

The Galactic Center

Jason Dexter
MPE Garching

with slides from R. Genzel, S. Gillessen, and the MPE GC group

mpe.mpg.de/ir/GC

The Galactic Center

1. Evidence for a massive black hole (14.11)
2. A paradox of youth (16.11)
3. Sgr A* and the faintest black holes (today)
4. Outbursts from Sgr A* and the high energy GC (23.11)

Seminar: strong gravity around Sgr A* (23.11)

About the lectures

- Selected topics: central parsec, highly biased
- Please ask questions!
- ~1 interactive Q / lecture:
~10 mins to think/calculate, discuss, share
- Further reading: Genzel+2010, Morris+2012, Falcke & Markoff 2013

About the lectures

- pdf of slides online:
mpe.mpg.de/~jdexter/GCslides
- Requests for topics: now or e-mail
jdexter@mpe.mpg.de
- Any other Q's: around after, Wed

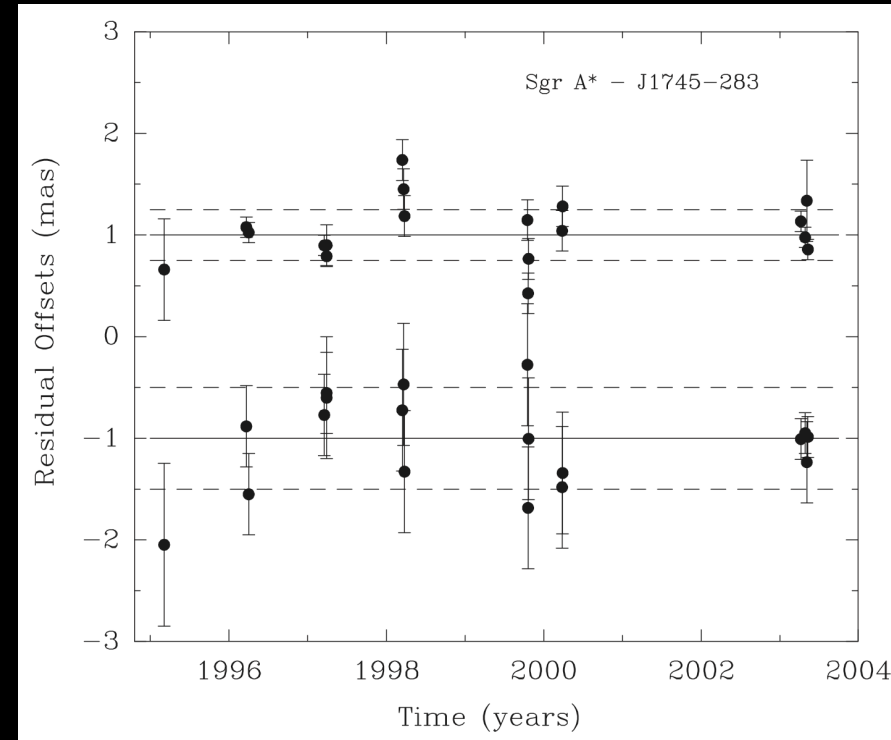


2. The Galactic Center: young stars near a massive black hole

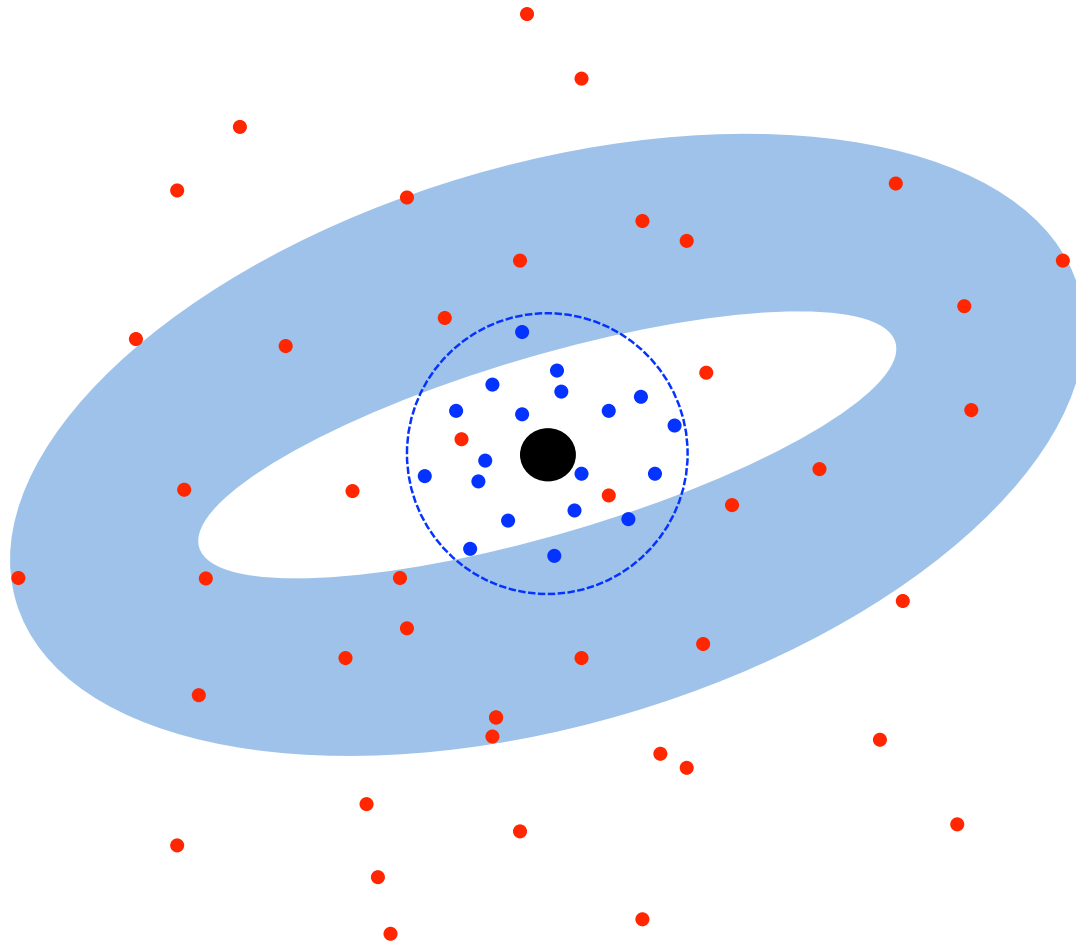
Recap

Sgr A* is a black hole

- $4 \times 10^6 M_{\text{sun}}$ inside of S2
(Schödel+ 2002, Gillessen+ 2009)
- $> 10\%$ of this is Sgr A*
(Reid & Brunthaler 2004)
- Sgr A* radio size: $\sim 4 R_s$
(Bower+2006, Doeleman+2008)
- density: $\sim 10^{-2}$ of black hole



Cartoon version of the stellar system

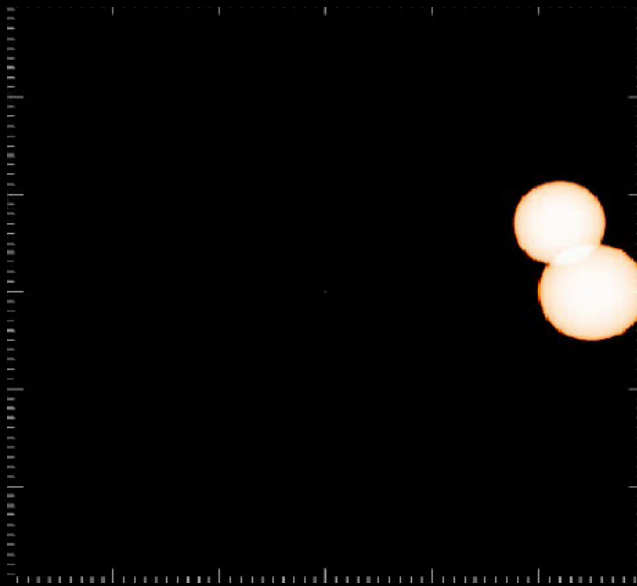


- S-stars
 - young, 10^8 yr
 - $r < 0.05$ pc
 - orbits
- Stellar disk
 - younger, 10^7 yr
 - $0.05 \text{ pc} < r < 0.5 \text{ pc}$
- Old stars
 - everywhere
- and more:
 - stellar black holes
 - neutron stars
 - white dwarfs
 - fainter MS stars

Two paradoxes of youth

O/WR stars

- 6 Myr young
- mean eccentricity 0.35
- disk configuration
- top-heavy IMF



Bonnell+ 2008, Hobbs+ 2009

B stars

- typical age 100 Myr
- eccentricities > 0.8
- randomly oriented
- normal IMF



Hills 1998, Perets+ 2007, Genzel+ 2010

Last time: tidal radius

- Binary a, m ; black hole M
- i) Forces: $F_t = F_g$ or $GMm a / r^3 = Gm^2/a^2$
- ii) Energy: $W = U_B$ or

