \quad (6.10)$$

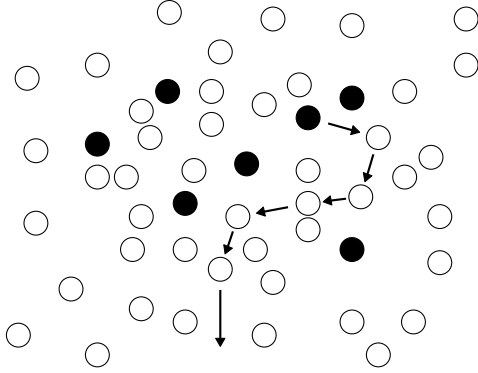


Figure 6.2: Illustrative concept of diffusion. An overabundance of the black “particles” is decreased in time via collisions with white “particles” and their local density thus decreases.

and hence

$$\mathbf{E} = \frac{D_i - D_e}{\mu_i + \mu_e} \frac{\nabla n}{n} \quad (6.11)$$

is the induced electric field triggered by the separation of charges. This is a repulsive electric field, which stops the charges from further separation due to diffusion and hence keeps plasma neutral on large scales. This electric field causes the electrons to slow down and accelerate ions simultaneously, so that the bulk of plasma (with charges separated on scales smaller than  $\lambda_D$ ) diffuses according to modified Fick’s law. This is the basic picture of *ambipolar diffusion*.

Then

$$\begin{aligned} \mathbf{\Gamma} &= \mathbf{\Gamma}_i = n\mu_i \frac{D_i - D_e}{\mu_i + \mu_e} \frac{\nabla n}{n} - D_i \nabla n = \\ &= \frac{\mu_i D_i - \mu_i D_e - \mu_i D_i - \mu_e D_i}{\mu_i + \mu_e} \nabla n = -\frac{\mu_i D_e + \mu_e D_i}{\mu_i + \mu_e} \nabla n = \\ &\equiv -D_a \nabla n, \end{aligned} \quad (6.12)$$

where  $D_a$  is the coefficient of the ambipolar diffusion. For  $\mu_e \gg \mu_i$  and using the definition of the mobility coefficient we have

$$D_a = \frac{\mu_i D_e + \mu_e D_i}{\mu_i + \mu_e} \approx D_i + \frac{\mu_i}{\mu_e} D_e = \left(1 + \frac{T_e}{T_i}\right) D_i. \quad (6.13)$$

Thus the coefficient of ambipolar diffusion is mostly determined by the diffusion coefficients for ions with some correction.

### Typical time evolution by diffusion

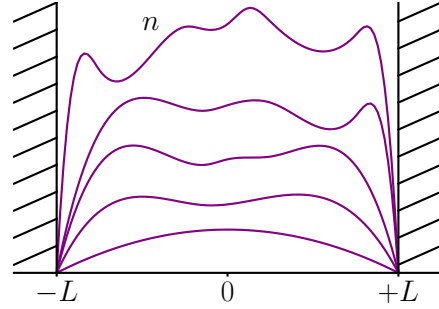
As an example, let’s take the continuum equation

$$\frac{\partial n}{\partial t} + \nabla \cdot n\mathbf{u} = 0, \quad (6.14)$$

and hence

$$\frac{\partial n}{\partial t} = -\nabla \cdot n\mathbf{u} = -\nabla \cdot \mathbf{\Gamma} = D_a \Delta n. \quad (6.15)$$

Figure 6.3: A model of plasma diffusion confined between two infinite plane-parallel walls. A series of plasma density snapshots is represented by the purple lines. At the beginning, the density profile is quite complex. Due to diffusion, it simplifies with time, large spatial frequencies (small-scale features) decay the fastest.



This equation is similar to the heat conduction equation. It may be solved by assuming the 1-D solution by separating the variables  $n(x, t) = T(t)S(x)$ , where  $T$  and  $S$  are some functions. Then

$$S \frac{dT}{dt} = D_a T \Delta S, \quad (6.16)$$

and hence

$$\frac{1}{T} \frac{dT}{dt} = \frac{D_a}{S} \Delta S = \text{const} = -\frac{1}{\tau}. \quad (6.17)$$

Both sides of the equations must equal the same constant because the left-hand side depends purely on time; thus, the right-hand side can't depend on time and vice versa for the spatial dependency of the right-hand side. For dimension reasons, this constant must have the dimension of reciprocal time. The equation may then be solved by parts

$$\frac{dT}{dt} = -\frac{T}{\tau} \quad \rightarrow \quad T = T_0 e^{-\frac{t}{\tau}} \quad (6.18)$$

and

$$\frac{d^2 S}{dx^2} = -\frac{1}{D_a \tau} S, \quad (6.19)$$

which has the solution in the series of harmonic functions,

$$S = S_0 \cos \frac{x}{\sqrt{D_a \tau}} \quad \text{and} \quad S = S_0 \sin \frac{x}{\sqrt{D_a \tau}}. \quad (6.20)$$

This solution must be supplied with the boundary conditions. Let's assume a simple model of plasma confined between two infinite plane-parallel walls (see Fig. 6.3), where the density of the plasma vanishes outside interval  $x \in (-L, L)$ . Then the solution to the problem is

$$n = n_0 \left\{ \sum_{l=0}^{\infty} a_{e,l} e^{-\frac{t}{\tau_{e,l}}} \cos \left[ \frac{(l + \frac{1}{2})\pi x}{L} \right] + \sum_{m=1}^{\infty} a_{o,m} e^{-\frac{t}{\tau_{o,m}}} \sin \left[ \frac{m\pi x}{L} \right] \right\} \quad (6.21)$$

The coefficients  $\tau_{e,l}$  and  $\tau_{o,m}$  may be determined by inserting the solution above into (6.16) and comparing the terms one by one. For the first term ( $l = 0$ ) in the cosine series we have

$$-a_{e,0}e^{-\frac{t}{\tau_{e,0}}}\frac{1}{\tau_{e,0}}\cos\frac{\pi x}{2L} = -D_a a_{e,0}e^{-\frac{t}{\tau_{e,0}}}\left(\frac{\pi}{2L}\right)^2\cos\frac{\pi x}{2L} \quad (6.22)$$

and hence

$$\tau_{e,0} = \frac{4L^2}{\pi^2 D_a}. \quad (6.23)$$

For a general mode  $l$ , we would similarly find

$$\tau_{e,l} = \frac{L^2}{(l + \frac{1}{2})^2 \pi^2 D_a}. \quad (6.24)$$

For the first term ( $m = 1$ ) in the sine series, we have

$$-a_{o,1}e^{-\frac{t}{\tau_{o,1}}}\frac{1}{\tau_{o,1}}\sin\frac{\pi x}{L} = -D_a a_{o,1}e^{-\frac{t}{\tau_{o,1}}}\left(\frac{\pi}{L}\right)^2\sin\frac{\pi x}{L} \quad (6.25)$$

and hence

$$\tau_{o,1} = \frac{L^2}{\pi^2 D_a}. \quad (6.26)$$

For the other modes  $m$  we similarly obtain

$$\tau_{o,m} = \frac{L^2}{m^2 \pi^2 D_a}. \quad (6.27)$$

The lifetime of various modes  $\tau_{e,l}$  and  $\tau_{o,m}$  shows that the higher modes are attenuated faster. Therefore, if there is a wild density profile at the beginning (in time  $t = 0$ ), then the profile is simplified with time. A set of snapshots of the evolution of an arbitrary density profile is given in Fig. 6.3.

