component	M ($10^9 { m M}_{\odot}$)	gas mass fraction
total H II	1.12	23%
total H I	2.9	60%
total H_2	0.84	17%
total gas including He	6.7	-
stars	50	-
dark halo	$pprox 2 imes 10^3$	

 The phases of the ISM behave to a certain degree as phases of a fluid: solid - liquid - gaseous.

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- Density (temperature) sets the state of the phase.
- Star forming is exclusively in the H_2 phase.

Star forming regions



(Rivera-Ingraham et al. 2013)

- In the Galaxy, star formation occurs within molecular clouds.
- Typical mass $10^4 {\rm M}_{\odot}$ to $5\times 10^6 {\rm M}_{\odot};$ self-gravity.
- Typical scales 5pc to 100pc.
- Molecular clouds are turbulent.
- Density typically above $400 {\rm cm}^{-3}$. Enveloped in H I gas.

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Star forming regions



(Rivera-Ingraham et al. 2013)

- Life-time of several Myr. They are transient, they are not long-living objects like many of star clusters.
- Molecular clouds are located in the discs of spiral galaxies, and gas rich dwarf galaxies. They are absent (or scarce) in current lenticular (S0) and elliptical galaxies.

Collapsing clouds



(Miyama et al. 1987)

- Self-gravity is important for collapse of parts of the molecular cloud.
- The collapsing cloud forms an intricated network of sheets and filaments (Herschel space telecope dust emission).

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Collapsing clouds



Figure 1. Left: $Herselut/SPHEE 220 \mu m date continuum map of a portion of the Folder three$ transmost color of the 10 por blaces are part of the HGES survey (e.g. Mirille Docklems et al.<math>Herselvel tab) (e.g. Mirille Docklems) and the HGE survey (e.g. Mirille Docklems et al. Herselvel tab) (Source 1 at 2003). The contrast of the Ramanna has been valued in sting. Herselvel tab) (Source 1 at 2003). The contrast of the Ramanna has been valued as thing, are consistent starsform (2.5 Source 4 at 2003). The location of the Ramanna has been valued as thing, are started as the result of the Ramanna has been at the result of the Ramanna has been at the result has an evolution of the Ramanna has been at the Ramanna has b

(Andre 2013)

- The situation within a molecular clouds is very different from the idealised spherically symmetric models; flows along filaments.
- Filaments have some remarkable and still elusive properties. E.g. their width is almost universally 0.1pc.
- Their density is $\propto
 ho_0/(1+(r/r_0)^2).$
- The majority (\gtrsim 80 %) of YSOs are found in filaments; in the field is only their minority.
- Star clusters form at junctions of filaments.



Figure 8. Total (thermal + nonthermal) velocity dipersion versus central column density for a sample of 64 finaments in northy interstellar clouds (from Arzonanian et al. 2013). The horizontal dashed line shows the value of the thermal sound speed ~ 0.2 km/s for T = 10 K. The vertical gray band marks the blocher zone between thermally subcritical and thermally structure of the value $M_{\rm eff} \sim 10$ M s/p for T = 10 K. The blue sold line shows the best powerlaw fit $m_{\rm eff} \sim N_{\rm eff} = 30$ M s/p for T = 10 K. The blue sold line shows the box powerlaw fit $m_{\rm eff} \sim N_{\rm eff} = 30$.

(Andre 2013)

- Filaments are observed both in star forming clouds (Aquila rift) and quiescent clouds (Polaris).
- In both regions, the filaments have characteristic radius of 0.1pc; however only the ones in Aquila are star forming.
- The line density of filaments decides about their ability to form stars.