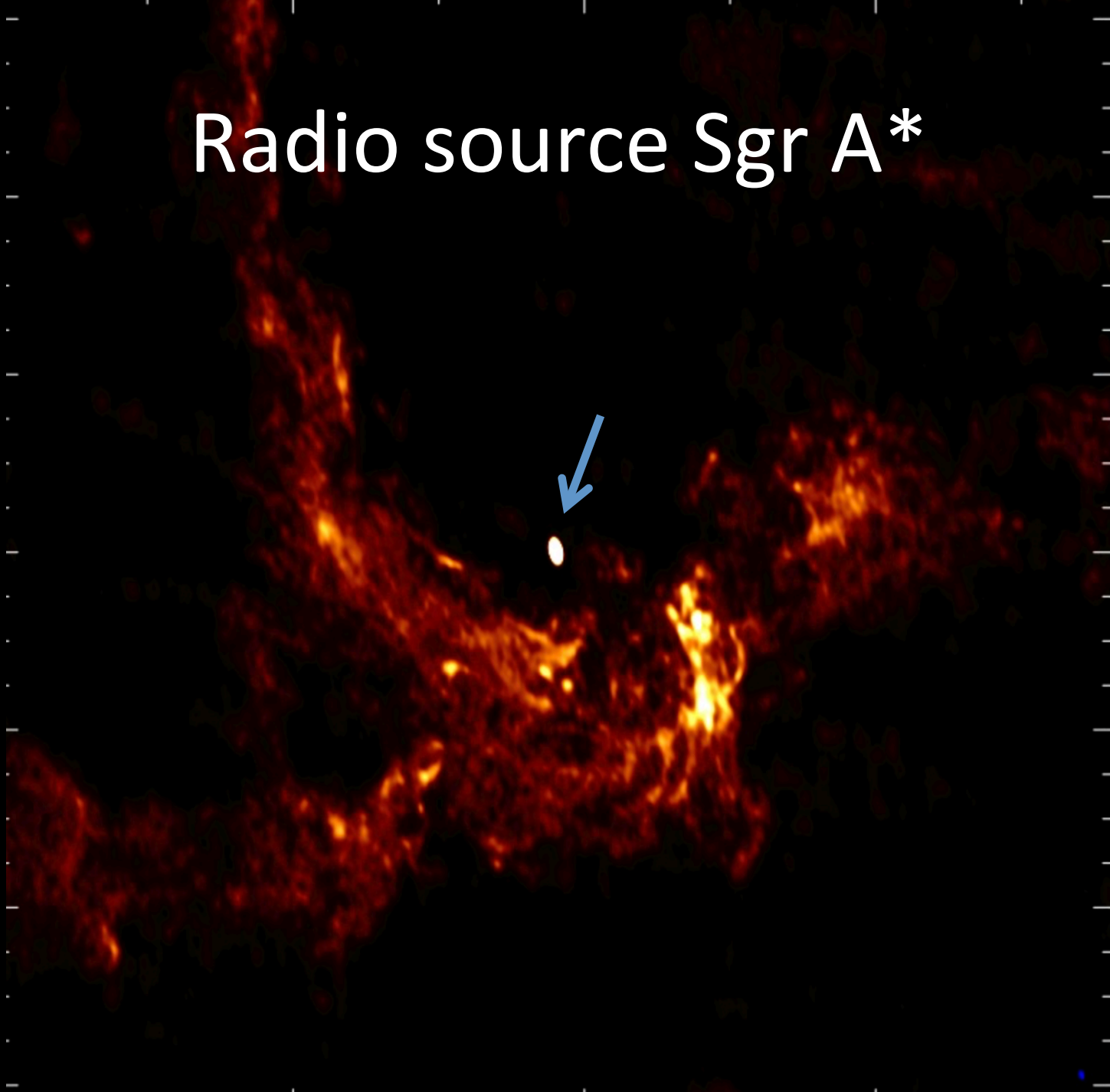
$$\int_0^a F_t dx = \int_0^a F_g dx$$

$$\rightarrow F_t = F_g$$

3. The Galactic Center

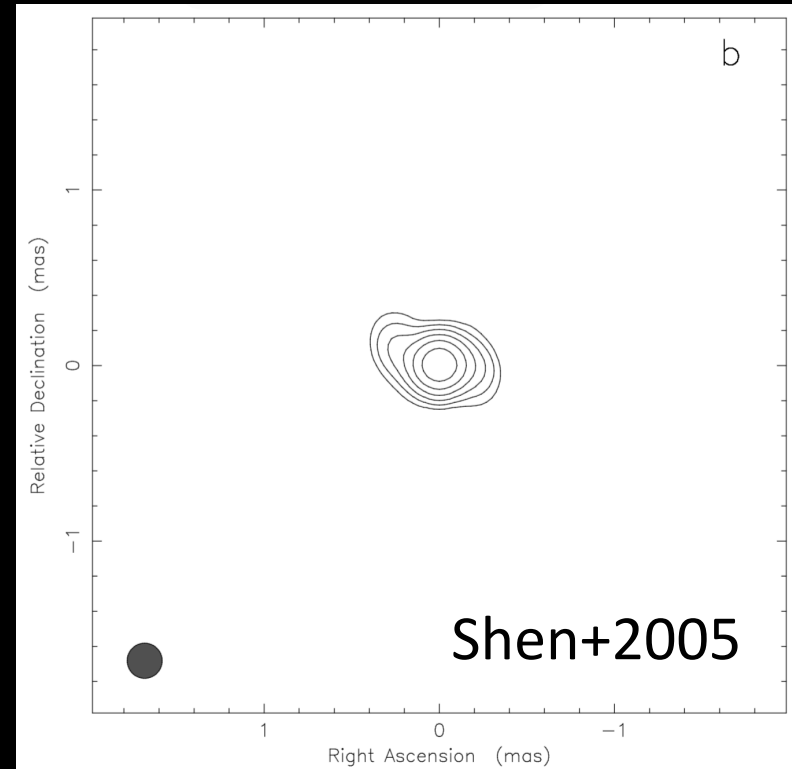
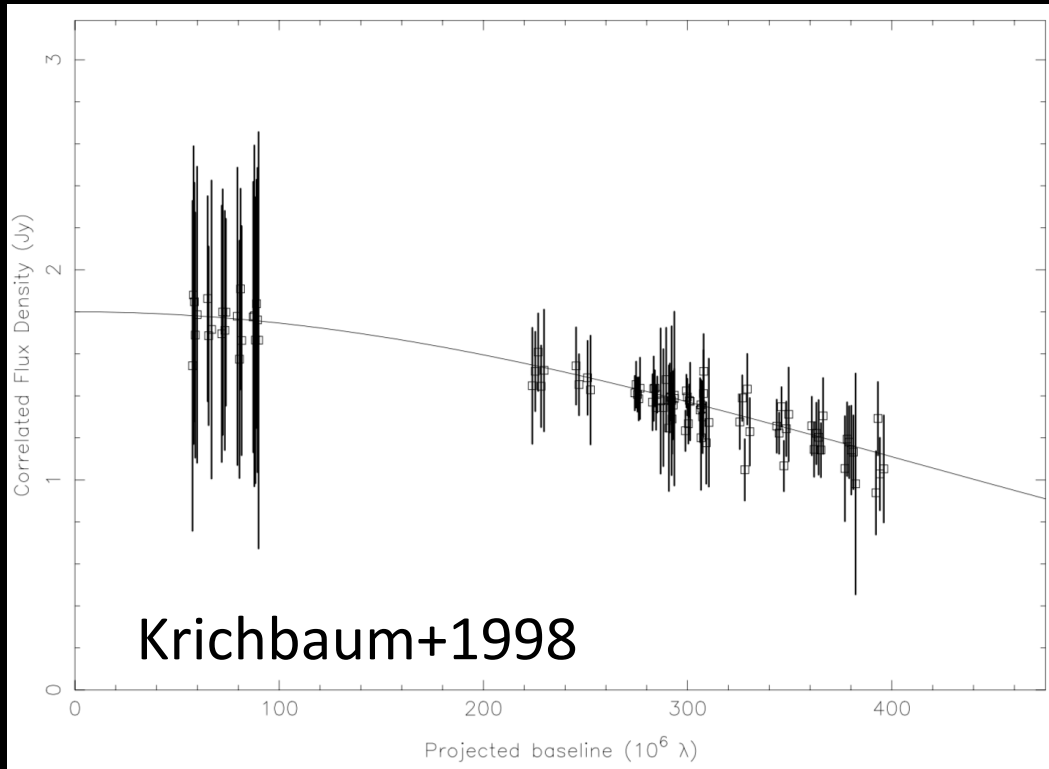
Sgr A* and the faintest black holes

Radio source Sgr A*

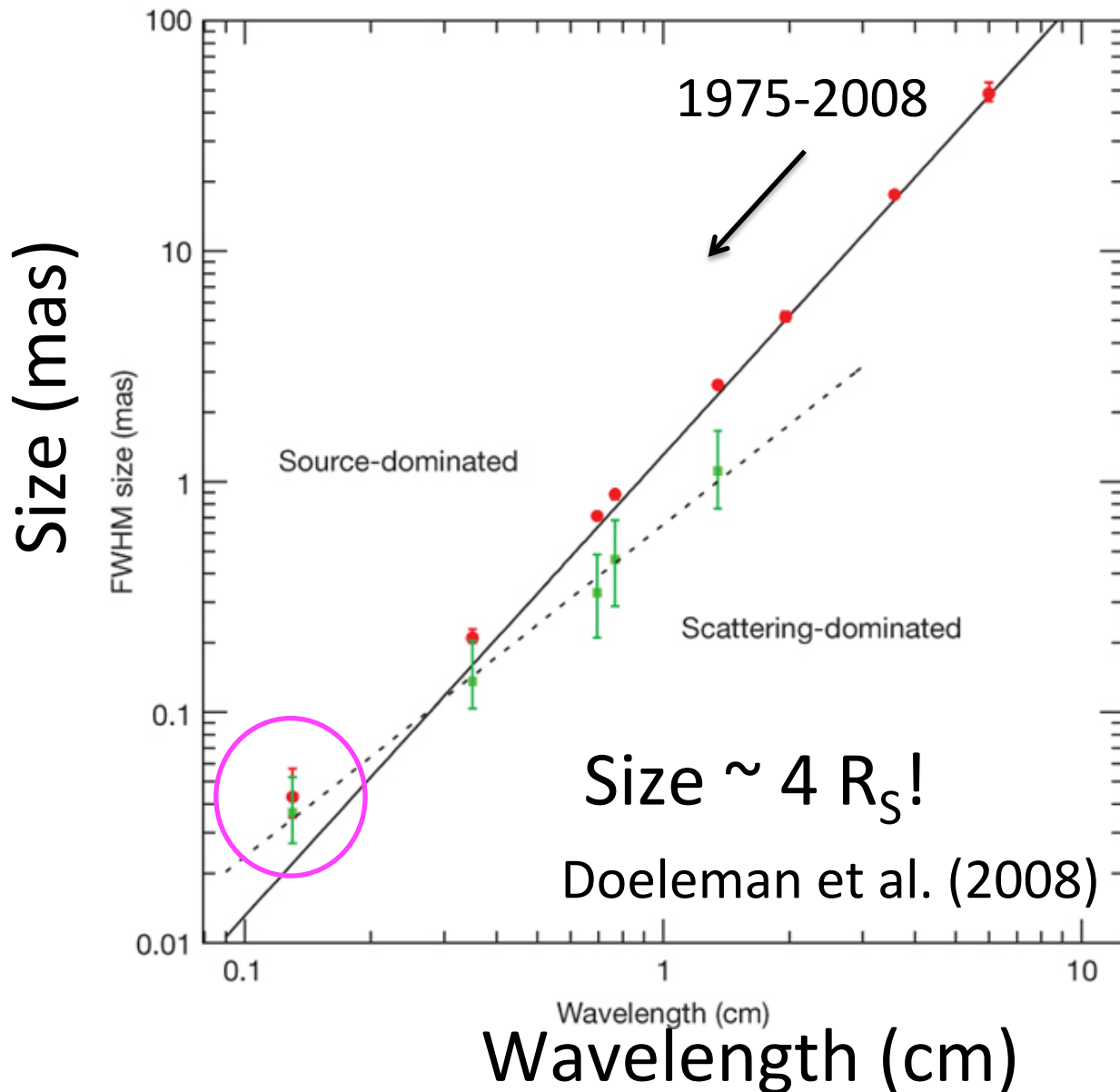


Push to higher resolution

- VLBI (VLBA, EVN) observations
(Lo, Moran, Krichbaum, Bower, ...)
- Atmosphere changes every ~ 10 -100s



The radio size of Sgr A*



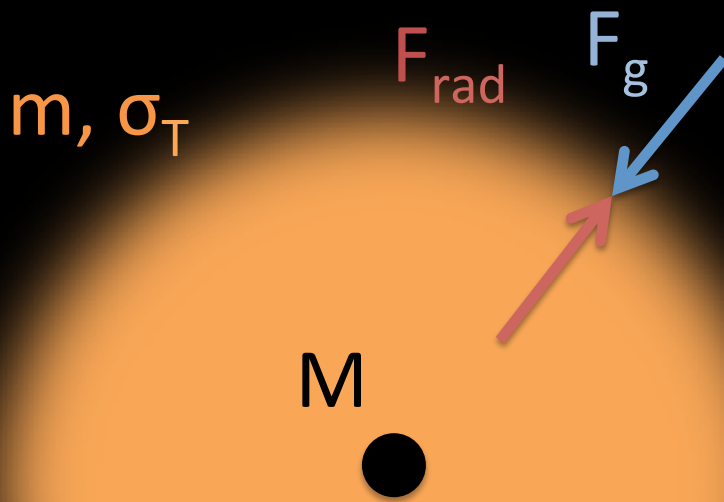
Accretion power

- Infalling gas radiates its gravitational energy



Eddington Luminosity

- How bright can accretion be?



$$F_g = F_{\text{rad}}$$
$$GMm_p/r^2 = L / 4\pi r^2 \sigma_T / c$$

$$L_{\text{edd}} = 4\pi Gm_p c M / \sigma_T$$
$$L_{\text{edd}} \sim 10^{38} \text{ ergs / s } (M/M_{\text{sun}})$$