### Other physics-motivated examples for diffusion

The question is whether it is possible to achieve stationarity in the presence of diffusion. It is not if the continuity equation holds. It is possible in the presence of the source term,

$$\nabla \cdot n\mathbf{u} = Q. \quad (6.28)$$

In case of collisional ionisation  $Q = Zn$ , where  $Z$  is the ionisation function, we have

$$\nabla \cdot n\mathbf{u} = -D_a \Delta n = Zn \quad \rightarrow \quad \Delta n = -\frac{Z}{D_a}n, \quad (6.29)$$

solution to which is again the series of harmonic functions.

The recombination, on the other hand, takes a different form because it does change the number of charged particles. The number density of the ionised particles naturally decays with  $n^2$  in recombination, hence the continuity equation may be written in the form

$$\frac{\partial n}{\partial t} - D_a \nabla^2 n = -\alpha n^2. \quad (6.30)$$

Similarly to the previous cases, we consider the solution separated to the temporal and spatial parts  $n = TS$ . Further, we will only be interested in the temporal part, which is equivalent to modification of (6.30) to

$$\frac{\partial n}{\partial t} = -\alpha n^2, \quad (6.31)$$

where  $n^2$  dependence is a consequence of the recombination being proportional to both the density of electrons and ions and the use of the plasma approximation. Then

$$\frac{1}{n^2} dn = -\alpha dt \quad \rightarrow \quad -\frac{1}{n} = -\frac{1}{n_0} - \alpha t. \quad (6.32)$$

For  $t \rightarrow \infty$  we have  $n \sim \frac{1}{\alpha t}$ , which says that due to the recombination in plasmas, the density of the charged particles drops as  $1/t$ . The spatial part of the continuity equation is not relevant to the problem; thus, we neglect it from the solution.

This solution is important for understanding which process drives the diffusion of charged particles in a cloud of plasma left-alone in time. Obviously, in the initial moment, recombination will be the dominant process, as its time dependence is steeper for small  $t$ . On the contrary, when the density of the plasma gets smaller, the diffusion with exponential time dependence prevails and further drives the diffusion.

## 6.2 Diffusion in highly ionised plasmas

We will show that in highly ionised plasmas, diffusion resembles the process of recombination rather than ambipolar diffusion. Let us use the fluid approximation to deal with the problem.

$$\rho \frac{\partial \mathbf{u}}{\partial t} = \mathbf{j} \times \mathbf{B} - \nabla p \quad (6.33)$$

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} = \eta \mathbf{j} \quad (6.34)$$

Note that generally, the specific resistivity of plasma  $\eta$  is not a scalar value but rather a tensor and thus may depend on direction. In plasma, this usually is the case. Therefore, we split the specific resistivity into the parallel and perpendicular components ( $\eta = \eta_{\parallel} + \eta_{\perp}$ ), where in the parallel direction Ohm's law simply is

$$\mathbf{E}_{\parallel} = \eta_{\parallel} \mathbf{j}_{\parallel}. \quad (6.35)$$

To obtain a similar relation in the direction perpendicular to the magnetic field, when assuming the stationarity, by vector-multiplying (6.34) by  $\mathbf{B}$  and using (6.33) we have

$$\mathbf{E} \times \mathbf{B} + (\mathbf{u} \times \mathbf{B}) \times \mathbf{B} = \mathbf{E} \times \mathbf{B} - B^2 \mathbf{u} + B^2 \mathbf{u}_{\parallel} = \eta_{\perp} \nabla p. \quad (6.36)$$

Since  $\mathbf{u} = \mathbf{u}_{\parallel} + \mathbf{u}_{\perp}$ , then the last two terms on the left-hand side together give  $B^2 \mathbf{u}_{\perp}$  and thus

$$\mathbf{u}_{\perp} = \frac{\mathbf{E} \times \mathbf{B}}{B^2} - \frac{\eta_{\perp} \nabla p}{B^2}. \quad (6.37)$$

The first term on the right-hand side is the same for both (electron and ion) fluids. The second one, on the other hand, may drive the diffusion. The perpendicular particle flux on top of the E-B drift is

$$\mathbf{\Gamma}_\perp = n\mathbf{u}'_\perp = -\frac{\eta_\perp n(K_B T_e + K_B T_i)}{B^2} \nabla n = -D_\perp \nabla n. \quad (6.38)$$

Note that the diffusion coefficient  $D_\perp$  is a function of particle density and for simpler manipulation may be rewritten as

$$D_\perp = \frac{\eta_\perp n K_B \sum_\alpha T_\alpha}{B^2} = 2An. \quad (6.39)$$

The factor of two is introduced so that  $\nabla \cdot (2An \nabla n) = 2A \nabla \cdot (n \nabla n) = A \Delta n^2$ .

The continuity equation

$$\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{u}) = 0 \quad (6.40)$$

then transforms to

$$\frac{\partial n}{\partial t} = A \Delta n^2. \quad (6.41)$$

Such an equation may be solved by separation of variables, hence  $n(\mathbf{r}, t) = T(t)S(\mathbf{r})$ . Then  $S \frac{dT}{dt} = AT^2 \Delta S^2$  and

$$\frac{1}{T^2} \frac{dT}{dt} = \frac{A}{S} \Delta S^2 = \text{const} = -\frac{1}{\tau^2} \quad (6.42)$$

and in the case we solved previously. Then the temporal part has the solution

$$\frac{1}{T} = \frac{t}{\tau^2} + \frac{1}{T_0} \quad (6.43)$$

and the form of  $\tau$  depends on the spatial part  $S$ . For  $t \rightarrow \infty$

$$T \propto \frac{\tau^2}{t}, \quad (6.44)$$

which reminds us of the the solution of the change of the particle density caused by recombinations. The difference between the ambipolar diffusion and the result we just obtained is in the diffusion coefficient depending on particle density. Note that the result we obtained is not consistent with the assumptions of the derivation of Fick's law, as we allow for non-stationarity here.