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Zooming to a filamentary hub



Fig. 1. Large scale view of the NGC1333 ridge in Breenes. The image represents the total column density map derived from our Herschel-Planck maps in order scales, including two comparison of equivalent column densities of A_{12}^{-2} and 10^{-2} , respectively. Offstes are referred to the maps center with coordinates (a_{12}^{-2}) maps $= 00^{22}$ /9709 A_{11}^{-2} 11 S^{-12}) in radio projection. The 4 anterpions around NGC1333 and their boundaries are average (contermined and A_{22000}^{-2}) result A_{2200}^{-2} (A_{2200}^{-2}) results of the content of the Coll 333 clump stated in this work, average (contermined and A_{22000}^{-2}) results of the content of the Coll 333 clump stated in this work.

(Hacar et al. 2017)

Zooming to a filamentary hub



Fig. 16. Core candidates identified in N₂H⁺ in NGC1333 (labels 1-50; see Table 5.5). The position of the 44 cores associated to fibers are indicated by open circles and are colos-coded similar to Fig. 12c. For efference, the fiber numbers are labeled using a larger fort. In addition, the 6 cores found in isolations are marked by white squares. Those cores with measure N ≥ 2.5 M₁₀ are highlighted with solal signations.

(Hacar et al. 2017)

- Filaments consist of fibres; fertile vs. sterile fibres.
- Protostars are found in fibres.
- The vas majority of the youngest protostars (class 0) are found in fibres, while older protostars (classes I-III) are found gradually more dispersed.

Embedded star clusters



- Contain protostars, often main sequence O and B stars but still contain gas.
- More massive stars reach the MS faster.
- Young objects (age \lesssim several Myr).
- Opaque in visible, but much more interesting in the infrared.
- At this time, planets already start forming around stars.



Core mass function halos histogram with error han) of the ~ 500 candidate prestata cores identified with *Herselob* in Aquilla (André et al. 2016). The IMF of single stars (corrected for binaries – e.g. Kronga 2001), the IMF of militple systems (e.g. Chaberr 2003), and the typical mass performs (GC Outney) (e.g. Kranne (e.g. Kronga 2001), the MMF of Market and the system stars performed (GC Outney) (e.g. Kranne (e.g. Kranne)), the particular stars of the system stars performed (GC Outney) (e.g. Kranne (e.g. Care)), it pasks at ~0.0 M_☉, chee to the Jeans mass within marginally critical filaments at $T \sim 10$ K (ef § 0.

(Alves et al. 2007; Ward-Thompson et al. 2007)

 Mass function tells us how many objects belongs to a given mass bin.

$$\xi(m) \equiv \frac{\mathrm{d}N}{\mathrm{d}m}$$

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- Core mass function is shifted by a factor of 3 towards higher masses \rightarrow might indicate that each core loses $\approx 2/3$ of its mass.
- One core might produce more than one star.

- Salpeter (1955) found that $\xi(m) \propto m^{-2.35}$ for stars with $4 M_{\odot} \lesssim m \lesssim 10 M_{\odot}$.
- This would indicate that there is infinite number of low mass objects \rightarrow cut-off.
- Some refinements came later: Massey 1998, Kroupa 1993, 2001, Chabrier 2005
- E.g. Kroupa 2001 suggests:

$$\frac{\mathrm{d}N}{\mathrm{d}m} = \begin{cases} m^{0.3}, & 0.01\mathrm{M}_\odot \lesssim m \lesssim 0.08\mathrm{M}_\odot \\ m^{1.3}, & 0.08\mathrm{M}_\odot \lesssim m \lesssim 0.5\mathrm{M}_\odot \\ m^{2.3}, & m \gtrsim 0.5\mathrm{M}_\odot. \end{cases}$$

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Mass	stellar type	number	mass
		fraction	fraction
$0.01 { m M}_\odot \lesssim m \lesssim 0.08 { m M}_\odot$	brown dwarfs	0.37	0.043
$0.08 { m M}_\odot \lesssim m \lesssim 0.5 { m M}_\odot$	M stars	0.48	0.28
$0.5 { m M}_\odot \lesssim {\it m} \lesssim 1 { m M}_\odot$	K stars	0.09	0.17
$1 { m M}_\odot \lesssim {\it m} \lesssim 8 { m M}_\odot$	G, F, A, late B	0.06	0.34
$\gtrsim 8 { m M}_{\odot}$	early B, O	0.004	0.17

- The mean mass of a star is $\langle m
 angle = 0.36 {
 m M}_{\odot}$
- The total stellar mass in objects of $m \lesssim 1 {
 m M}_{\odot}$ is almost the same as the total stellar mass in objects of $m \gtrsim 1 {
 m M}_{\odot}$.