Sgr A*:

$$L_{\text{edd}} \sim 10^{45} \text{ ergs / s}$$

$$L_{\text{bol}} < 10^{37} \text{ ergs / s}$$

Eddington accretion rate

- What is minimum amount of infalling material to get L ?

m, σ_T



Maximum L :

$$L = dE/dt = d/dt(mc^2)$$

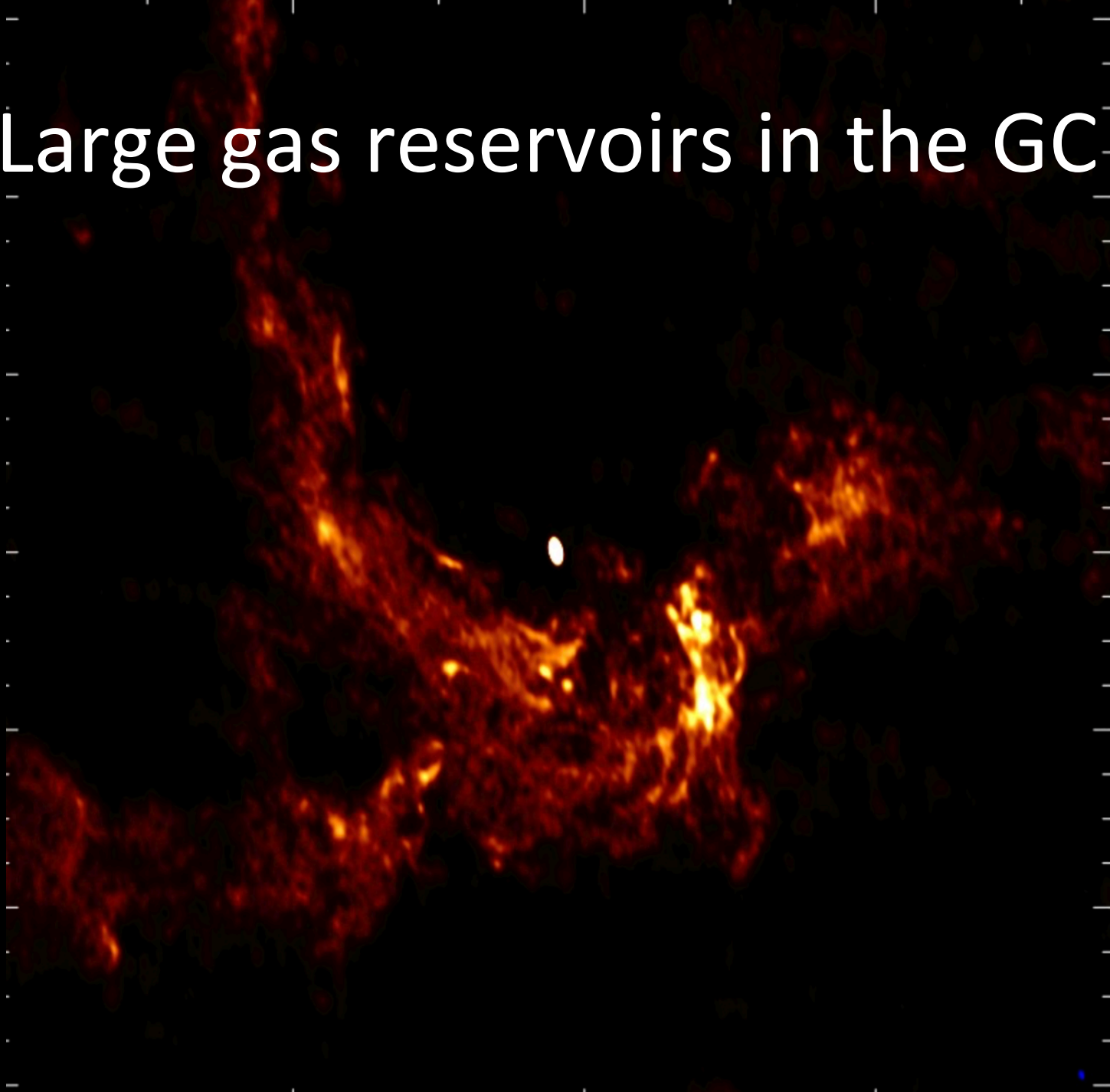
$$L = \dot{M} c^2$$

General L :

$$L = \epsilon \dot{M} c^2$$

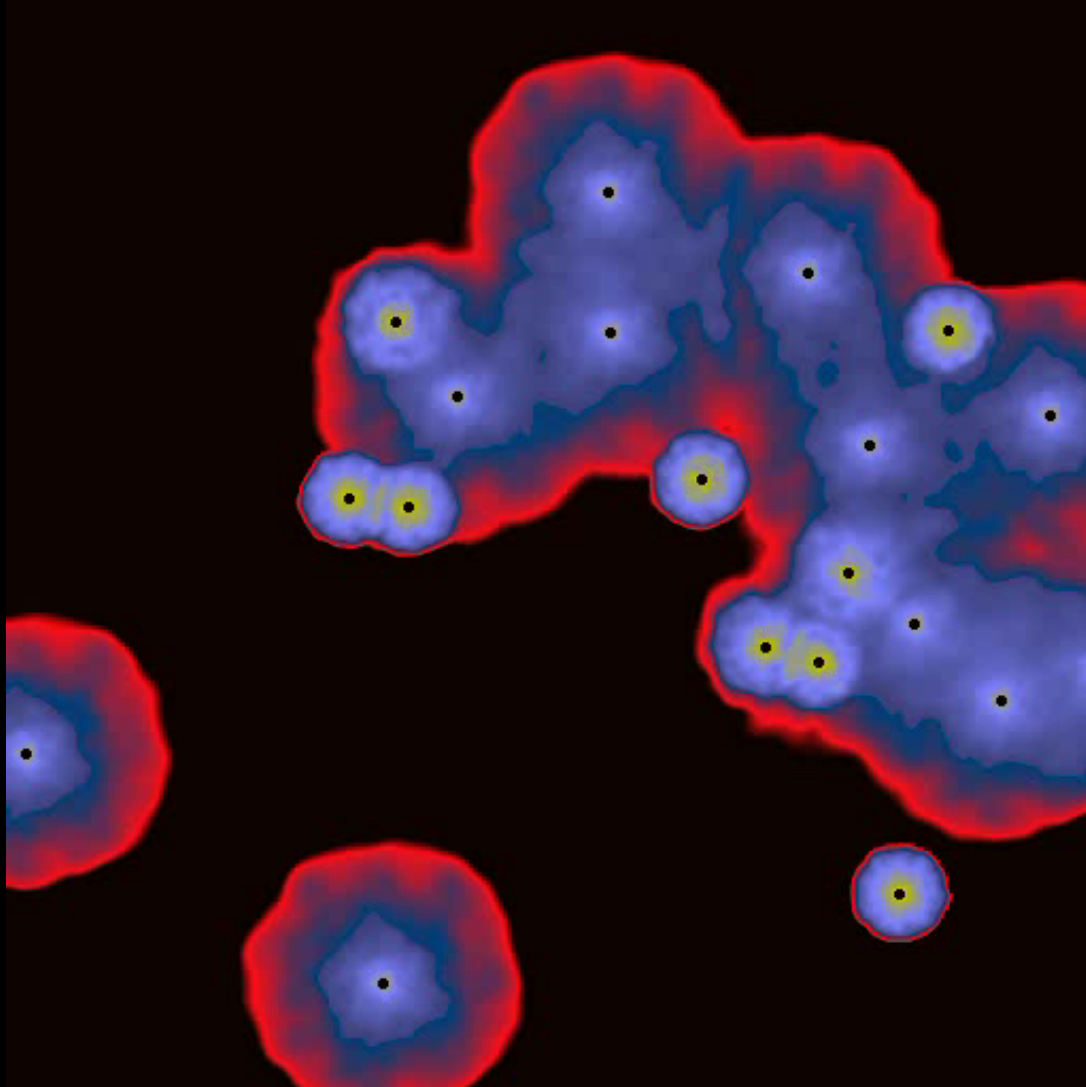
ϵ = “accretion efficiency”

Large gas reservoirs in the GC



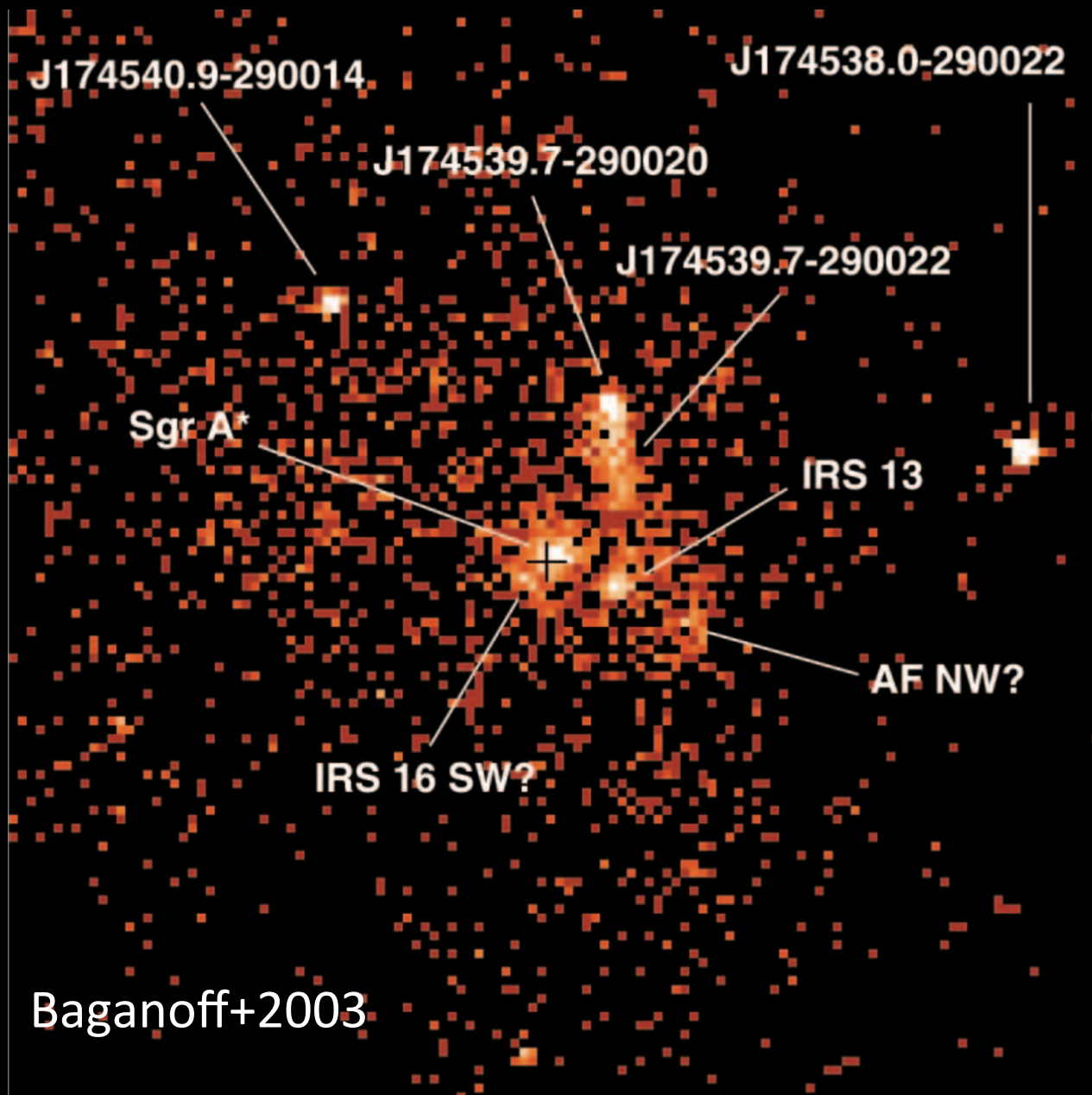
Stellar winds: black hole fuel

- $\dot{M} \sim 10^{-4} - 10^{-3} M_{\text{sun}} / \text{yr}$ (Martins+2007)