### 6.3 Diffusion in a magnetic field

In presence of a magnetic field, the Euler equation retains an additional term. Let us see what happens in the presence of a magnetic field. We must go back to the double-fluid approximation since collisions are essential in the process of diffusion.

$$mn \left[ \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right] = \text{sgn } q en(\mathbf{E} + \mathbf{u} \times \mathbf{B}) - K_B T \nabla n - mn \nu \mathbf{u} \quad (6.45)$$

Let's assume that the collisions are much faster than the overall evolution of the system in time, or in other words  $\frac{d}{dt} \ll \omega_c, \nu$ , where  $\nu$  is the collision frequency with neutrals and  $\tau \equiv 1/\nu$  is the corresponding time between collisions. Then the left-hand side may be neglected. Thus in components of the Cartesian system with the  $z$  axis oriented along the magnetic field ( $\mathbf{B} = B\mathbf{e}_z$ ) we have

$$mn\nu u_x = \text{sgn } q enE_x - K_B T \frac{\partial n}{\partial x} + \text{sgn } q enu_y B, \quad (6.46)$$

$$mn\nu u_y = \text{sgn } q enE_y - K_B T \frac{\partial n}{\partial y} - \text{sgn } q enu_x B. \quad (6.47)$$

Along the magnetic field, a classic diffusion as derived above occurs.

By introducing  $\mu = \frac{e}{m\nu}$ ,  $D = \frac{K_B T}{m\nu} = \frac{\omega_c}{\nu} \frac{K_B T}{eB}$  we have

$$u_x = \text{sgn } q \mu E_x - \frac{D}{n} \frac{\partial n}{\partial x} + \text{sgn } q \frac{\omega_c}{\nu} u_y, \quad (6.48)$$

$$u_y = \text{sgn } q \mu E_y - \frac{D}{n} \frac{\partial n}{\partial y} - \text{sgn } q \frac{\omega_c}{\nu} u_x, \quad (6.49)$$

and (using  $\tau \equiv 1/\nu$ )

$$u_x(1 + \omega_c^2 \tau^2) = \text{sgn } q \mu E_x - \frac{D}{n} \frac{\partial n}{\partial x} + \omega_c^2 \tau^2 \frac{E_y}{B} - \text{sgn } q \omega_c^2 \tau^2 \frac{K_B T}{neB} \frac{\partial n}{\partial y}, \quad (6.50)$$

$$u_y(1 + \omega_c^2 \tau^2) = \text{sgn } q \mu E_y - \frac{D}{n} \frac{\partial n}{\partial y} - \omega_c^2 \tau^2 \frac{E_x}{B} + \text{sgn } q \omega_c^2 \tau^2 \frac{K_B T}{neB} \frac{\partial n}{\partial x}. \quad (6.51)$$

The last two terms on the right-hand side obviously indicate the already known E-B and diamagnetic drifts. Hence

$$\mathbf{u}_\perp = \frac{\mathbf{v}_E + \mathbf{v}_D}{1 + 1/(\omega_c^2 \tau^2)} + \text{sgn } q \frac{\mu}{1 + \omega_c^2 \tau^2} \mathbf{E} - \frac{D}{1 + \omega_c^2 \tau^2} \frac{\nabla n}{n}, \quad (6.52)$$

where  $\mathbf{v}_E = \mathbf{E} \times \mathbf{B}/B^2$  and  $\mathbf{v}_D = -\nabla p \times \mathbf{B}/(qnB^2)$ . Let us further identify

$$\mu_\perp \equiv \frac{\mu}{1 + \omega_c^2 \tau^2} \quad \text{and} \quad D_\perp \equiv \frac{D}{1 + \omega_c^2 \tau^2}. \quad (6.53)$$

Hence generalised Fick's law transforms into

$$\mathbf{u}_\perp = \text{sgn } q \mu_\perp \mathbf{E} - D_\perp \frac{\nabla n}{n} + \frac{\mathbf{v}_E + \mathbf{v}_D}{1 + 1/(\omega_c^2 \tau^2)}. \quad (6.54)$$

The last term represents the already familiar drifts in the plasma fluid, only with the correction term for the diffusion. Obviously, the correcting term is important for  $\nu > \omega_c$ . Hence in the weak field regime, drifts are suppressed by collisions.

The correction term for the mobility and the diffusion coefficient act in opposite ways. In the strong field, i.e., when  $\nu < \omega_c$ , diffusion is suppressed. In particular, we have

$$D_\perp \propto \frac{K_B T \nu}{m \omega_c^2} \quad \text{and} \quad D_\parallel = \frac{K_B T}{m \nu}, \quad (6.55)$$

where  $D_{\parallel}$  is the non-modified diffusion coefficient which appears in the component of the Euler equation parallel to the magnetic field. Hence the decay caused by diffusion in plasmas has different rates in different directions. In case the collisional frequency is proportional<sup>1</sup> to  $m^{-1/2}$  then the diffusion depends also on the species of the particles. Then

$$D_{\perp} \propto m^{1/2} \quad \text{and} \quad D_{\parallel} \propto m^{-1/2}. \quad (6.56)$$

In the perpendicular direction, the diffusion is faster for ions than electrons, while in the parallel direction the electrons diffuse faster. It has to be noted that generally speaking, collisions are important for particles to be able to penetrate across the magnetic field.

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<sup>1</sup>As we discussed at the beginning of this chapter,  $\nu = n_n \langle \sigma v \rangle$ . Leaving out a possible dependence of the effective cross-section on velocity, we have  $\nu \propto \langle v \rangle \propto \sqrt{\langle v^2 \rangle}$ . Since  $T \propto m \langle v^2 \rangle$ , for a fixed temperature  $T$  we have  $\langle v^2 \rangle \propto m^{-1}$  and finally  $\nu \propto m^{-1/2}$ .