Spectral Type	Mass (M_{\odot})	$\frac{\log \mathcal{N}_*}{(\mathrm{s}^{-1})}$	$\log \mathcal{N}_{\rm FUV} \ ({\rm s}^{-1})$
O4	70	49.9	49.5
O5	60	49.4	49.2
O6	40	48.8	48.8
07	30	48.5	48.6
08	23	48.2	48.4
09	20	47.8	48.2
B0	18	47.1	48.1
B1	13	45.4	47.5
B2	10	44.8	47.1

Table 15.1 Ultraviolet Radiation from Massive Stars

 N_* is the number of ionising photons (E > 13.6 eV), N_{EUV} is the number of EUV photons (6 eV < E < 13.6 eV). (Stahler & Palla 2005)

- Early feedback (occurs immediately after star formation): Photoionising, stellar winds, radiation pressure.
- Late feedback (when the star leaves the main sequence): red supergiant winds, AGB winds, SNe

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Simple picture of early feedback bubbles: Photoionising radiation



- Photoionising feedback. The bubbles are at temperature of $\approx 10^4 {\rm K}$ due to an efficient thermostat.
- Cooling via collisional excitation of ions of heavier elements (mainly oxygen, smaller contribution from neon and nitrogen).
- Simple structure: Shock front and ionisation front.

Simple picture of early feedback bubbles: Stellar wind



- Free wind region; reverse shock; contact discontinuity; shock front.
- The hot bubble in wind feedback reach temperatures $\gtrsim 10^7 {\rm K}.$
- Numerical models indicate that the photoionising feedback impacts the interstellar medium more than stellar wind feedback.

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- Energy released by stars of given mass per $1 M_{\odot}$ of coeval stellar population.
- More massive stars are more numerous, ut they produce the most photons and wind.
- With demise of the most massive stars, feedback energy drops abruptly.

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Real picture of early feedback



H II region NGC 602

- Don't be confused by distant galaxies.
- Cluster of young stars containing massive stars.
- Gas is being expelled.
- Sharp ionisation front.
- Structures (cones) in the ionisation front.

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- Each SN of type II releases suddenly $\approx 10^{51} {\rm erg.}$
- SNe are important for shaping the structure of galactic disc, regulate formation of molecular clouds, create the hot phase.
- SNe launch galactic outflows, fountains.
- SNe occur after 3 Myr; the life-time of the most massive stars.
- At the time when SNe occur, embedded clusters have already been cleared by early feedback.

Open Star clusters



Open star cluster M 44.

- After several Myr, the cluster is gas free.
- Most of open star clusters dissolve early, only 10% of open star clusters reach the age of 10 Myr and 4% reach the age of 100 Myr.
- However, there are some old open star clusters, age up to 10 Gyr (NGC 188, M 67).
- Occupy the disc of the Galaxy.

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Open Star clusters

Figure 2

Hubble Space Telescope optical/IR image of the dense massive young cluster R136/30Dor (courtesy of M.J. McCaughrean: FOV ~30 arcsec × 30 arcsec or 7.5 pc × 7.5 pc). Dozens of massive O stars are found crowded within the half-light radius of 2 pc (Brandl et al. 1996). (a) A VLT image of NGC 3603 (Brandl et al. 1999) and (b) a VLT image of the Tranezium Cluster in Orion (McCaughrean 2001) are shown as these two galactic clusters would be seen if they were located at the distance of R136 in the Large Magellanic Cloud (50 kpc) and imaged with similar angular resolution (see Zinnecker 2002).