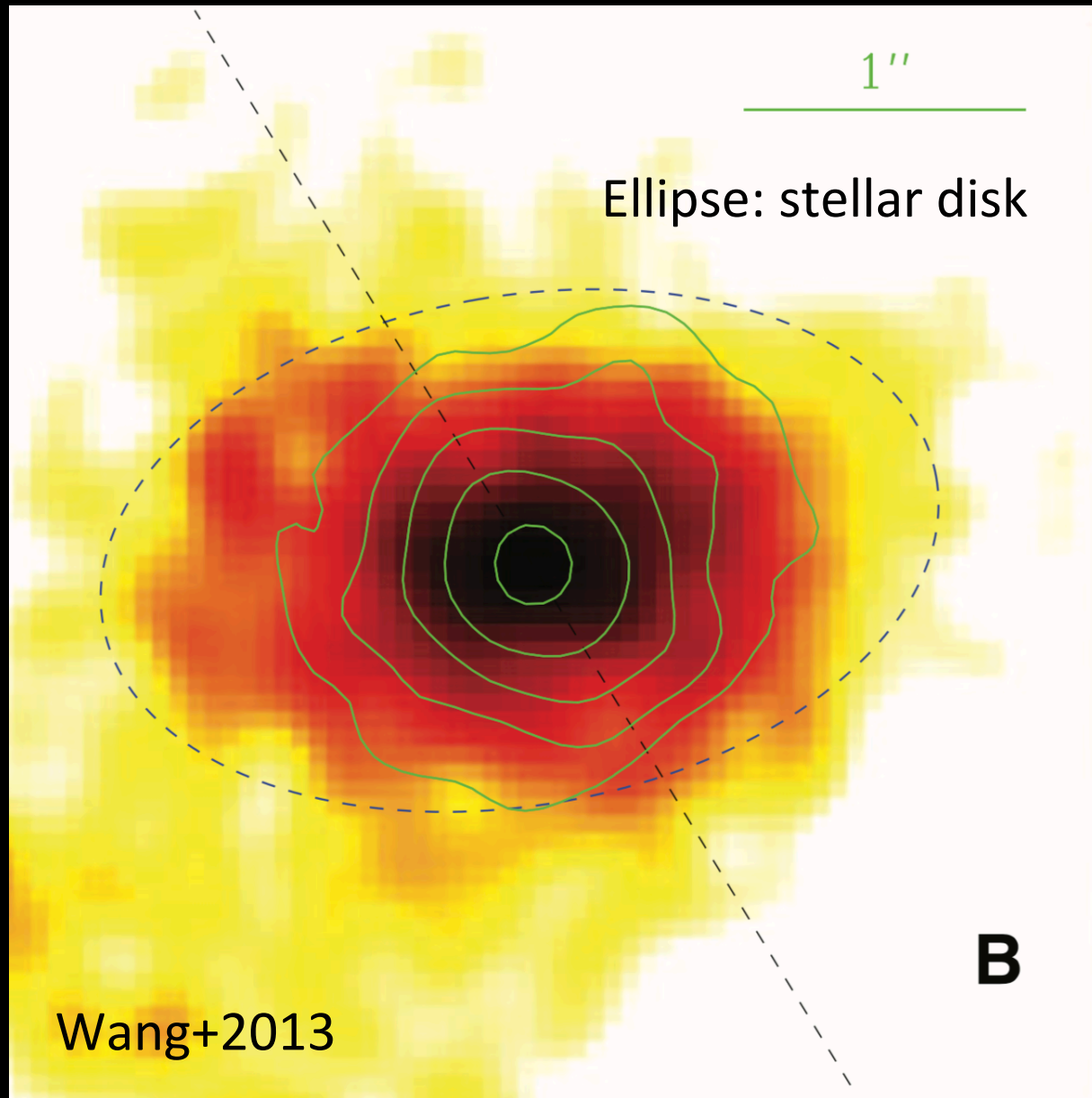
Cuadra
+2006, 2008

X-rays from Sgr A*



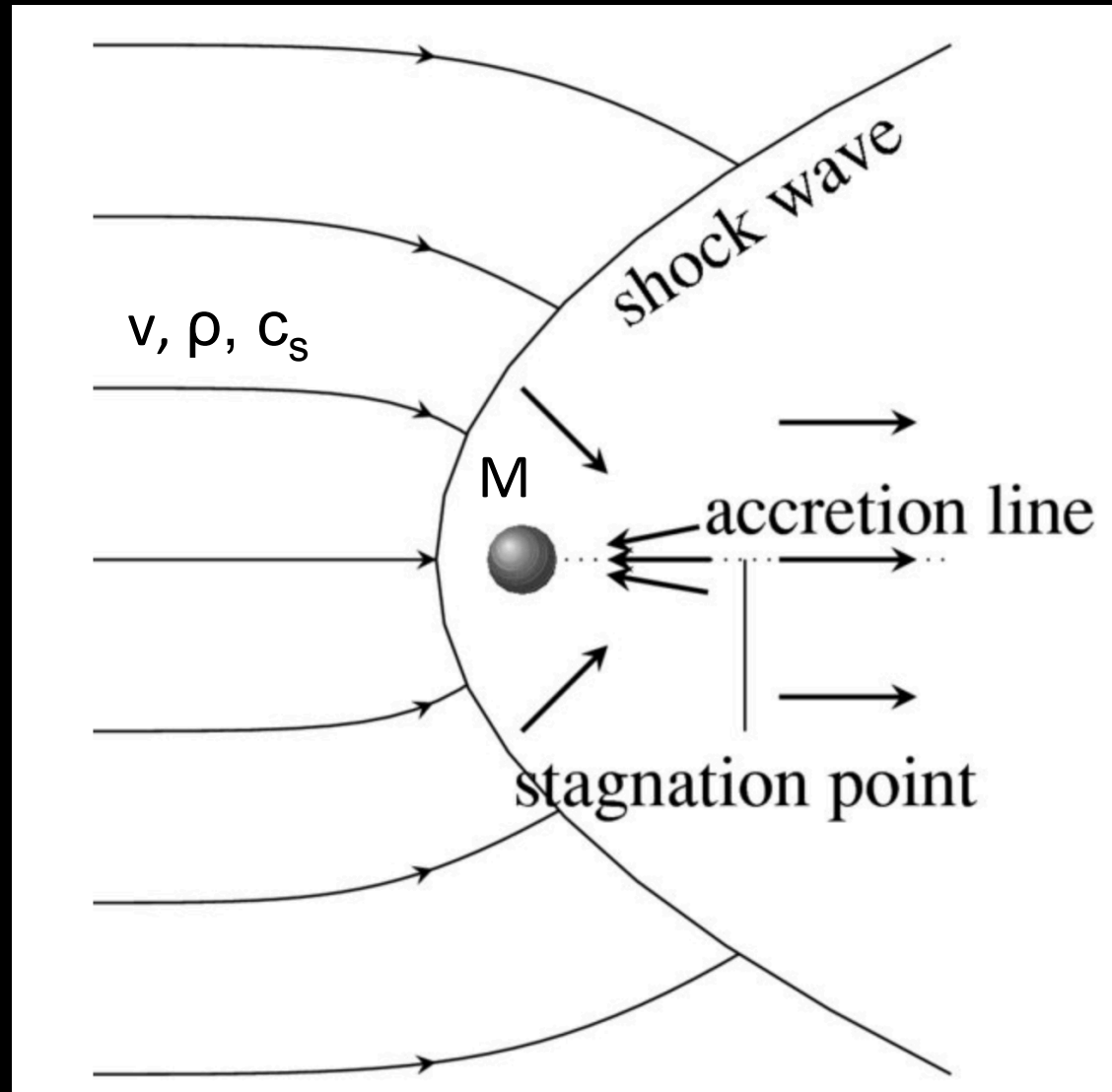
Baganoff+2003

X-rays from Sgr A*



Bondi-Hoyle accretion

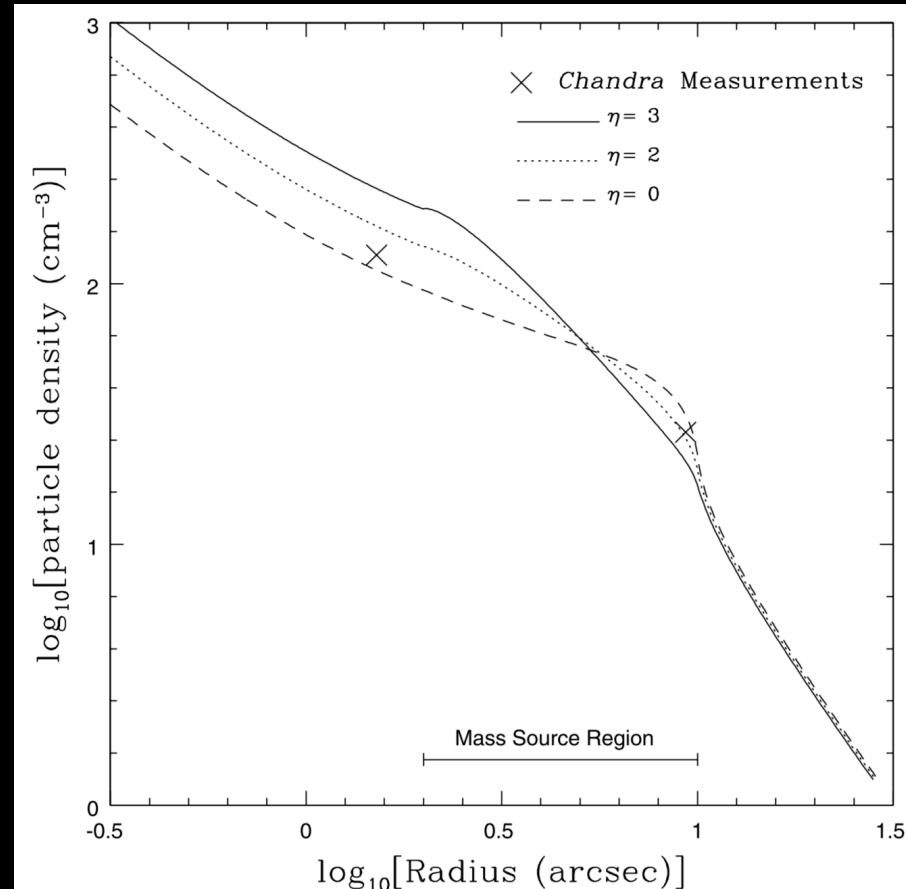
- $\dot{M} = \frac{4\pi \rho G^2 M^2}{(v^2 + c_s^2)^{3/2}}$



Bondi-Hoyle accretion

- $\dot{M} = 4\pi \rho G^2 M^2 / (v^2 + c_s^2)^{3/2}$
- X-rays:
 $n \sim 100 \text{ cm}^{-3}$,
 $T \sim 10^7 \text{ K}$
- Winds: $v \sim 1000 \text{ km/s}$
- $\dot{M} \sim 10^{-5} M_{\text{sun}} / \text{yr}$

Quataert+1999, 2004

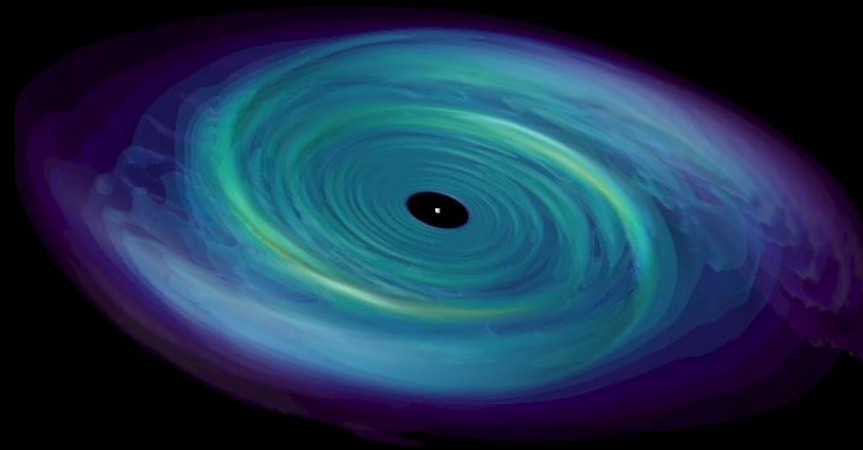


Q: Why is it difficult for gas to fall into a black hole?