## Chapter 7

# Equilibrium and stability

### 7.1 Equilibrium

$$\rho \frac{\partial \mathbf{u}}{\partial t} = \mathbf{j} \times \mathbf{B} - \nabla p. \quad (7.1)$$

$$\nabla p = \mathbf{j} \times \mathbf{B} \quad (7.2)$$

$$\mathbf{B} \times \nabla p = \mathbf{B} \times (\mathbf{j} \times \mathbf{B}) = B^2 \mathbf{j} - \mathbf{B} \cdot \mathbf{j} \mathbf{B} = B^2 (\mathbf{j} - \mathbf{j}_{\parallel}). \quad (7.3)$$

$$\mathbf{j}_{\perp} = \frac{\mathbf{B} \times \nabla p}{B^2} = (K_B T_i + K_B T_e) \frac{\mathbf{B} \times \nabla n}{B^2}, \quad (7.4)$$

### 7.2 Plasma $\beta$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j} \quad (7.5)$$

$$\begin{aligned} \nabla p = \mathbf{j} \times \mathbf{B} &= \frac{1}{\mu_0} (\nabla \times \mathbf{B}) \times \mathbf{B} = \\ &= \frac{1}{\mu_0} [-(\nabla \mathbf{B}) \cdot \mathbf{B} + (\mathbf{B} \cdot \nabla) \mathbf{B}] = \frac{1}{\mu_0} \left[ (\mathbf{B} \cdot \nabla) \mathbf{B} - \frac{1}{2} \nabla B^2 \right] \end{aligned} \quad (7.6)$$

$$\nabla \left( p + \frac{1}{\mu_0} B^2/2 \right) = \frac{1}{\mu_0} (\mathbf{B} \cdot \nabla) \mathbf{B}. \quad (7.7)$$

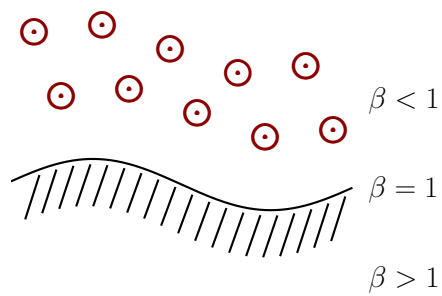
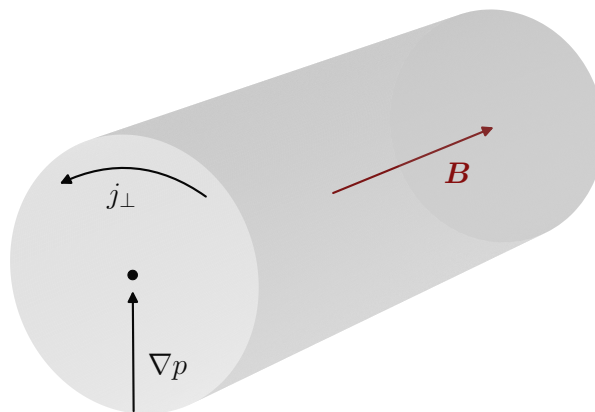
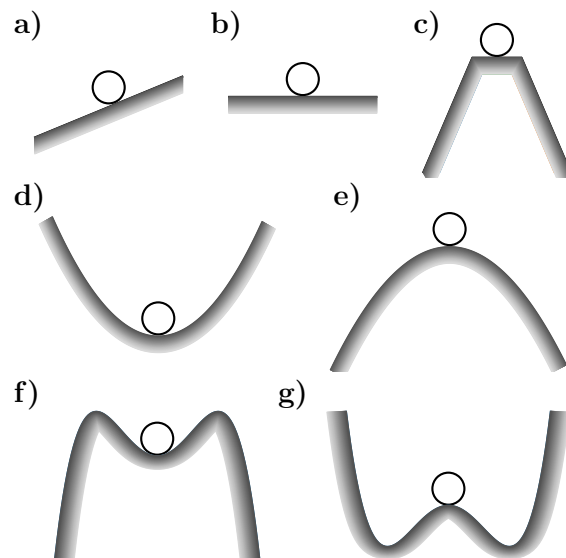
$$p + \frac{B^2}{2\mu_0} = \text{const.} \quad (7.8)$$

$$\beta \equiv \frac{p}{p_{\text{mag}}} = \frac{n \sum_{\alpha} K_B T_{\alpha}}{\frac{B^2}{2\mu_0}}, \quad (7.9)$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad (7.10)$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j} \quad (7.11)$$

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} = \eta \mathbf{j} \quad (7.12)$$



$$\mathbf{u} = \frac{1}{\rho}(n_i M \mathbf{u}_i + n_e m \mathbf{u}_e) \sim \frac{M \mathbf{u}_i + m \mathbf{u}_e}{M + m}, \quad (7.13)$$

$$\mathbf{j} = e(n_i \mathbf{u}_i - n_e \mathbf{u}_e) \sim ne(\mathbf{u}_i - \mathbf{u}_e). \quad (7.14)$$

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \eta \mathbf{j} = -\frac{\eta}{\mu_0} [\nabla(\nabla \cdot \mathbf{B}) - \Delta \mathbf{B}] = \frac{\eta}{\mu_0} \Delta \mathbf{B}, \quad (7.15)$$

$$\frac{\partial \mathbf{B}}{\partial t} = \frac{\eta}{\mu_0 L^2} \mathbf{B}, \quad (7.16)$$

$$\mathbf{B} = \mathbf{B}_0 \exp \left[ \pm \frac{t}{\tau} \right], \quad (7.17)$$

$$\tau = \frac{\mu_0 L^2}{\eta} \quad (7.18)$$

$$P = \eta j^2. \quad (7.19)$$

$$\Delta W = \eta j^2 \tau. \quad (7.20)$$

$$\Delta W = \eta j^2 \tau = \eta \left( \frac{B}{\mu_0 L} \right)^2 \frac{\mu_0 L^2}{\eta} = \frac{B^2}{\mu_0}. \quad (7.21)$$

### 7.3 Instabilities

$$M n_0 \frac{\partial \mathbf{u}_{i1}}{\partial t} = e n_0 \mathbf{E}_1, \quad (7.22)$$

$$m n_0 \left[ \frac{\partial \mathbf{u}_{e1}}{\partial t} + (\mathbf{u}_0 \cdot \nabla) \mathbf{u}_{e1} \right] = -e n_0 \mathbf{E}_1. \quad (7.23)$$

$$-i\omega M u_{i1} = e E_1, \quad (7.24)$$

$$m(-i\omega + i k u_0) u_{e1} = -e E_1. \quad (7.25)$$

$$\frac{\partial n_{i1}}{\partial t} + n_0 \nabla \cdot \mathbf{u}_{i1} = 0, \quad (7.26)$$

$$n_{i1} = \frac{k}{\omega} n_0 u_{i1} = \frac{i e n_0 k}{M \omega^2} E_1. \quad (7.27)$$

$$\frac{\partial n_{e1}}{\partial t} + n_0 \nabla \cdot \mathbf{u}_{e1} + (\mathbf{u}_0 \cdot \nabla) n_{e1} = 0, \quad (7.28)$$

$$(-i\omega + i k u_0) n_{e1} + i k n_0 u_{e1} = 0, \quad (7.29)$$

$$n_{e1} = \frac{k n_0 u_{e1}}{\omega - k u_0} = -\frac{i e k n_0}{m(\omega - k u_0)^2} E_1. \quad (7.30)$$

$$\nabla \cdot \mathbf{E}_1 = \frac{e}{\varepsilon_0} (n_{i1} - n_{e1}), \quad (7.31)$$