Zinnecker & Yorke 2007

- Comparison of sizes of open star clusters.
- (b) ONC mass $\sim 10^3 {\rm M}_\odot$, (a) NGC 3603 - mass $\sim 10^4 {\rm M}_\odot$, background: R 136 (in LMC) - mass $\sim 10^5 {\rm M}_\odot$.
- Mass ranges from tens of M_{\odot} to 10^{6} or $10^{7}M_{\odot}$ in star bursting galaxies (e.g. Antennae).
- Radii are several pc.
- Metallicity is approximately solar Z_☉.

Open Star clusters



Johnson et al. 2017 (for M 31)

- Cluster initial mass function has a slope of ≈ -2 , which is similar to the IMF.
- Slope -2: the same mass of objects in each logarithmic mass bin.
- The most massive clusters are underrepresented.

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More massive open star clusters



Figure 3: Top view of the Milky Way Galaxy, with the known spiral pattern (Valké 2008) and young star clusters more massive than $\gtrsim 10^4 M_{\odot}$ identified. Basic cluster parameters are listed in Tab.21 The location of the Sun is indicated by \odot , and its orbit by the dotted circle. Dashed lines indicate circles of 1 kpc and 2 kpc around the sun.

(Zwart et al. 2010)

- Distribution of massive open star clusters within the Galaxy.
- The most massive open star clusters currently forming in the Galaxy are of the mass $1-2.10^4 M_{\odot}$.
- Clusters are at the tip of the bar, galactic centre, but also in the disc at relatively large $R_{\rm gal}$.

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Massive stars in star clusters



Weidner et al. 2013

- Massive stars are usually found only in the most massive star clusters.
- Perhaps, massive stars need special (high density, pressure) environment to form.
- Massive stars are often found mass segregated, i.e. more massive stars are found closer to the cluster centre.

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• Mass segregation can be primordial or dynamical.

Theories for open star cluster formation



Figure 7: Deviation for the "extended system A Bill (b) ~ 5 yes, tor Table 3). An expected, the infall of the sub-transmission mech mass adveg than this for the comparison of Fig. 50. The present of the system is all highly advanced 100 km in 4 × 20 per label and a system of the present of

(Banerjee & Kroupa 2015)

- Formation of star clusters is not well understood.
- Two classes of models: coalescence of subclusters; monolithic scenario.
- Coalescence of subclusters leaves signatures in the merged objects, and constraints the possible compactness of subclusters.

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Theories of open star cluster formation

- Competitive accretion (in monolithic star cluster formation; Bonnell et al. 1997).
- Consider a gaseous cloud containing protostars, each protostar is of the same mass.
- Some protostars are located at denser medium, move at different relative velocities \rightarrow different accretion rates.
- Protostars accreting more lose their momentum → sink towards the cluster centre → are located at denser medium → accrete at even higher rate.
- Competitive accretion naturally explains mass segregation as most massive stars are slowed down at the cluster centre.

Binary stars in open star clusters



The Trapezium in the Orion Nebula Cluster (ONC)

- Multiplicity increases with the mass of the primary: O stars -2.1; early B stars - 1.6; A stars -0.9; G stars - 0.5.
- More massive stars are found in tighter binaries.
- The presence of binary stars is important dynamically.

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Binary stars in open star clusters



The distribution of orbital periods for stars of various masses (spectral types).

- Multiplicity increases with the mass of the primary: O stars -2.1; early B stars - 1.6; A stars -0.9; G stars - 0.5.
- More massive stars are found in tighter binaries.
- The presence of binary stars is important dynamically.

- Gravitationally unbound objects, usually expanding.
- Typical size 50pc.
- Contain O, B stars, but also lower mass stars (T associations).
- Two possibilities of their origin: expanding outer halos of open star clusters after gas expulsion; formed throughout a molecular cloud, were never gravitationally bound together.
- OB associations are associated with the so called runaway phenomenon → almost any association has some runaway stars.