- Sgr A* Bondi radius ~ 0.1 pc: what is the angular momentum of a circular orbit?
- What fraction of this angular momentum must be lost to reach the event horizon, $R \sim 10^{-6}$ pc?
- Or: what eccentricity is needed for orbit from Bondi radius to reach event horizon?

Accretion theory

- Hydrodynamics equations
 - Collisional, “viscous” infalling material
- Approximations:
 - Stationary, axisymmetric ($d_t = d_\phi = 0$)
 - $\dot{M} = \text{constant}$
 - Vertically integrated
 - “Thin” disk $H/R \ll 1$,
in practice $H/R \leq 1$



Standard accretion theory

- Mass:

After Blaes (2004)

$$\dot{M} = 4\pi R H \rho v$$

- Momentum:

$$\rho v \, d_R v = \rho(\Omega^2 - \Omega_K^2) R - d_R p$$

- Angular momentum:

$$\dot{M} \, d_R(\Omega R^2) = d_R(4\pi R^2 H \tau_{R\phi})$$

- Energy:

$$\frac{\dot{M}}{2\pi \rho R} T \, d_R s = 2R H \tau_{R\phi} \, d_R \Omega + 2F^-$$

Thin disk accretion

Shakura &
Sunyaev 1973)

- $\Omega = \Omega_K$, $\dot{M} = \text{constant}$, cooling is fast, no torque at inner edge (R_i)

- Angular momentum:

$$\dot{M} d_R(\Omega R^2) = d_R(4\pi R^2 H \tau_{R\phi})$$

$$\dot{M}(\Omega_K R^2 - \Omega_{K,i} R_i^2) = 4\pi H R^2 \tau_{R\phi}$$

- Energy:

$$\frac{\dot{M}}{2\pi \rho R} T d_R s = 2R H \tau_{R\phi} d_R \Omega + 2F^-$$

$$F^- = -R H \tau_{R\phi} d_R \Omega_K$$

Thin disk accretion

Shakura &
Sunyaev 1973)

- Solve for flux emerging from disk, F^-

$$\dot{M}(\Omega_K R^2 - \Omega_{K,i} R_i^2) = 4\pi H R^2 \tau_{R\phi}$$

$$F^- = -R H \tau_{R\phi} d_R \Omega_K$$

- $\Omega_K^2 = GM/R^3$, $d_R \Omega_K = -3/2 \Omega_K / R$:

$$F^-(r) = \frac{3GM\dot{M}}{8\pi R^3} \left(1 - \sqrt{\frac{R_{in}}{R}} \right)$$

$$F = 2 \int_{R_i}^{\infty} dR 2\pi R F^-(R)$$

Thin disk accretion

Shakura &
Sunyaev 1973)

- Solve for total luminosity integrated over disk:

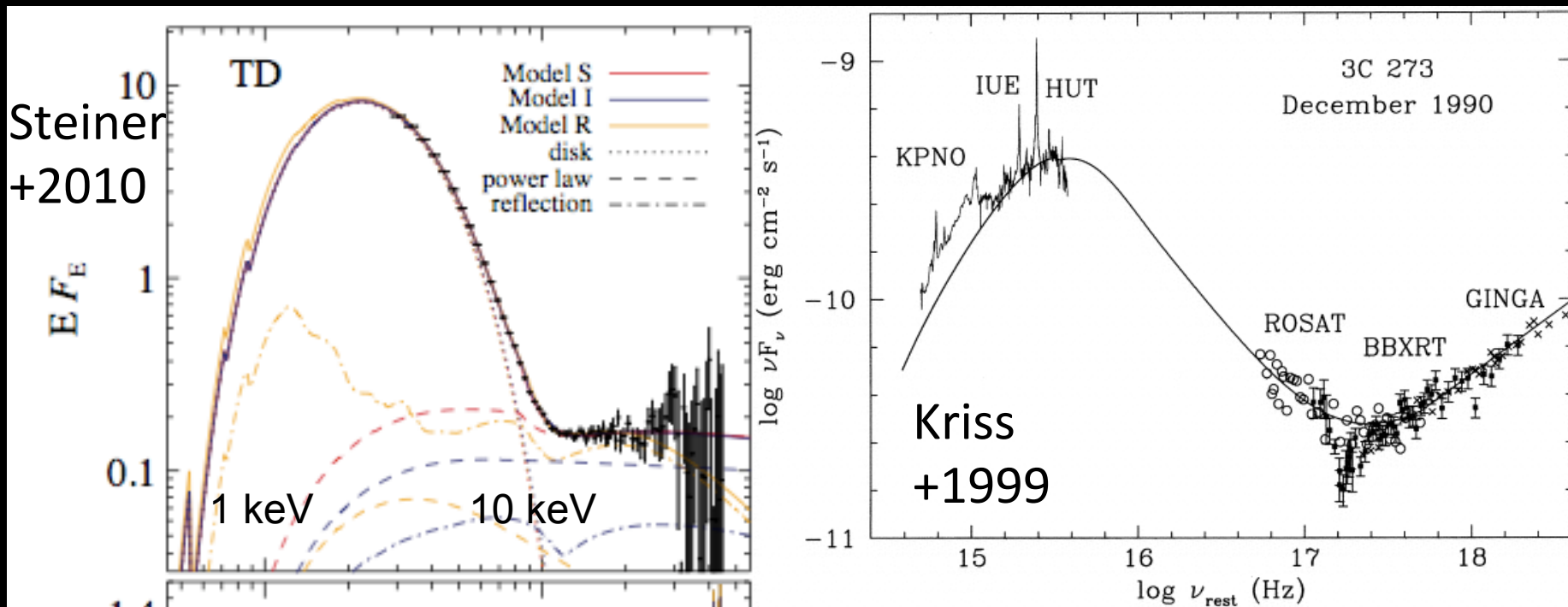
$$L = 2 \int_{R_i}^{\infty} dR \, 2\pi R \, F^{-}(R)$$
$$L = \frac{GM\dot{M}}{2R_i}$$

- $R_i \sim$ innermost stable circular orbit of a black hole = $6 GM/c^2$ for spin zero
- $L \sim 10\% \, \dot{M} c^2$!
- Compare: nuclear fusion $\sim 0.1\% \, \dot{M} c^2$

Thin disk accretion

Shakura &
Sunyaev 1973)

- opt. thick, so $F^-(R) = \sigma T_{\text{eff}}(R)^4$
- $R \sim M$, $\dot{m} = \dot{M} c^2 / L_{\text{edd}}$
- $T_{\text{eff}}(r) \sim \dot{m}^{1/4} M^{-1/4} r^{-3/4}$
- Stellar black holes in X-rays, AGN in UV