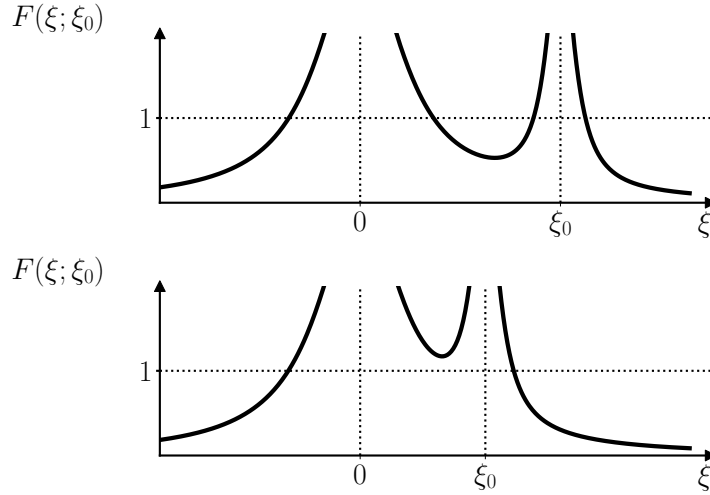
$$i k E_1 = \frac{e}{\varepsilon_0} i e k n_0 \left[ \frac{1}{M \omega^2} + \frac{1}{m(\omega - k u_0)^2} \right] E_1 \quad (7.32)$$

$$1 = \frac{e^2 n_0}{\varepsilon_0 m} \left[ \frac{m/M}{\omega^2} + \frac{1}{(\omega - k u_0)^2} \right] = \omega_p^2 \left[ \frac{m/M}{\omega^2} + \frac{1}{(\omega - k u_0)^2} \right]. \quad (7.33)$$

$$\omega_j = \alpha_j + i\gamma_j, \quad \text{where} \quad \alpha_j = \Re[\omega_j] \quad \text{and} \quad \gamma_j = \Im[\omega_j]. \quad (7.34)$$

$$E_1 = E e_x \exp[i(kx - \omega t)] e^{\gamma_j t}. \quad (7.35)$$

$$1 = \frac{m/M}{\xi^2} + \frac{1}{(\xi - \xi_0)^2} = F(\xi; \xi_0). \quad (7.36)$$



$$M n_{i0} [(\mathbf{u}_{i0} \cdot \nabla) \mathbf{u}_{i0}] = e n_{i0} \mathbf{u}_{i0} \times \mathbf{B}_0 + M n_{i0} \mathbf{g}. \quad (7.37)$$

$$0 = e n_{i0} \mathbf{u}_{i0} \times \mathbf{B}_0 + M n_{i0} \mathbf{g} \quad (7.38)$$

$$0 = e \mathbf{B}_0 \times (\mathbf{u}_{i0} \times \mathbf{B}_0) + M \mathbf{B}_0 \times \mathbf{g} = e B_0^2 \mathbf{u}_{i0} - e \mathbf{B}_0 \cdot \mathbf{u}_{i0} \mathbf{B}_0 + M \mathbf{B}_0 \times \mathbf{g}, \quad (7.39)$$

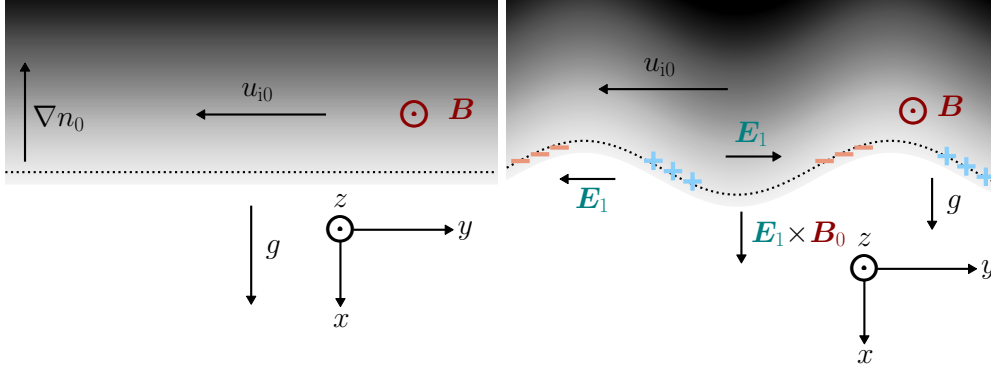
$$\mathbf{u}_{i0} = -\frac{M}{e} \frac{\mathbf{B}_0 \times \mathbf{g}}{B_0^2} = -\frac{g}{\Omega_c} \mathbf{e}_y, \quad (7.40)$$

$$M(n_{i0} + n_{i1}) \left\{ \frac{\partial}{\partial t} (\mathbf{u}_{i0} + \mathbf{u}_{i1}) + [(\mathbf{u}_{i0} + \mathbf{u}_{i1}) \cdot \nabla] (\mathbf{u}_{i0} + \mathbf{u}_{i1}) \right\} = e(n_{i0} + n_{i1}) [\mathbf{E}_1 + (\mathbf{u}_{i0} + \mathbf{u}_{i1}) \times \mathbf{B}_0] + M(n_{i0} + n_{i1}) \mathbf{g} \quad (7.41)$$

$$M(n_{i0} + n_{i1}) \left\{ \frac{\partial}{\partial t} (\mathbf{u}_{i0} + \mathbf{u}_{i1}) + (\mathbf{u}_{i0} \cdot \nabla) \mathbf{u}_{i0} + (\mathbf{u}_{i0} \cdot \nabla) \mathbf{u}_{i1} + (\mathbf{u}_{i1} \cdot \nabla) \mathbf{u}_{i0} \right\} = e(n_{i0} + n_{i1}) [\mathbf{E}_1 + (\mathbf{u}_{i0} + \mathbf{u}_{i1}) \times \mathbf{B}_0] + M(n_{i0} + n_{i1}) \mathbf{g} \quad (7.42)$$

$$M(n_{i0} + n_{i1}) (\mathbf{u}_{i0} \cdot \nabla) \mathbf{u}_{i0} = (n_{i0} + n_{i1}) e \mathbf{u}_{i0} \times \mathbf{B}_0 + M(n_{i0} + n_{i1}) \mathbf{g}. \quad (7.43)$$

$$M n_{i0} \left[ \frac{\partial \mathbf{u}_{i1}}{\partial t} + (\mathbf{u}_{i0} \cdot \nabla) \mathbf{u}_{i1} \right] = e n_{i0} (\mathbf{E}_1 + \mathbf{u}_{i1} \times \mathbf{B}_0). \quad (7.44)$$