Globular star clusters

- Very different from open star clusters.
- The SFR per unit volume in the cosmic history peaked at z = 2 (10 Gyr ago).
- Very old objects, age pprox 12 Gyr.
- $\bullet\,$ Typical mass is 10^4 to $10^6 {\rm M}_\odot.$
- Almost spherically symmetric, of coeval age.
- But not that "simple" might not be appropriate: multiple stellar populations; their formation.
- \bullet Dynamically old ($\sim 10^4$ orbits), they provide nice examples for N-body problem.

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- Located in galactic halo, often at eccentric orbits.
- Denser (radii $pprox 3 \mathrm{pc}$) than open star clusters.
- Of subsolar metallicity (typically Z = -2 to Z = -0.8).

Nuclear star clusters



- Very different from the previous two cases.
- Very dense systems $\sim 10^6 {\rm M}_\odot {\rm pc}^{-3}$,old.
- In some cases, there is possible interaction with the supermassive black hole.
- Hosted mostly by galaxies of stellar mass 10^8 to $10^{10} {\rm M}_\odot.$

Neumeyer et al. 2020

Illustration of thermodynamics of self-gravitating systems

- Consider a cluster composed of *N* stars, each of the same mass.
- A temperature at position r is defined as

$$\frac{1}{2}m\langle v^2(r)\rangle = \frac{3}{2}k_BT(r) \tag{1}$$

• The mean temperature of the cluster is then

$$\overline{T} = \frac{\int 4\pi r^2 \rho(r) T(r) dr}{\int 4\pi r^2 \rho(r) dr}$$
(2)

• From the virial theorem (2K + W = 0), it follows

$$E = -\frac{3}{2}Nk_{B}\overline{T}.$$
(3)

Illustration of thermodynamics of self-gravitating systems

• This implies that the heat capacity is

$$C \equiv \frac{\mathrm{d}E}{\mathrm{d}\overline{T}} = -\frac{3}{2}Nk_B. \tag{4}$$

- The heat capacity is negative.
- If we transfer a small amount of heat from the system, the system gets hotter → releases heat → gets hotter further.
- There is no equilibrium in such a simple model.
- Binary stars stop the collapse.
- Evolved clusters tend to form bound cores with high velocity dispersion enveloped in a low density halos.

- Relevant for open star clusters (some similarities to globular star clusters, but we do not know how the initial conditions of their formation).
- Early gas expulsion.
- Evaporation and ejections.
- Core collapse.
- Dissolution in the external field of the galaxy.
- Characteristic time-scales are usually expressed in the relaxation time

$$t_{\rm rlx} = \frac{0.34\sigma^3}{G^2 m \rho \log \Lambda}.$$
 (5)

- Large difference between the adiabatic and impulsive approximation, they depend on the time-scale.
- $t_{\rm cross} < t_{\rm ge} \rightarrow$ adiabatic. Adiabatic invariants (e.g. $aM_{\rm cl}$) are conserved.
- During impulsive gas expulsion, star clusters are typically impacted more.
- Gas expulsion leads to cluster expansion, and many stars (approx. 1/2 to 2/3) are lost for the star formation efficiency of 30%.

- Core collapse occurs at typically $16t_{\rm rlx}$.
- Clusters are dissolved at typically $20 60t_{\rm rlx}$, depending on the strength of the external field of the galaxy.

$$r_{\rm t} \simeq R_{\rm gal} (M_{\rm cl}/3M_{\rm gal})^{1/3}.$$
 (6)

- The presence of the external gravitational field facilitates cluster evaporation.
- During evaporation, clusters lose preferentially lower mass stars, they tend to retain compact remnants (WDs, NSs).

Runaway stars



- Runaway stars move at velocity $\gtrsim 30 {\rm km} \, {\rm s}^{-1}$.
- Can be produced by two mechanisms: Close interaction between three or more stars; binary supernova scenario.
- In the binary supernova scenario, the observed star is the former secondary in a binary. When the SN exploded, it decreased suddenly its mass, and also the NS might get a kick.

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Runaway stars originating from the ONC.