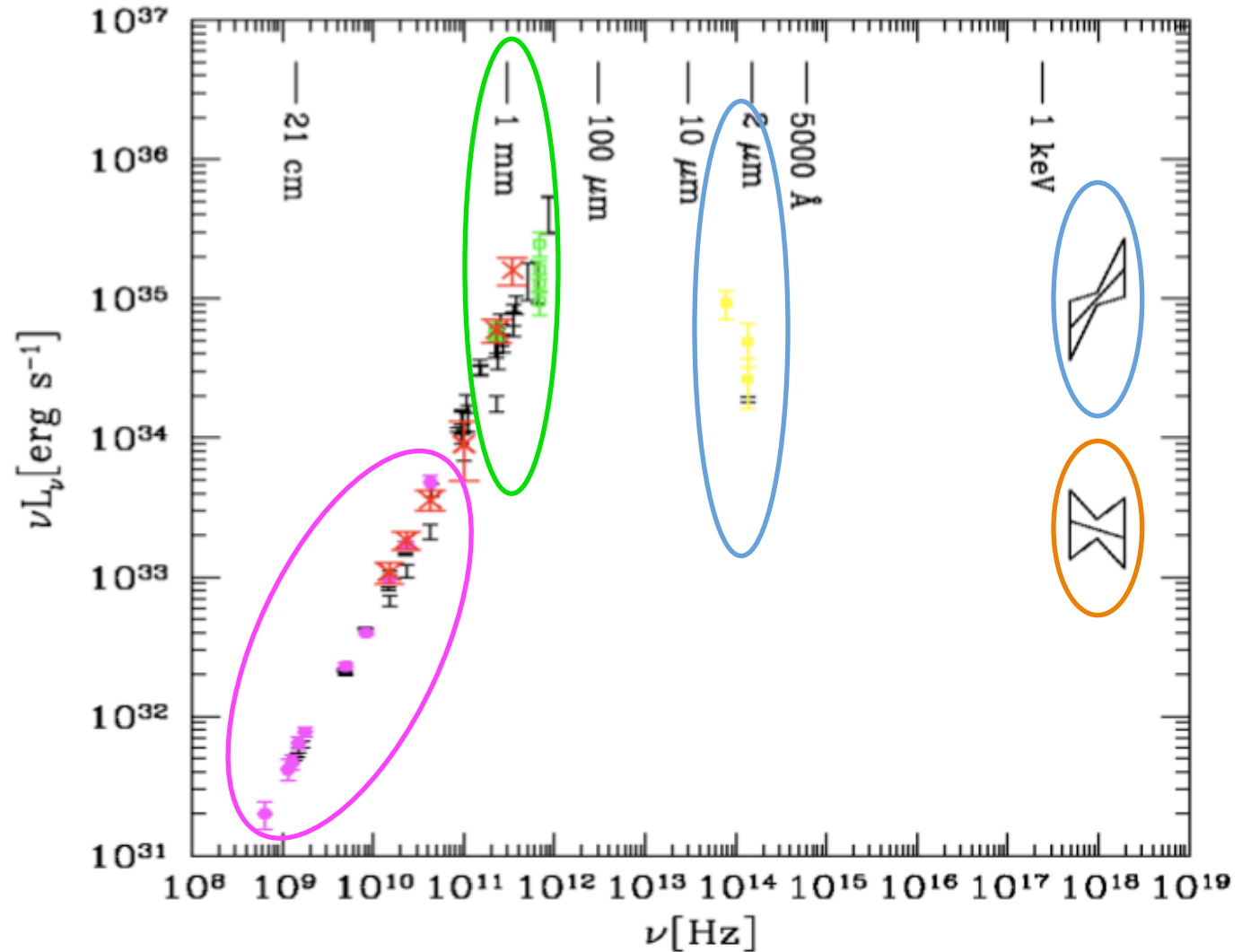
Sgr A* spectrum

Radio: Balick &
Brown 1974

mm: Zylka &
Mezger 1988

NIR/X-ray:
Genzel+2003;
Ghez+2004;
Baganoff+2001

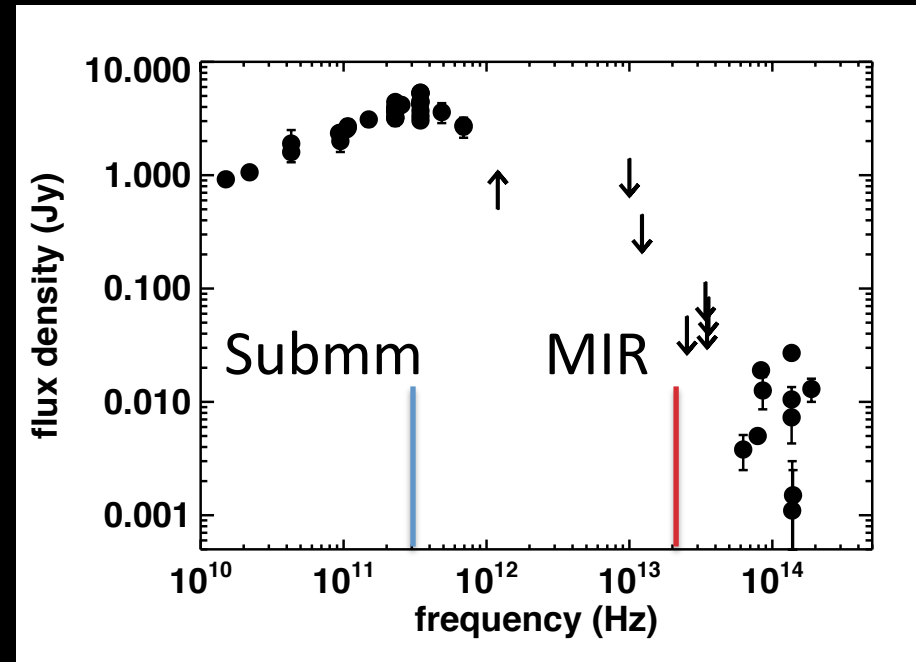
X-ray:
Baganoff 2003



Charles Gammie, Rainer Schödel

Thin disks don't work for Sgr A*

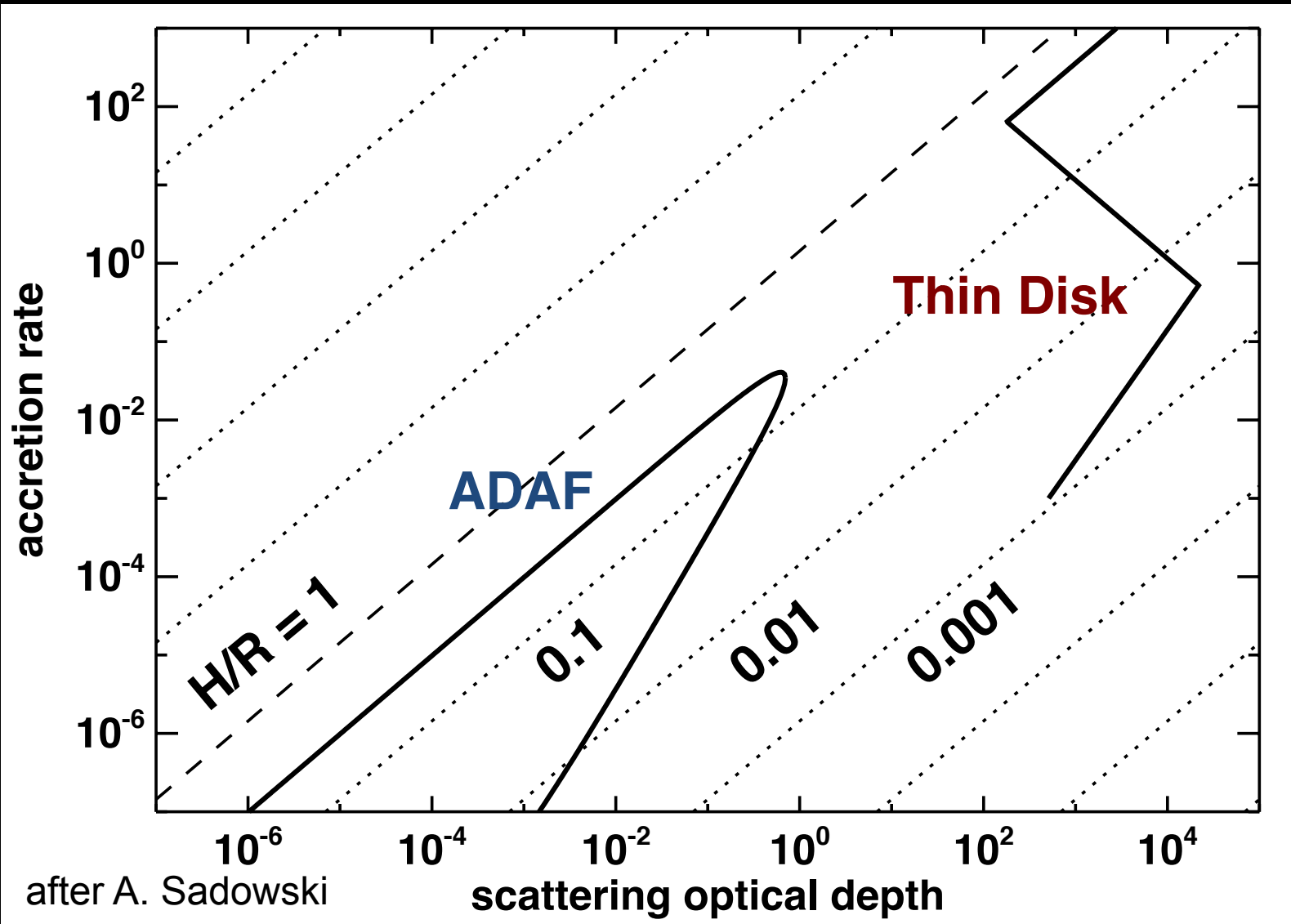
- Thin disk:
 $L \sim 10^5$ too bright;
spectrum should peak in IR not submm
- If $\dot{M} \sim \text{Bondi}$,
 $\epsilon = L / \dot{M} c^2 \sim 10^{-5}$!
- Why is Sgr A* so faint?



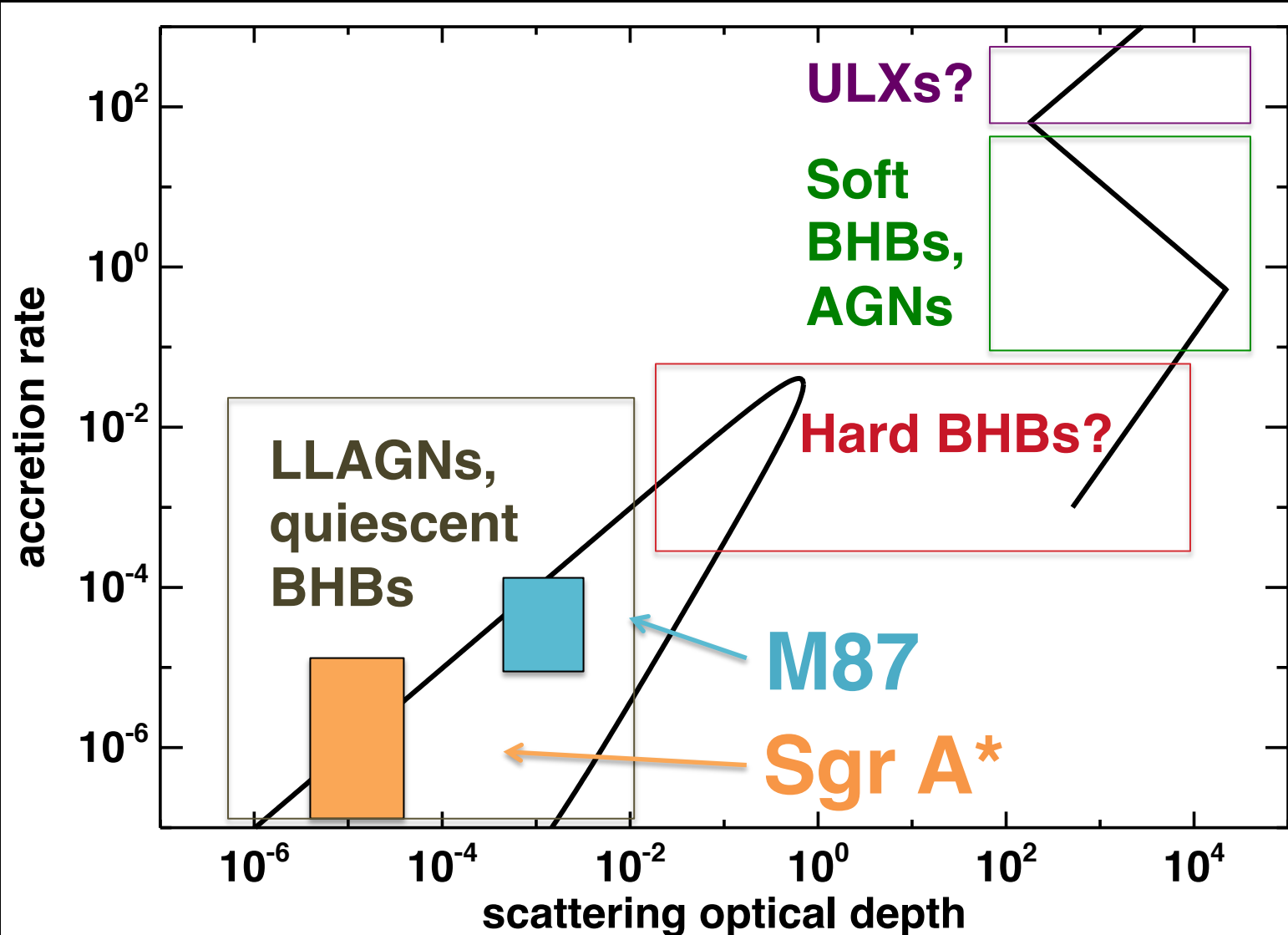
Standard Accretion Theory

- **ADAF** (Narayan & Yi 1994, also Ichimaru 1977, Rees+1982)
 - Parameterize fraction of local heating advected:
 - $f=0$: $\Omega = \Omega_K$, thin disk
 - $f=1$: $\Omega = 0$, spherical (Bondi) accretion
 - Self-similar scalings like spherical accretion:
$$\rho, \Omega \propto R^{-3/2}; \quad v, c_s \propto R^{-1/2}$$
 - Low density, “collisionless,” $\rightarrow$
different proton and electron temperatures!