$$M(\omega - \mathbf{k} \cdot \mathbf{u}_{i0})\mathbf{u}_{i1} = ie(\mathbf{E}_1 + \mathbf{u}_{i1} \times \mathbf{B}_0). \quad (7.45)$$

$$u_{i1x} = \frac{ie}{M(\omega - ku_{i0})}u_{i1y}B_0, \quad (7.46)$$

$$u_{i1y} = \frac{-ie}{M(\omega - ku_{i0})}(E_{1y} - u_{i1x}B_0). \quad (7.47)$$

$$u_{i1y} = \frac{-ie}{M(\omega - ku_{i0})} \left( E_{1y} + \frac{-ieB_0^2}{M(\omega - ku_{i0})}u_{i1y} \right) \quad (7.48)$$

$$u_{i1y} = \frac{-ieE_{1y}}{M(\omega - ku_{i0})} \left[ 1 - \frac{e^2B_0^2/M^2}{(\omega - ku_{i0})^2} \right]^{-1} = \frac{-ieE_{1y}}{M(\omega - ku_{i0})} \left[ 1 - \frac{\Omega_c^2}{(\omega - ku_{i0})^2} \right]^{-1}. \quad (7.49)$$

$$u_{i1x} = -\frac{e\Omega_c E_{1y}}{M(\omega - ku_{i0})^2} \left[ 1 - \frac{\Omega_c^2}{(\omega - ku_{i0})^2} \right]^{-1}. \quad (7.50)$$

$$u_{i1x} = \frac{E_{1y}}{B_0}, \quad u_{i1y} = -i\frac{\omega - ku_{i0}}{\Omega_c} \frac{E_{1y}}{B_0} \quad (7.51)$$

$$u_{e1x} = \frac{E_{1y}}{B_0}, \quad u_{e1y} = -i\frac{\omega - ku_{e0}}{\omega_c} \frac{E_{1y}}{B_0}, \quad (7.52)$$

$$\frac{\partial n_{i1}}{\partial t} + \nabla \cdot (n_{i0}\mathbf{u}_{i0}) + (\mathbf{u}_{i0} \cdot \nabla)n_{i1} + n_{i1} \nabla \cdot \mathbf{u}_{i0} + (\mathbf{u}_{i1} \cdot \nabla)n_{i0} + n_{i0} \nabla \cdot \mathbf{u}_{i1} + \nabla \cdot (n_{i1}\mathbf{u}_{i1}) = 0. \quad (7.53)$$

$$-i\omega n_{i1} + iku_{i0}n_{i1} + u_{i1x} \frac{\partial n_{i0}}{\partial x} + ikn_{i0}u_{i1y} = 0. \quad (7.54)$$

$$-i\omega n_1 + u_{e1x} \frac{\partial n_0}{\partial x} = 0. \quad (7.55)$$

$$-i\omega n_1 + iku_{i0}n_1 + \frac{E_{1y}}{B_0} \frac{\partial n_0}{\partial x} - i^2kn_0 \frac{\omega - ku_{i0}}{\Omega_c} \frac{E_{1y}}{B_0} = 0, \quad (7.56)$$

$$(\omega - ku_{i0})n_1 + i\frac{E_{1y}}{B_0} \frac{\partial n_0}{\partial x} + ikn_0 \frac{\omega - ku_{i0}}{\Omega_c} \frac{E_{1y}}{B_0} = 0. \quad (7.57)$$

$$-i\omega n_1 + \frac{E_{1y}}{B_0} \frac{\partial n_0}{\partial x} = 0, \quad (7.58)$$

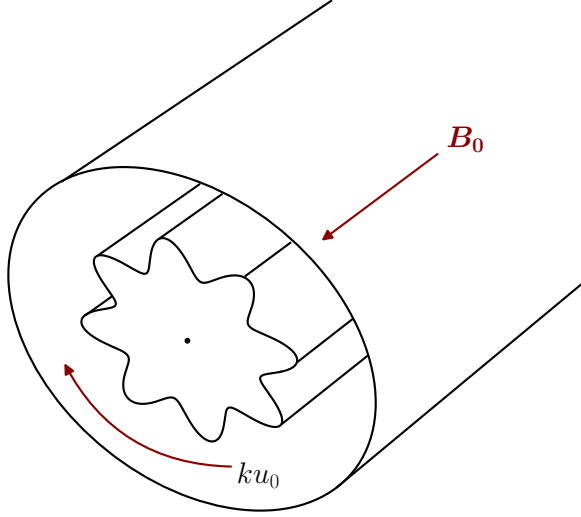


Figure 7.1: A flute instability.

$$(\omega - ku_{i0})n_1 - \frac{\omega n_1}{\frac{\partial n_0}{\partial x}} \frac{\partial n_0}{\partial x} - kn_0 \frac{\omega - ku_{i0}}{\Omega_c} \frac{\omega n_1}{\frac{\partial n_0}{\partial x}} = 0. \quad (7.59)$$

$$\omega(\omega - ku_{i0}) = -\frac{u_{i0}\Omega_c}{n_0} \frac{\partial n_0}{\partial x}. \quad (7.60)$$

$$\omega(\omega - ku_{i0}) = \frac{g}{n_0} \frac{\partial n_0}{\partial x}, \quad (7.61)$$

$$\omega^2 - ku_{i0}\omega - g \left( \frac{\frac{\partial n_0}{\partial x}}{n_0} \right) = 0. \quad (7.62)$$

$$\omega = \frac{1}{2}ku_{i0} \pm \left[ \frac{1}{4}k^2u_{i0}^2 + g \left( \frac{\frac{\partial n_0}{\partial x}}{n_0} \right) \right]^{\frac{1}{2}}. \quad (7.63)$$

$$-g \left( \frac{\frac{\partial n_0}{\partial x}}{n_0} \right) > \frac{1}{4}k^2u_{i0}^2, \quad (7.64)$$

$$\mathbf{u}_{i0} = \mathbf{u}_{Di} = \frac{K_B T_i}{eB_0} \frac{1}{n_0} \frac{\partial n_0}{\partial x} \mathbf{e}_y, \quad (7.65)$$

$$\mathbf{u}_{e0} = \mathbf{u}_{De} = -\frac{K_B T_e}{eB_0} \frac{1}{n_0} \frac{\partial n_0}{\partial x} \mathbf{e}_y, . \quad (7.66)$$

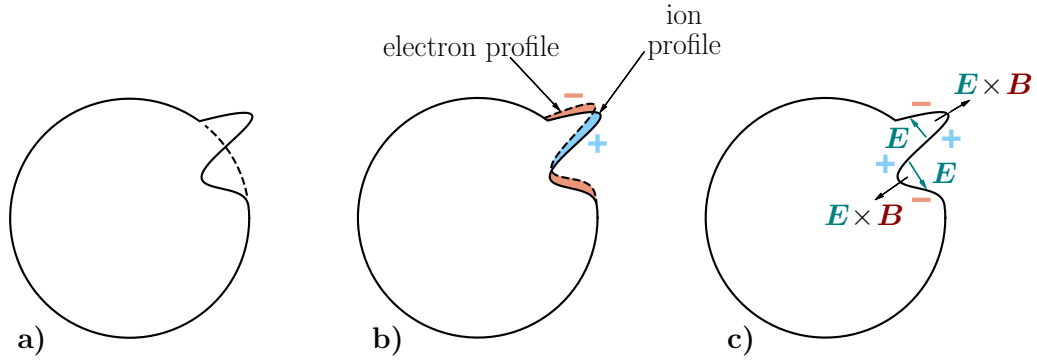
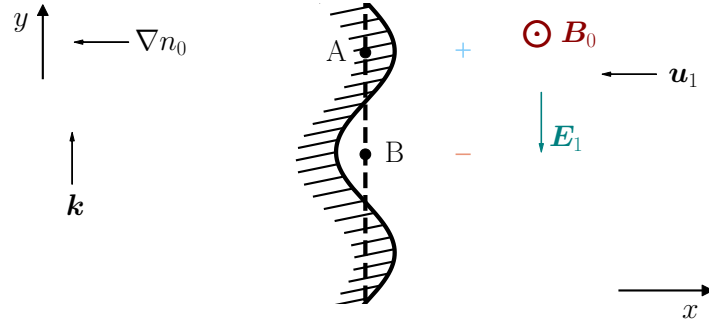


Figure 7.2: Development of the flute instability. a) Initial disturbance. b) Effect of ion and electron azimuthal drifts. c) Resulting  $\mathbf{E} \times \mathbf{B}$  drifts increase the amplitude.



$$\frac{n_{e1}}{n_0} = \frac{e\phi_1}{K_B T_e}. \quad (7.67)$$

$$\mathbf{u}_1 = \frac{\mathbf{E}_1 \times \mathbf{B}_0}{B_0^2}. \quad (7.68)$$

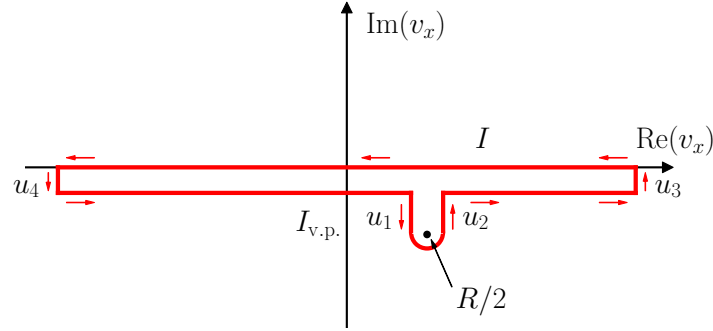
$$u_{1x} = \frac{E_y}{B_0} = -\frac{ik_y \phi_1}{B_0}, \quad (7.69)$$

$$\frac{\partial n_1}{\partial t} = -u_{1x} \frac{\partial n_0}{\partial x}, \quad (7.70)$$

$$-i\omega n_{e1} = -u_{e1x} \frac{\partial n_0}{\partial x} = \frac{ik_y \phi_1}{B_0} \frac{\partial n_0}{\partial x} = -i\omega n_0 \frac{e\phi_1}{K_B T_e}, \quad (7.71)$$

$$\frac{\omega}{k} = -\frac{K_B T_e}{eB_0} \frac{1}{n_0} \frac{\partial n_0}{\partial x}, \quad (7.72)$$

$$\frac{\partial f}{\partial t} + \frac{\partial h}{\partial p_i} \frac{\partial f}{\partial r_i} - \frac{\partial h}{\partial r_i} \frac{\partial f}{\partial p_i} = \left( \frac{\partial f}{\partial t} \right)_{\text{coll}}, \quad (7.73)$$



$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla f + \frac{F}{m} \cdot \frac{\partial f}{\partial \mathbf{v}} = \left( \frac{\partial f}{\partial t} \right)_{\text{coll}}, \quad (7.74)$$

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla f + \frac{\langle F \rangle}{m} \cdot \frac{\partial f}{\partial \mathbf{v}} = \left( \frac{\partial f}{\partial t} \right)'_{\text{coll}}, \quad (7.75)$$

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla f + \frac{q}{m} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f}{\partial \mathbf{v}} = 0, \quad (7.76)$$

$$-i\omega f_1 + ik f_1 v_x = \frac{e}{m} E_{1x} \frac{\partial f_0}{\partial v_x}, \quad (7.77)$$

$$f_1 = \frac{ie E_{1x}}{m} \frac{\partial f_0 / \partial v_x}{\omega - k v_x}. \quad (7.78)$$

$$\nabla \cdot \mathbf{E}_1 = ik E_{1x} = -\frac{en_1}{\varepsilon_0} = -\frac{e}{\varepsilon_0} \int_{-\infty}^{+\infty} f_1 d^3 \mathbf{v}. \quad (7.79)$$

$$\int_{-\infty}^{+\infty} f_1 d^3 \mathbf{v} = \frac{ie E_{1x}}{m} n_0 \int_{-\infty}^{+\infty} \frac{\partial \hat{f}_0 / \partial v_x}{\omega - k v_x} d^3 \mathbf{v} = -\frac{ik E_{1x} \varepsilon_0}{e}, \quad (7.80)$$

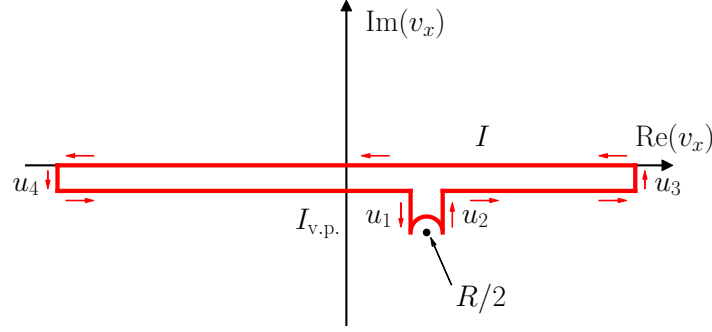
$$1 = -\frac{e^2 n_0}{\varepsilon_0 m k} \int_{-\infty}^{+\infty} \frac{\partial \hat{f}_0 / \partial v_x}{\omega - k v_x} d^3 \mathbf{v} = +\frac{\omega_p^2}{k^2} \int_{-\infty}^{+\infty} \frac{\partial \hat{f}_0 / \partial v_x}{v_x - \frac{\omega}{k}} d^3 \mathbf{v}. \quad (7.81)$$