Standard Accretion Theory

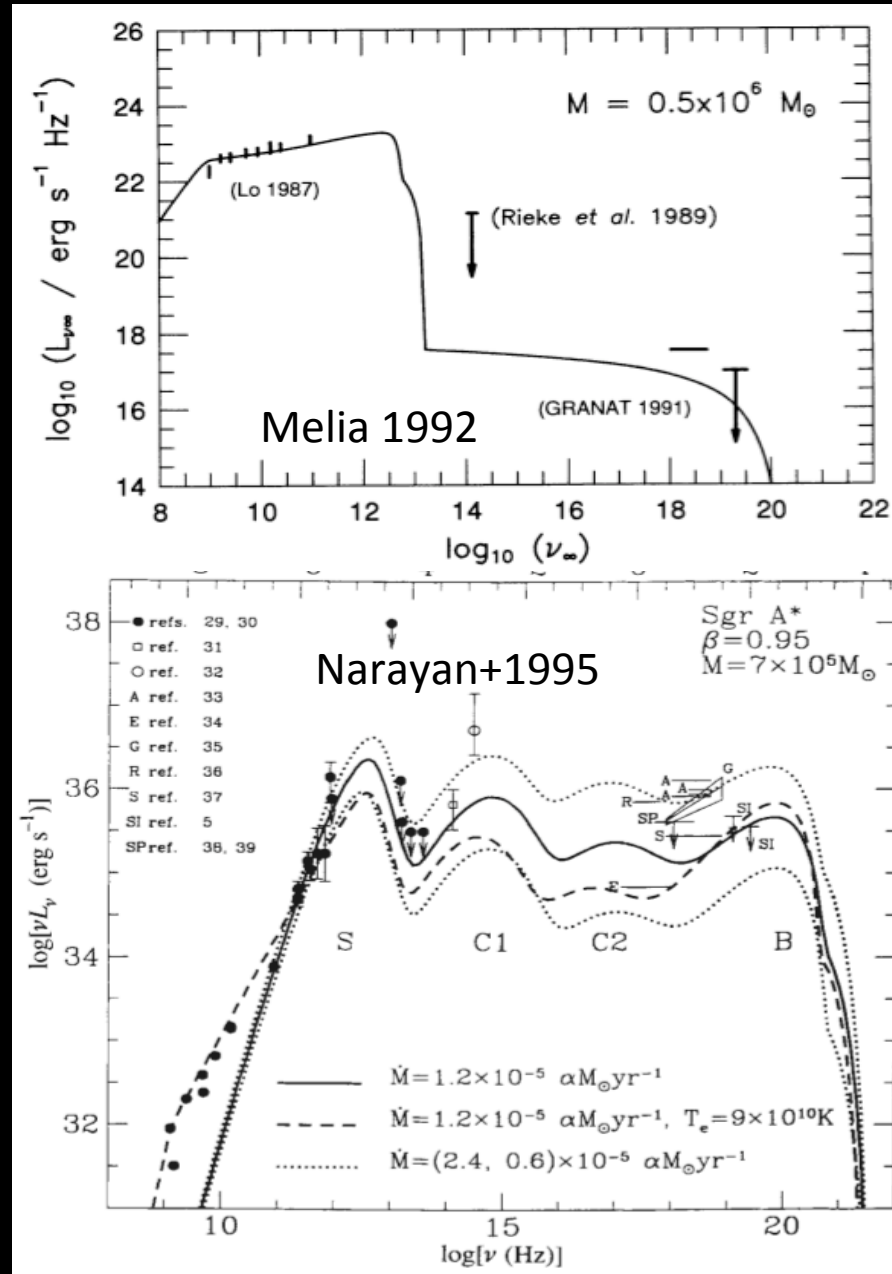


Standard Accretion Theory



Accretion Flow Models of Sgr A*

- Spherical (Melia 1992)
ADAF (Narayan+1995)
 - $dM/dt \approx dM/dt_{\text{Bondi}}$
 - Lower T_e , large central density ($n \sim r^{-3/2}$)
- Accretion energy is carried into the black hole!

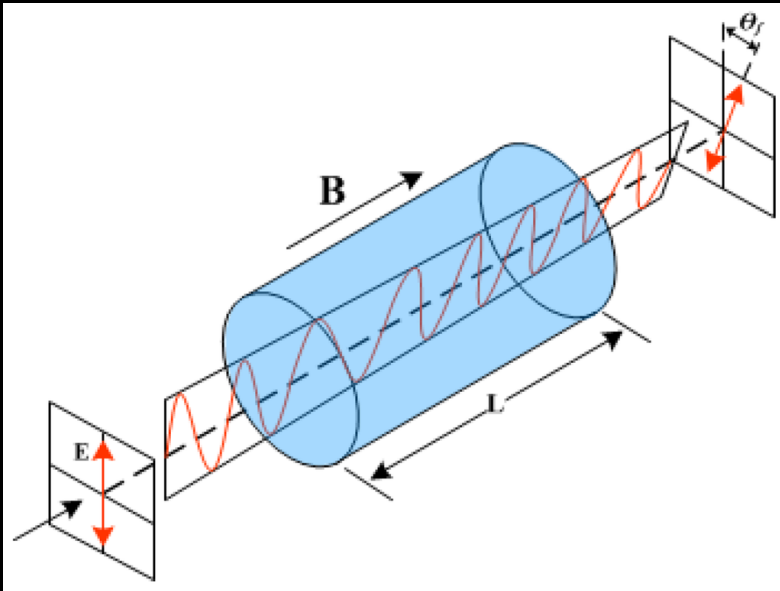


Polarization, Faraday rotation, and the Sgr A* accretion rate

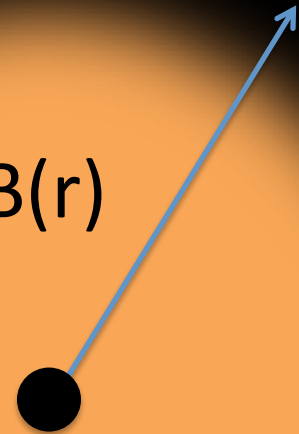
- Magnetic field rotates polarization direction

$$\chi(\nu) \sim \nu_0 + \text{RM} \, c^2/\nu^2$$

$$\text{RM} \simeq 2.6 \times 10^{-13} \int dl \cdot B \, n \, \text{rad/m}^2$$



$n(r), B(r)$



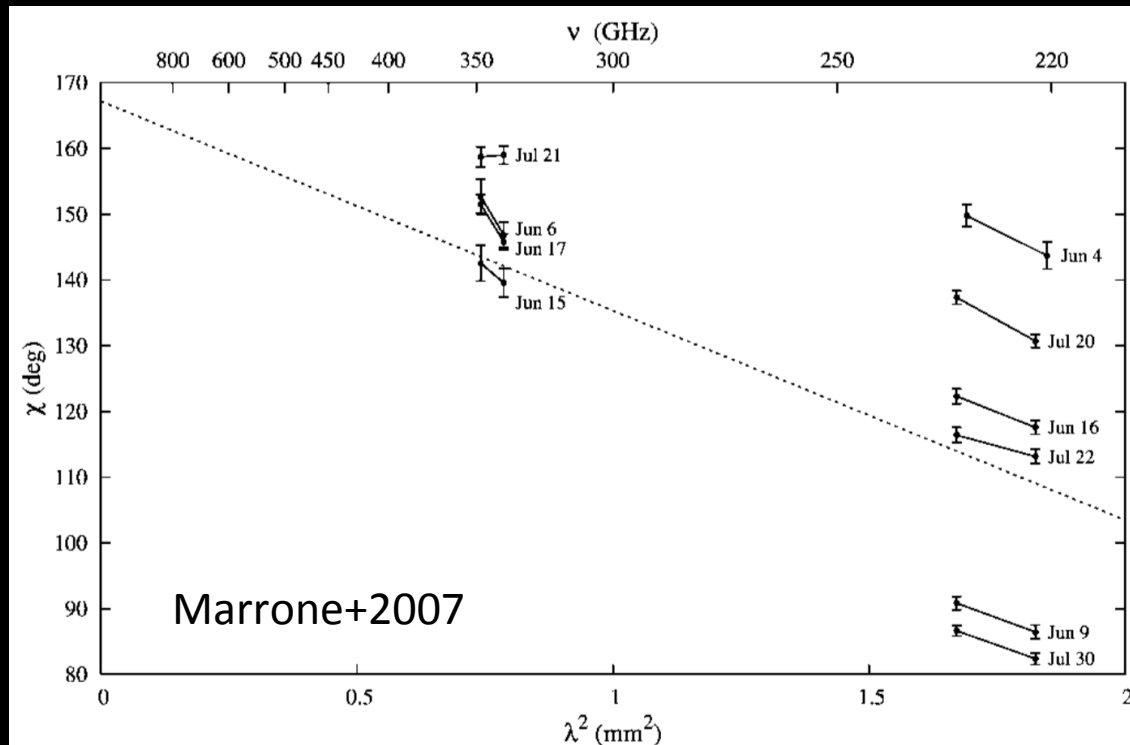
Polarization, Faraday rotation, and the Sgr A* accretion rate

- $RM \sim M^{-2} \dot{M}^{3/2}$ (Marrone et al. 2006)

$$\dot{M} \simeq 10^{-9} r_{\text{NR}}^{7/6} M_{\odot} \text{yr}^{-1}$$

$$r \sim 1-100 R_s$$

- > 99% of the mass doesn't make it!
(Aiken+2000, Agol 2000, Quataert & Gruzinov 2000, Bower+2003, Marrone+2006,2007)



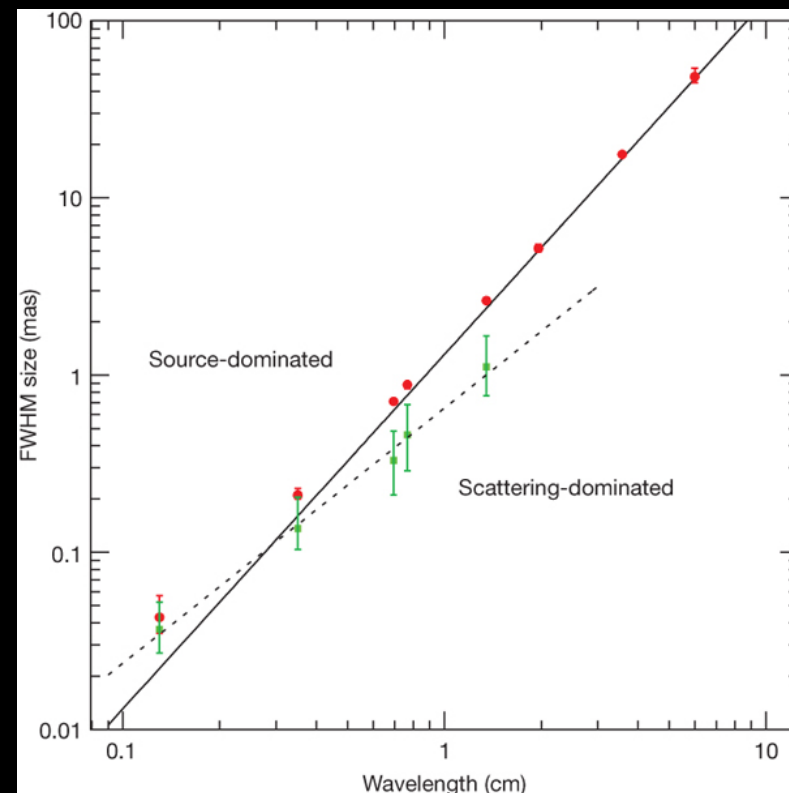
Sgr A* e- temperature

- For known flux, angular size “brightness” temperature is $B_\nu (T_b) = I_\nu$
- Sgr A* at 230 GHz:

$$T_b = \frac{c^2 I_\nu}{2k\nu^2}$$
$$\simeq 6 \times 10^{10} \text{ K}$$

cf. $T_i \approx T_{\text{vir}} \approx 10^{12} \text{ K}$

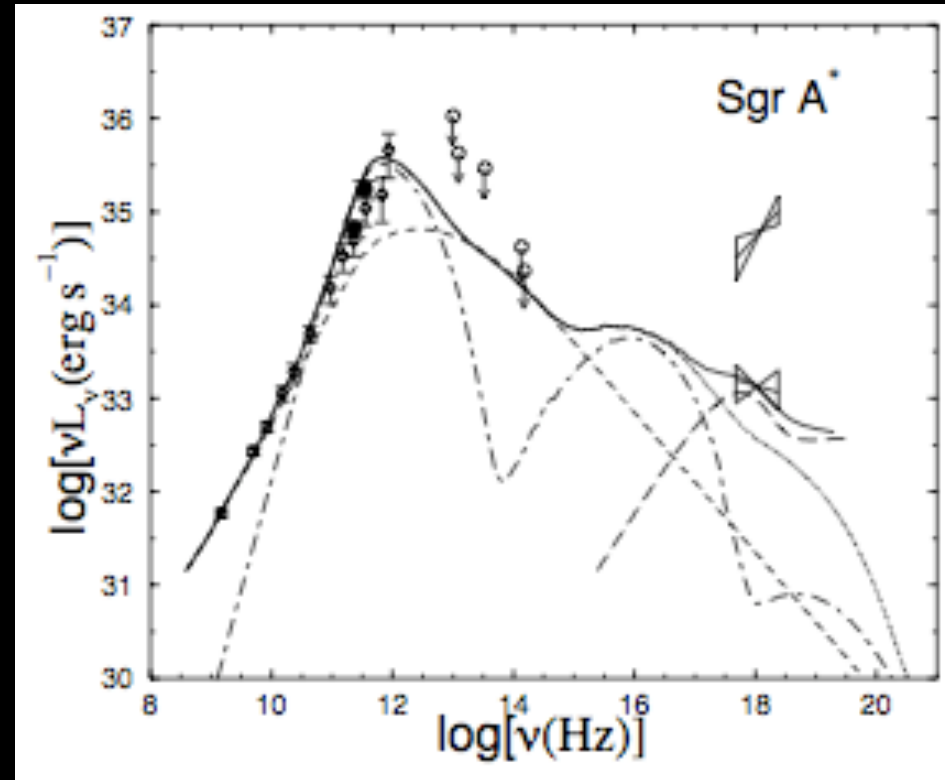
- Lower limit to T_e ,
 $\gg 10^9 \text{ K}$ in ADAF



Doeleman et al. 2008

Accretion flow models of Sgr A*

- Outflows $\rightarrow$ ADIOS
(Blandford & Begelman 1999)
- Convection $\rightarrow$ CDAF
(Quataert & Gruzinov 2000)
- ADAF/CDAF/ADIOS/
... $\rightarrow$ RIAF!
 - Mass loss!
 - Non-thermal e- or jet
for polarization



Yuan et al. (2003)

Why is Sgr A* so faint?

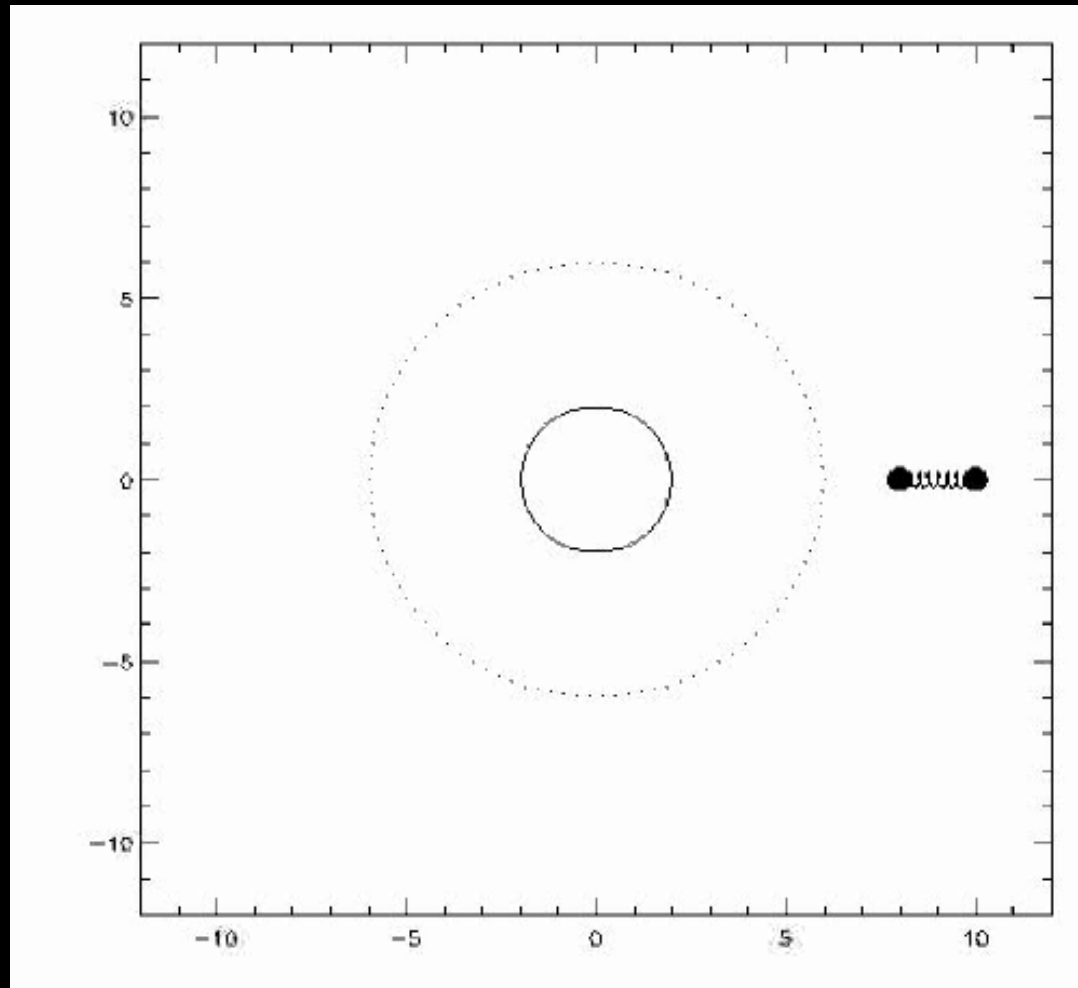
	Radius (R_s or pc)	Accretion rate ($M_{\text{sun}} / \text{yr}$)	Accretion efficiency
Giant molecular clouds	10-100 pc	10^{-2}	10^{-8}
Circumnuclear disk	1-7 pc	10^{-3}	10^{-7}
Winds from massive stars	$< 0.5 \text{ pc} \sim 10^6 R_s$	$10^{-4} - 10^{-3}$	$10^{-6} - 10^{-7}$
Winds at Bondi radius	$< 0.1 \text{ pc} \sim 10^5 R_s$	10^{-5}	10^{-5}
Inner accretion flow	$1-10^3 R_s$	$10^{-9} - 10^{-7}$	$10^{-1} - 10^{-3}$

Sgr A* is faint because

- Gas supply (stellar winds) is not large enough
- A tiny fraction of gas supplied reaches the black hole!
- The accretion flow is inefficient at radiating away its gravitational binding energy

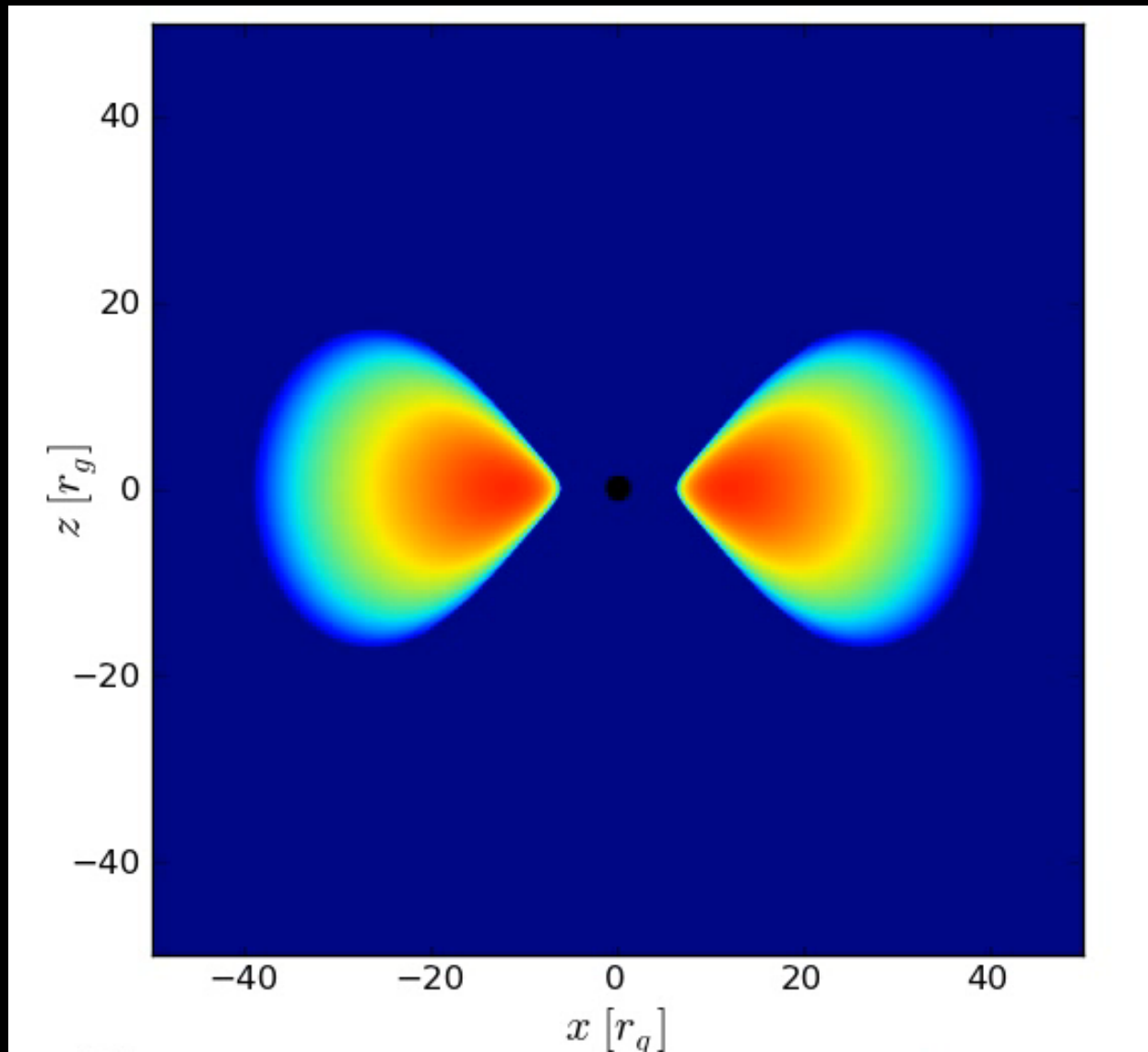
How does gas fall into a black hole?

- Weakly magnetized gas: field is dragged along, restoring force when stretched (torque!)



The MRI can cause accretion

- Instability $\rightarrow$ turbulence, stresses $\rightarrow$ torque



McKinney &
Blandford
2009

MHD simulations of BH accretion

- Good: physical theory of accretion!
- Bad: turbulence, magnetic fields, time-dependence, 3D → numerical simulations
- Still missing: plasma physics
- Sgr A* is a great laboratory for MRI accretion theory

Next time:

Has Sgr A* always been so faint?