$$\hat{f}_0(\mathbf{v}) = \tilde{f}_0(v_x) \tilde{f}_0(v_y) \tilde{f}_0(v_z), \quad (7.82)$$

$$1 = \frac{\omega_p^2}{k^2} \int_{-\infty}^{+\infty} dv_x \frac{\partial \tilde{f}_0(v_x) / \partial v_x}{v_x - \frac{\omega}{k}} \int_{-\infty}^{+\infty} dv_y \tilde{f}_0(v_y) \int_{-\infty}^{+\infty} dv_z \tilde{f}_0(v_z) = \frac{\omega_p^2}{k^2} \int_{-\infty}^{+\infty} \frac{\partial \tilde{f}_0 / \partial v_x}{v_x - \frac{\omega}{k}} dv_x. \quad (7.83)$$

$$-I + I_{\text{v.p.}} + u_1 + u_2 + \frac{R}{2} + u_3 + u_4 = 0, \quad (7.84)$$

$$1 = \frac{\omega_p^2}{k^2} \left\{ \text{v.p.} \int_{-\infty}^{+\infty} \frac{\partial \tilde{f}_0 / \partial v_x}{v_x - \frac{\omega}{k}} dv_x + \left[ i\pi \frac{\partial \tilde{f}_0}{\partial v_x} \right]_{\frac{\omega}{k}} \right\}. \quad (7.85)$$



$$\int_{-\infty}^{+\infty} \frac{\partial \tilde{f}_0}{\partial v_x} \frac{dv_x}{v_x - \omega/k} = \left[ \frac{\tilde{f}_0}{v_x - \omega/k} \right]_{+\infty}^{-\infty} + \int_{-\infty}^{+\infty} \frac{\tilde{f}_0 dv_x}{(v_x - \omega/k)^2}. \quad (7.86)$$

$$(v_x - v_\varphi)^{-2} \sim v_\varphi^{-2} \left( 1 + \frac{2v_x}{v_\varphi} + \frac{3v_x^2}{v_\varphi^2} \right). \quad (7.87)$$

$$\begin{aligned} \int_{-\infty}^{+\infty} \frac{\tilde{f}_0 dv_x}{(v_x - v_\varphi)^2} &\sim v_\varphi^{-2} \left[ \int_{-\infty}^{+\infty} \tilde{f}_0 dv_x + \int_{-\infty}^{+\infty} \frac{2v_x}{v_\varphi} \tilde{f}_0 dv_x + \int_{-\infty}^{+\infty} \frac{3v_x^2}{v_\varphi^2} \tilde{f}_0 dv_x \right] = \\ &= v_\varphi^{-2} \left[ 1 + \frac{2}{v_\varphi} \langle v_x \rangle + \frac{3 \langle v_x^2 \rangle}{v_\varphi^2} \right] = v_\varphi^{-2} \left( 1 + \frac{3 \langle v_x^2 \rangle}{v_\varphi^2} \right), \end{aligned} \quad (7.88)$$

$$1 = \frac{\omega_p^2}{k^2} \frac{k^2}{\omega^2} \left\{ 1 + \frac{3 \langle v_x^2 \rangle}{v_\varphi^2} + i\pi \left[ \frac{\omega^2}{k^2} \frac{\partial \tilde{f}_0}{\partial v_x} \right]_{\frac{\omega}{k}} \right\}. \quad (7.89)$$

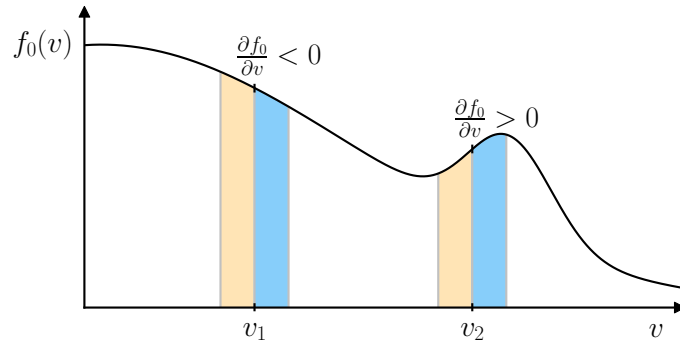
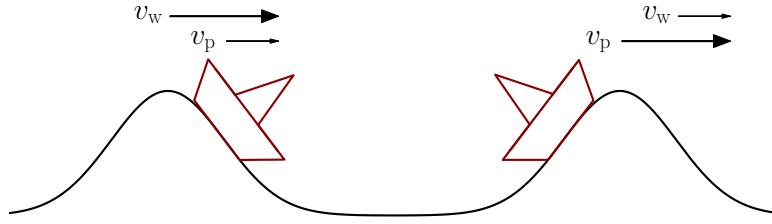
$$1 = \frac{\omega_p^2}{\omega^2} \left( 1 + 3 \frac{k^2}{\omega^2} \frac{K_B T_e}{m} \right). \quad (7.90)$$

$$\omega^2 = \omega_p^2 + 3k^2 \frac{K_B T_e}{m}, \quad (7.91)$$

$$\omega^2 = \omega_p^2 \left\{ 1 + i\pi \left[ \frac{\omega^2}{k^2} \frac{\partial \tilde{f}_0}{\partial v_x} \right]_{\frac{\omega}{k}} \right\}. \quad (7.92)$$

$$\omega \sim \omega_p \left\{ 1 + \frac{1}{2} i\pi \left[ \frac{\omega^2}{k^2} \frac{\partial \tilde{f}_0}{\partial v_x} \right]_{\frac{\omega}{k}} \right\}. \quad (7.93)$$

$$\frac{\omega^2}{\omega_p^2} = \left\{ 1 + i\pi \left[ \frac{\omega^2}{k^2} \frac{\partial \tilde{f}_0}{\partial v_x} \right]_{\frac{\omega}{k}} \right\} \sim 1 \quad (7.94)$$



$$\omega \sim \omega_p \left\{ 1 + \frac{1}{2} i \pi \left[ \frac{\omega_p^2}{k^2} \frac{\partial \tilde{f}_0}{\partial v_x} \right]_{\frac{\omega}{k}} \right\}. \quad (7.95)$$

$$\Im \left( \frac{\omega}{\omega_p} \right) \sim 0.22 \sqrt{\pi} \left( \frac{\omega_p}{k v_T} \right)^3 \exp \left[ -\frac{1}{2 k^2 \lambda_D} \right], \quad (7.96)$$